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# Micromagnetic Dynamics in Thin Films

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# Web Pages

- Home Page:

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- OOMMF:

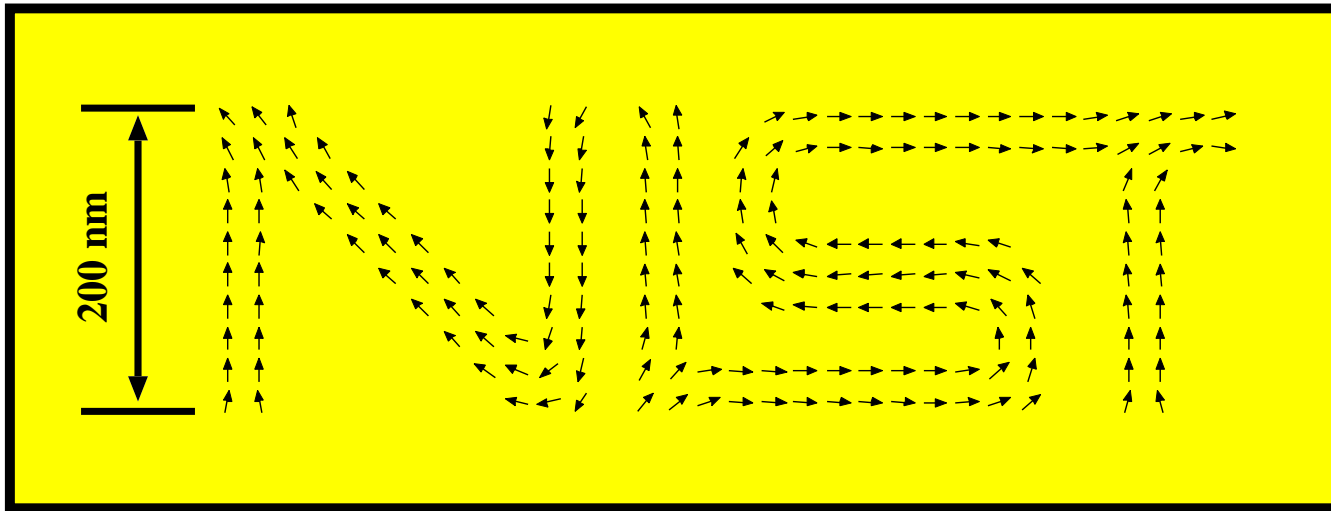
<http://math.nist.gov/oommf/>

- $\mu$ MAG:

<http://www.ctcms.nist.gov/~rdm/mumag.org.html>

# Outline

- Background
- $\mu$ MAG Standard Problems
- High Speed Switching (Std. Prob. 4)
- Dynamics Near Quasi-Static Coercive Field
- Dynamics at Remanence



**Micromagnetics is the study, modeling and simulation of magnetic materials and their behavior at the sub-micron scale.**

# Micromagnetic Equations

Landau-Lifshitz-Gilbert:

$$\frac{d\mathbf{M}}{dt} = -\gamma\mathbf{M} \times \mathbf{H}_{\text{eff}} - \frac{\gamma\alpha}{M_s}\mathbf{M} \times (\mathbf{M} \times \mathbf{H}_{\text{eff}})$$

$$\mathbf{H}_{\text{eff}} = -\frac{1}{\mu_0} \frac{\partial E_{\text{density}}}{\partial \mathbf{M}}$$

Energies:

$$E_{\text{total}} = E_{\text{exchange}} + E_{\text{anisotropy}} + E_{\text{Zeeman}} + E_{\text{demag}}$$

# Micromagnetic Energy Equations

$$E_{\text{exchange}} = \int_V \frac{A}{M_s^2} (|\nabla M_x|^2 + |\nabla M_y|^2 + |\nabla M_z|^2) d^3 r$$

$$E_{\text{anisotropy}} = - \int_V \frac{K_1}{M_s^2} (\mathbf{M} \cdot \mathbf{u})^2 d^3 r$$

$$E_{\text{Zeeman}} = -\mu_0 \int_V \mathbf{M} \cdot \mathbf{H}_{\text{applied}} d^3 r$$

$$E_{\text{demag}} = -\frac{\mu_0}{2} \int_V \mathbf{M} \cdot \mathbf{H}_{\text{demag}} d^3 r$$

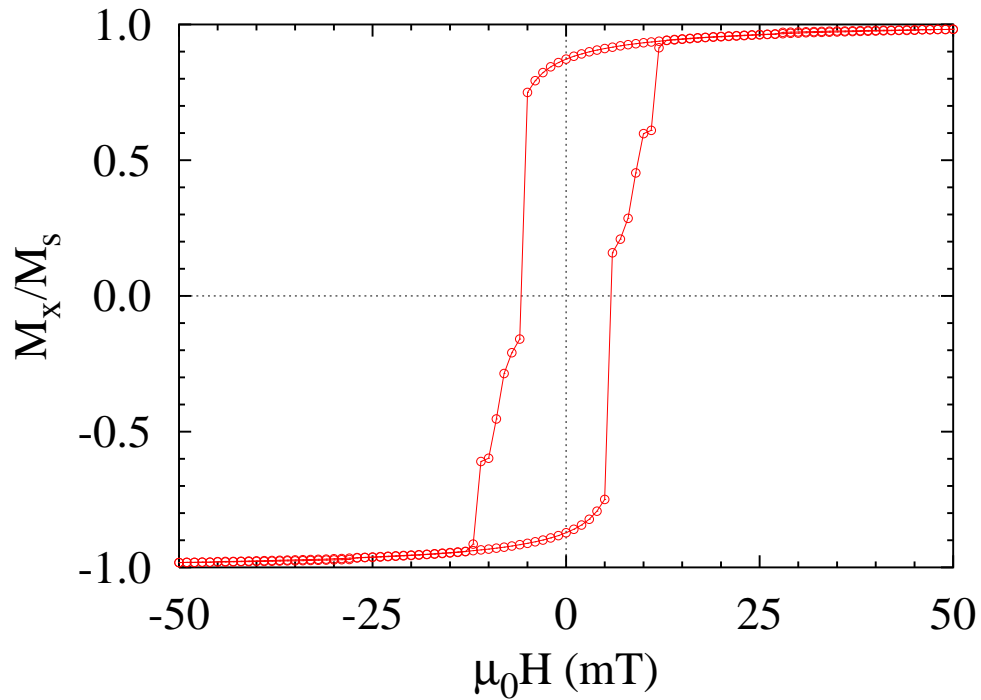
$$\begin{aligned} \text{where } \mathbf{H}_{\text{demag}}(r) = & -\frac{1}{4\pi} \int_V \nabla \cdot \mathbf{M}(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|^3} d^3 r' \\ & + \frac{1}{4\pi} \int_S \hat{\mathbf{n}} \cdot \mathbf{M}(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|^3} d^2 r' \end{aligned}$$

## Comparative Material Constants

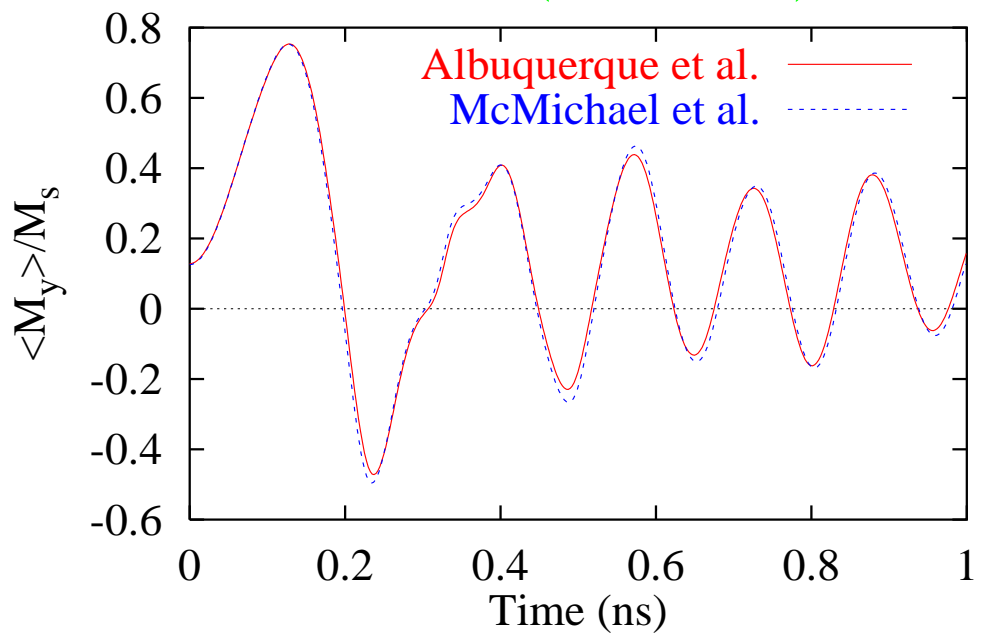
Material	$M_s$ (kA/m)	$K_1$ (kJ/m <sup>3</sup> )	$A$ (pJ/m)	$l_{\text{ex}}$ (nm)
Co	1400	520	30	4.9
Fe	1700	48	21	3.4
Ni	490	-5.7	9	7.7
Permalloy	860	0	13	5.3
Nd <sub>2</sub> Fe <sub>14</sub> B	1280	4500	13	3.5

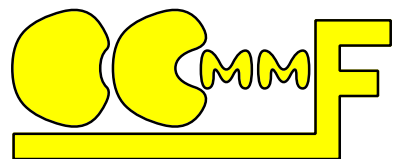
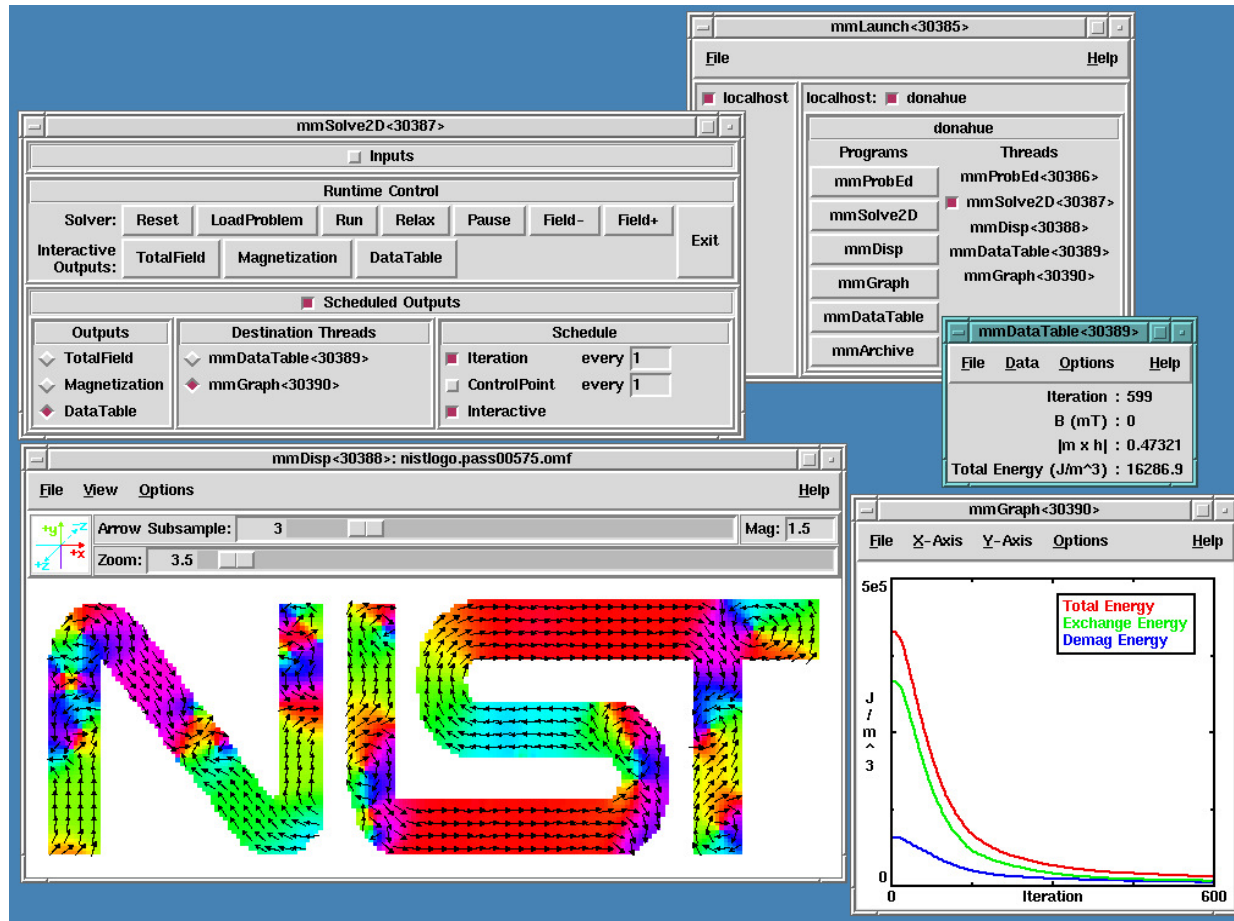


## Hysteresis Loops (quasi-static)



## Time Series (dynamic)

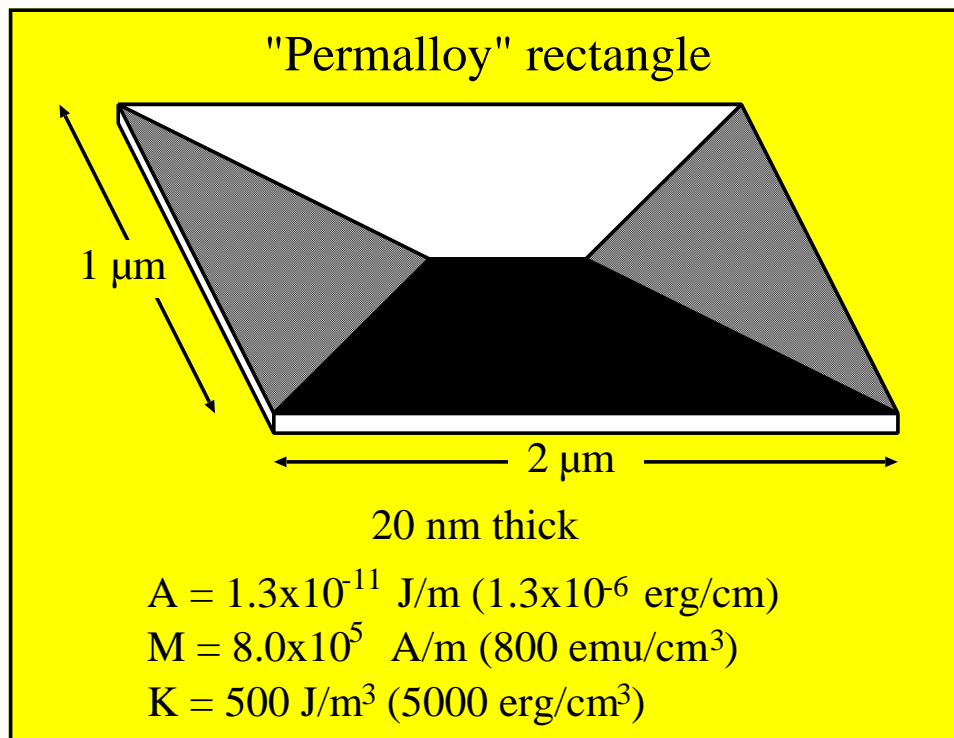




Public Code

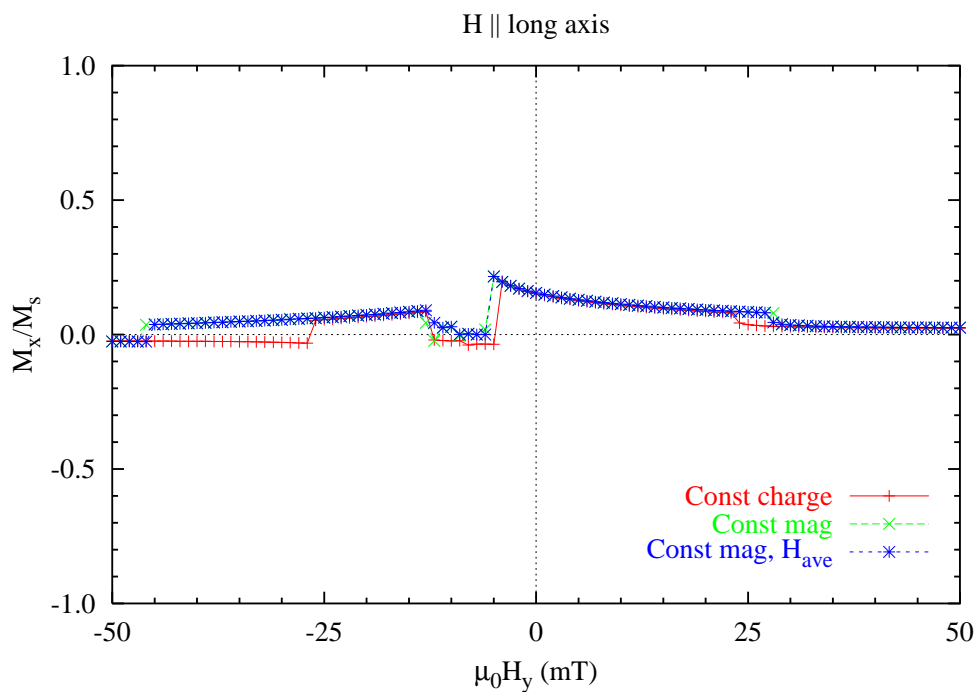
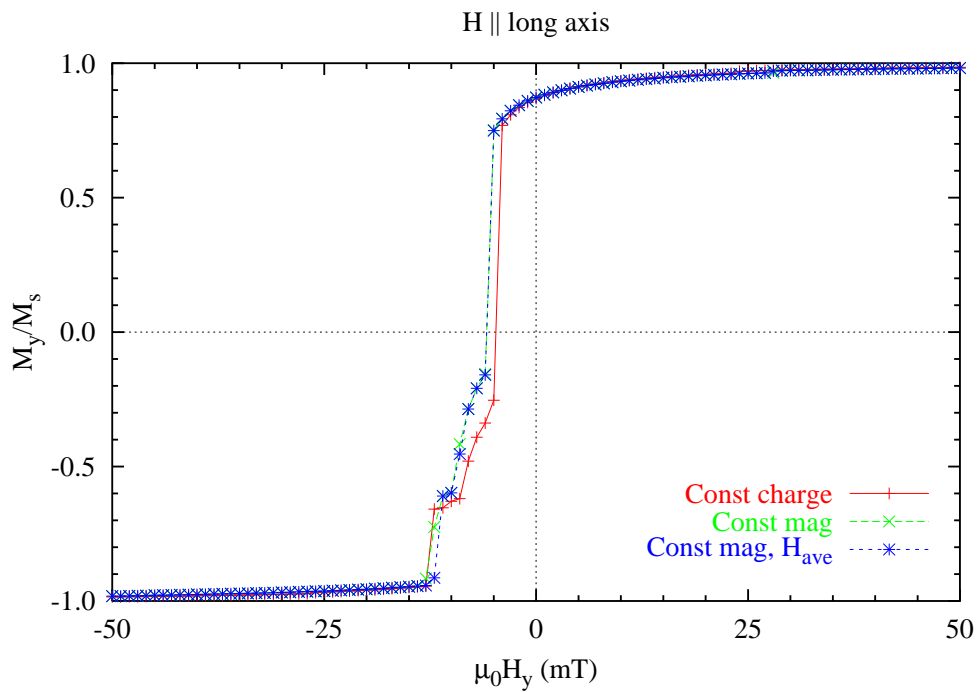
# $\mu$ MAG Problem #1

## Specification



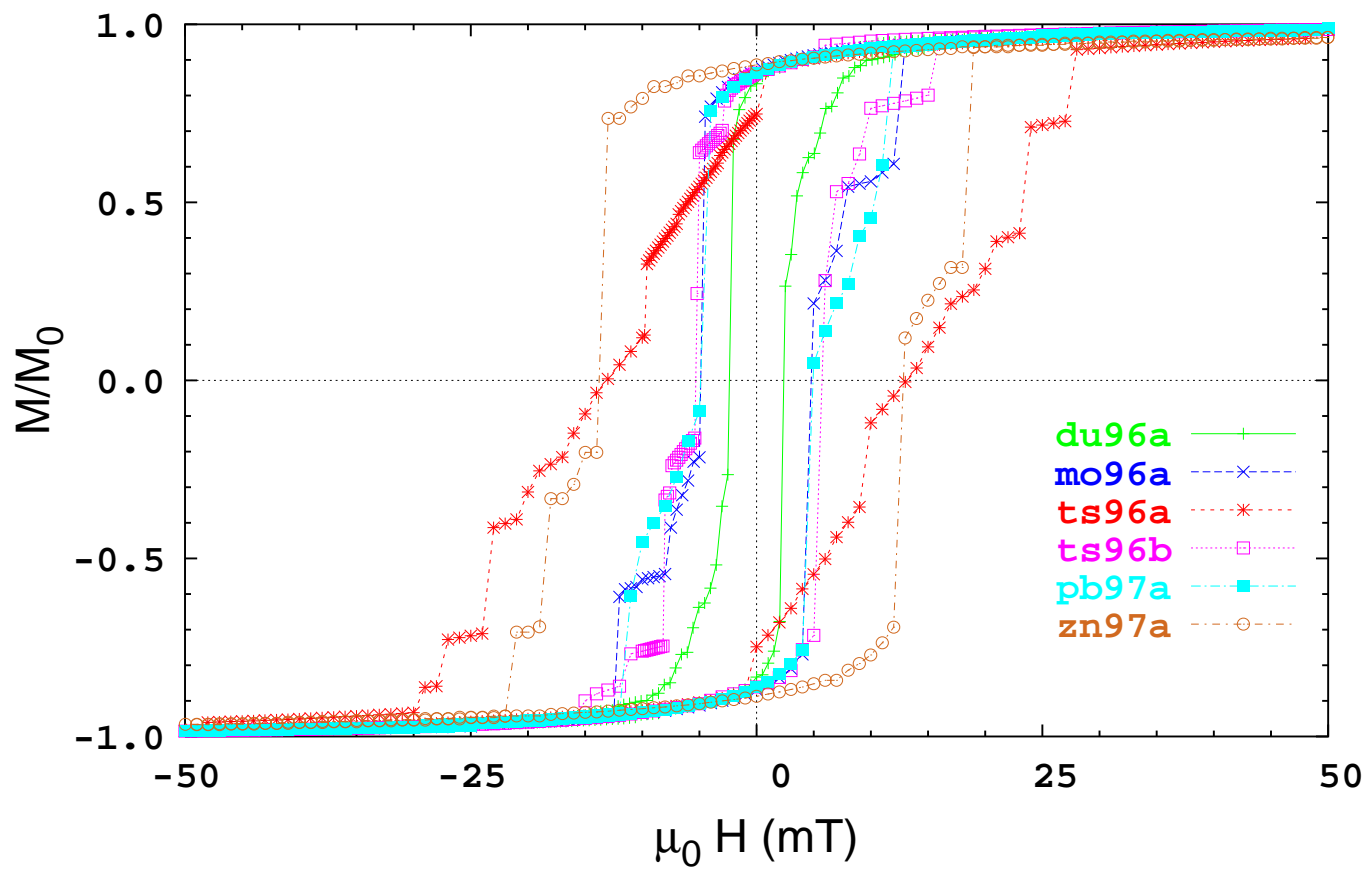
# $\mu$ MAG Problem #1

## OOMMF Results (20 nm cells)

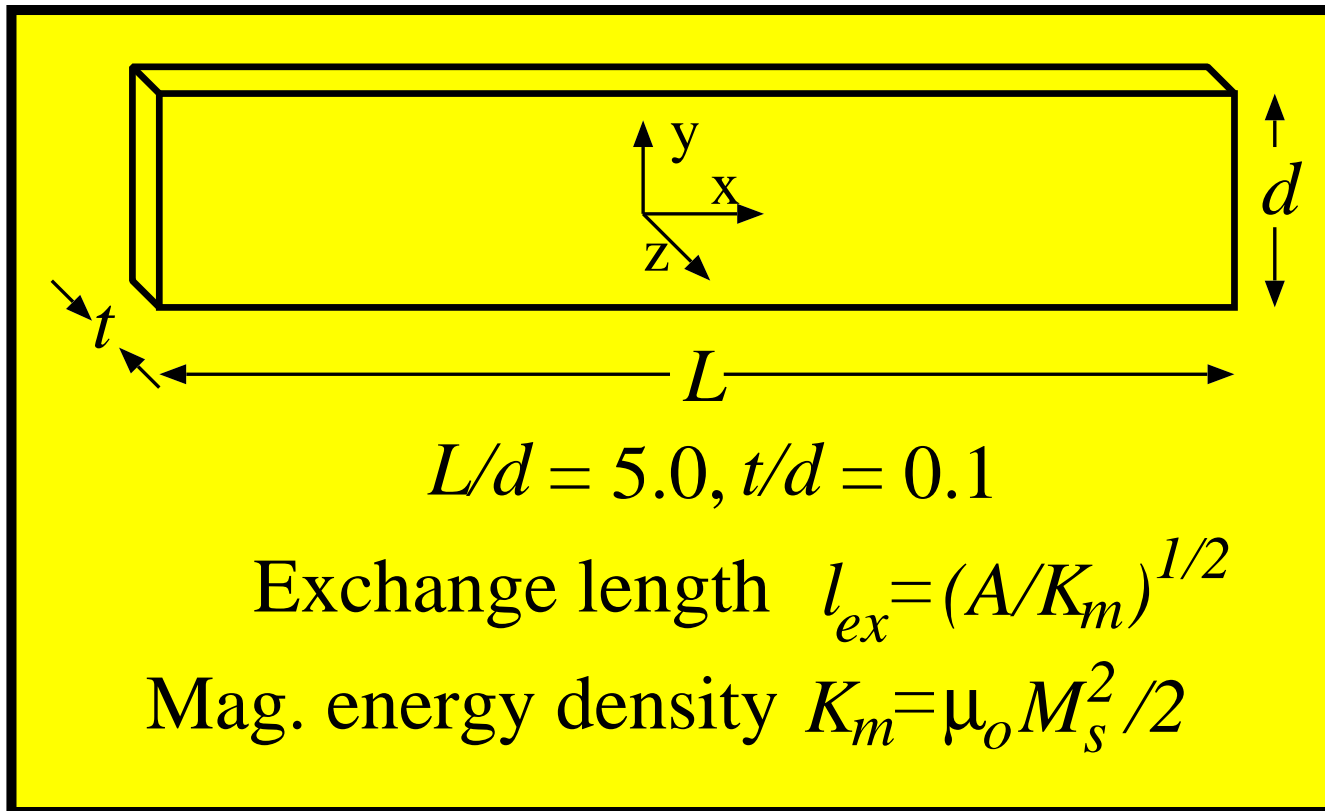


# $\mu$ MAG Problem #1 Results

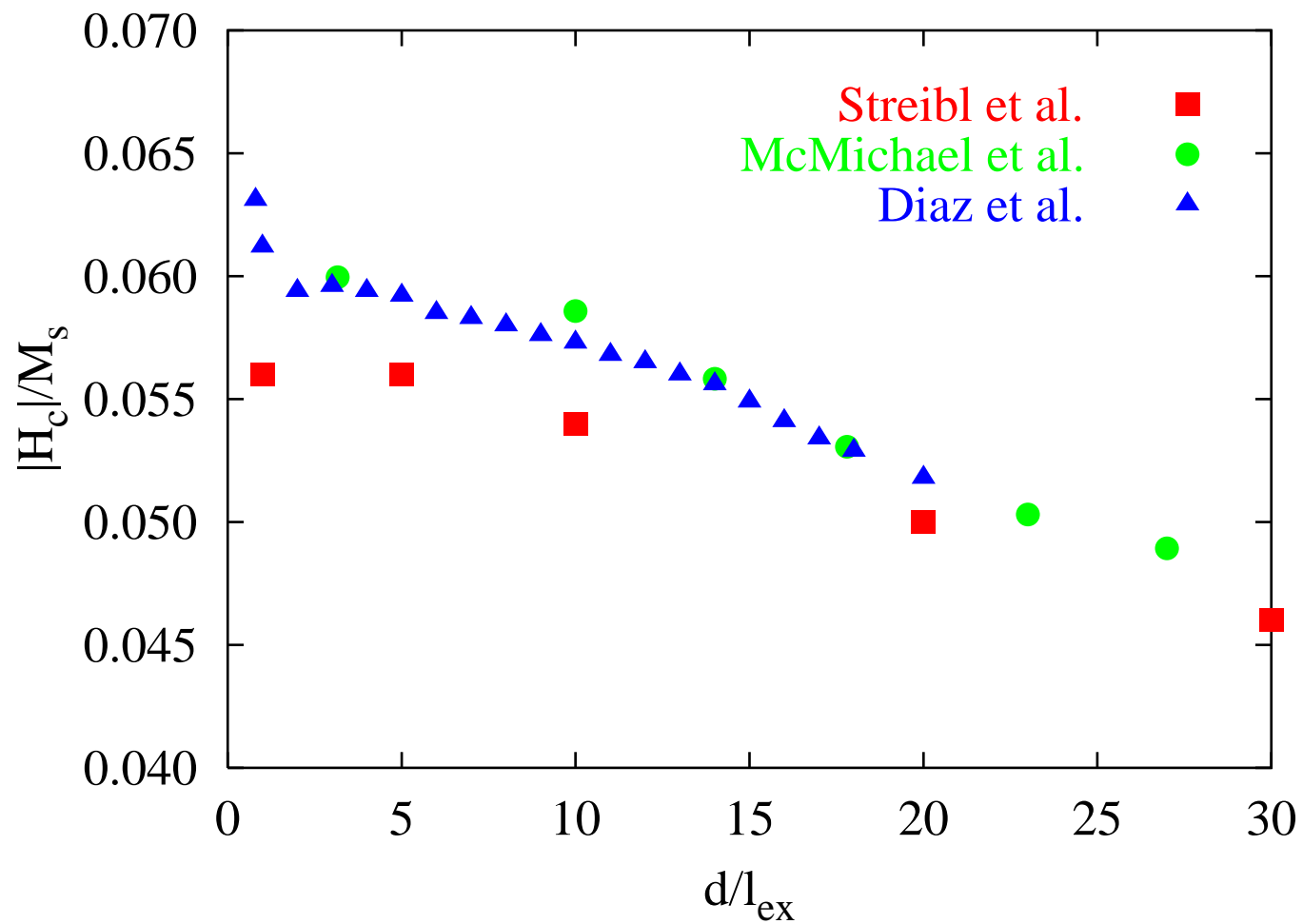
H || y, y-component



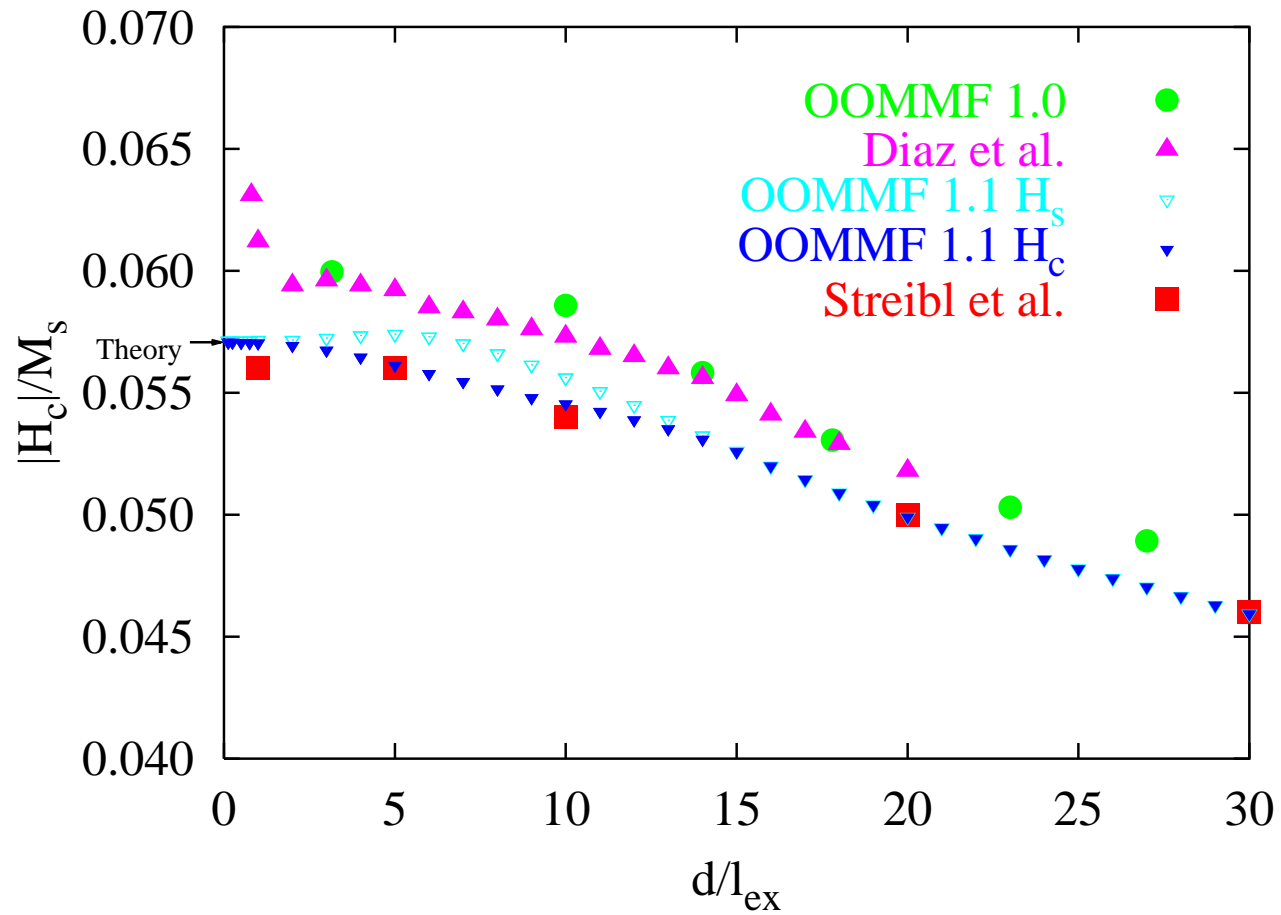
## $\mu$ MAG Problem #2 Specification



## $\mu$ MAG Problem #2, Coercivities

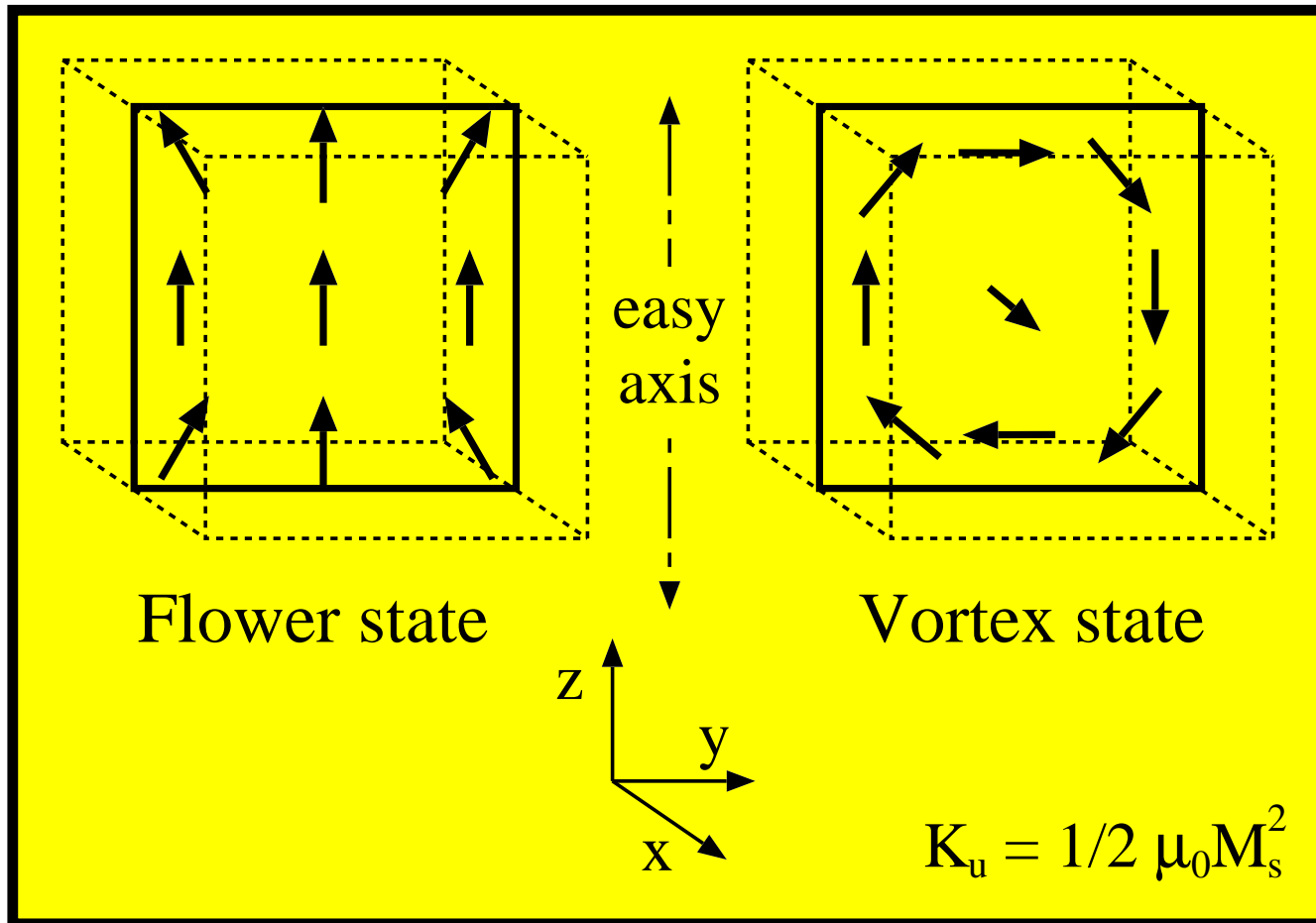


## $\mu$ MAG Problem #2, Coercivities

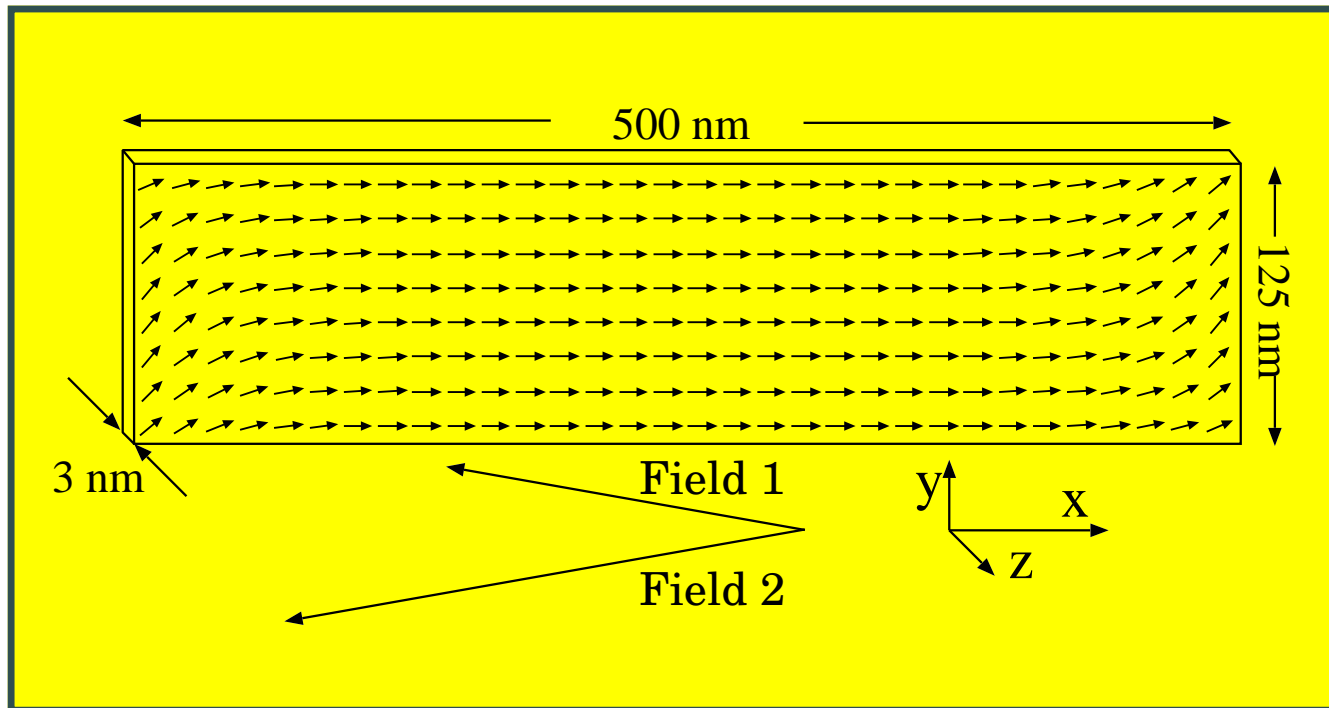




# $\mu$ MAG Problem #3 Specification



## $\mu$ MAG Problem #4 Specification



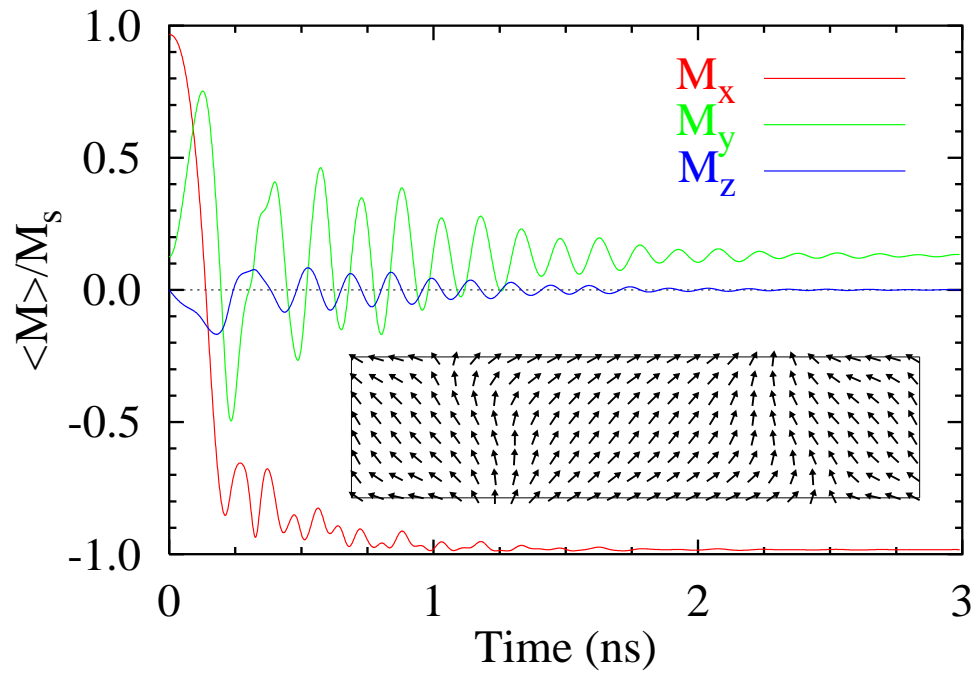
Permalloy parameters

Field 1:  $\mu_0 H = 25$  mT at  $170^\circ$

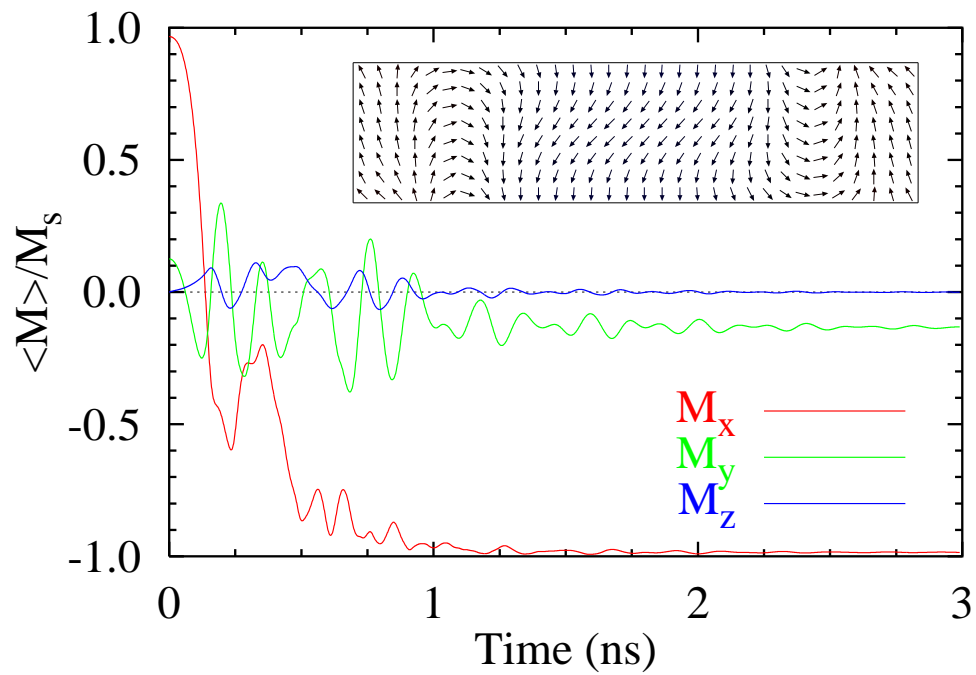
Field 2:  $\mu_0 H = 36$  mT at  $190^\circ$

(Fields are approximately 1.5 times the quasi-static coercivity.)

## Field 1, 170°

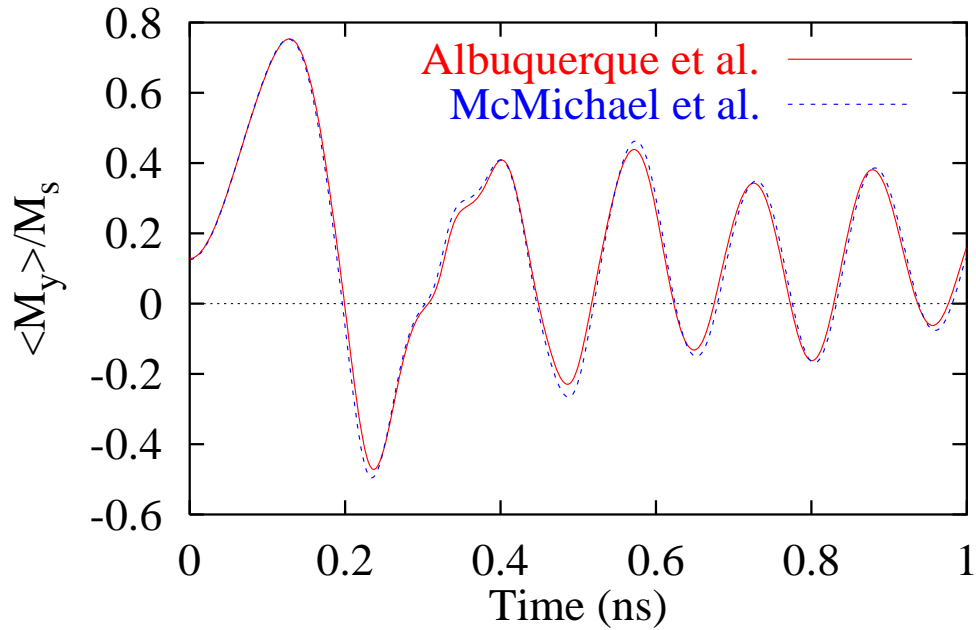


## Field 2, 190°

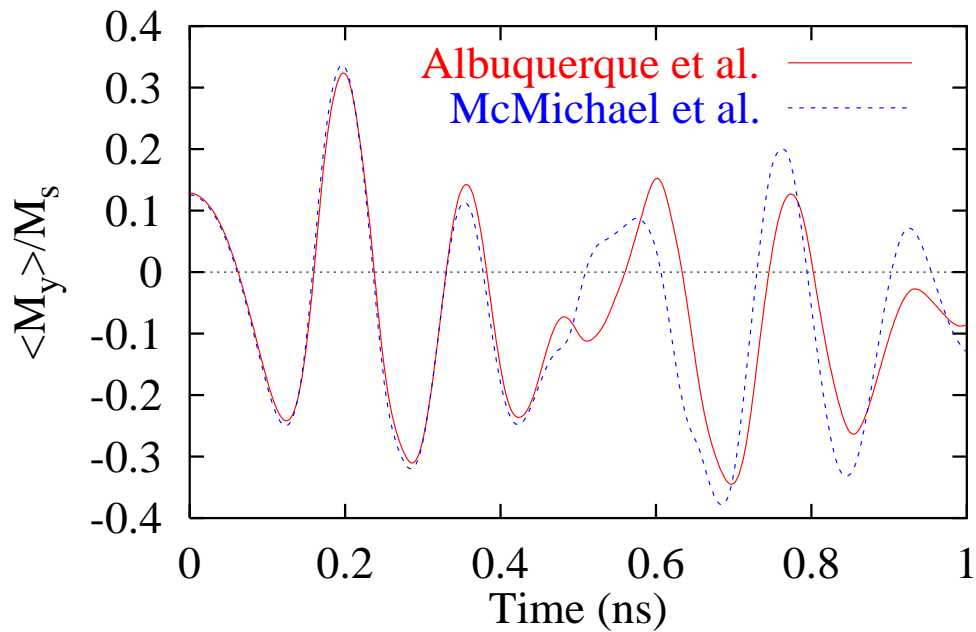


# Prob. 4 Comparisons

Field 1 ( $170^\circ$ )



Field 2 ( $190^\circ$ )



## Max angle data

- Field 1, 2.5 nm cells, 8-dot exchange:  $22.3^\circ$
- Field 2, 3.125 nm cells, 8-dot exchange:  $118^\circ$
- Field 2, 2.5 nm cells, 8-dot exchange:  $71.9^\circ$
- Field 2, 3.125 nm cells, 4-ang exchange:  $48.1^\circ$

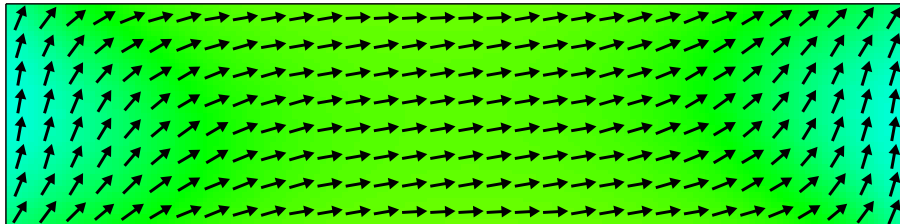
Exchange reference:

**Exchange Energy Representations  
in Computational Micromagnetics,**  
M. J. Donahue and R. D. McMichael,  
*Physica B*, **233**, pp 272-278 (1997).

# Quasi-Static Equilibria

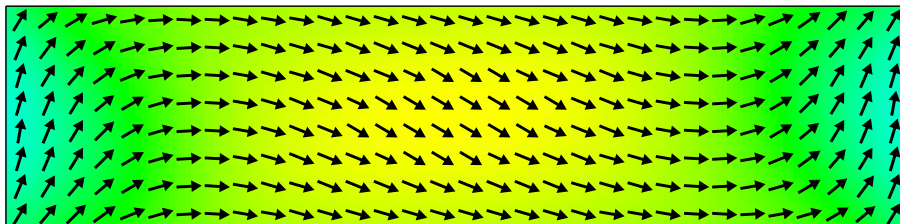
Field 1,  $H_{\text{app}} = 15 \text{ mT}$

$\mu_0 H = 15 \text{ mT}$   
x



Field 2,  $H_{\text{app}} = 21.6 \text{ mT}$

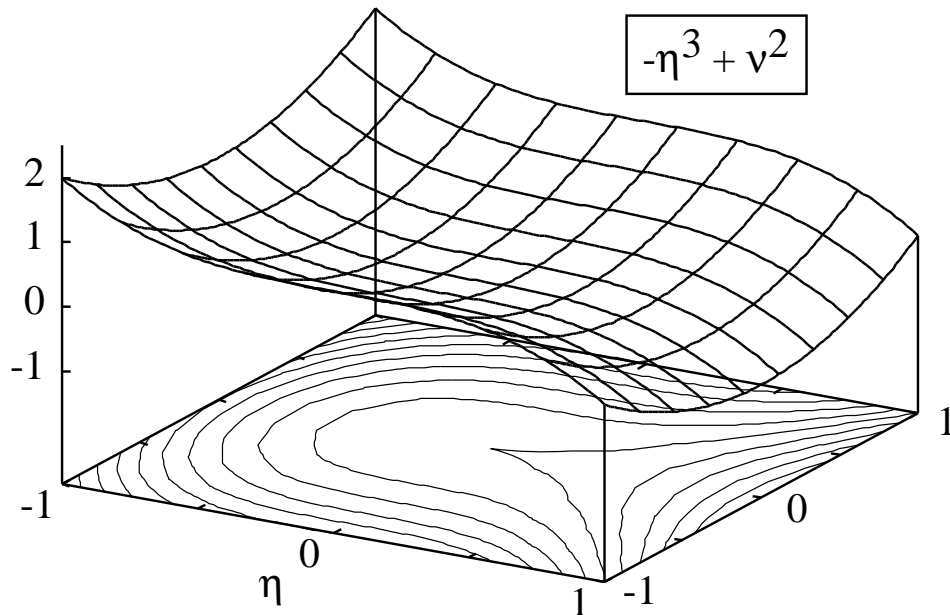
$\mu_0 H = 21.6 \text{ mT}$   
x



# Single Spin: Energy Surface

At critical field  $\mathbf{H}_c$  assume

$$E_c = -U\eta^3 + K\nu^2$$



Let  $\mathbf{h} = \mathbf{H} - \mathbf{H}_c$ , then

$$E = -U\eta^3 + K\nu^2 - \mu_0 M \left[ h_\eta \eta + h_\nu \nu + h_\zeta \left( 1 - \frac{1}{2}\eta^2 - \frac{1}{2}\nu^2 \right) \right]$$

and

$$\eta_{\text{eq}} = \frac{\mu_0 M}{6U} \left[ h_\zeta - \sqrt{h_\zeta^2 - \frac{12U h_\eta}{\mu_0 M}} \right].$$

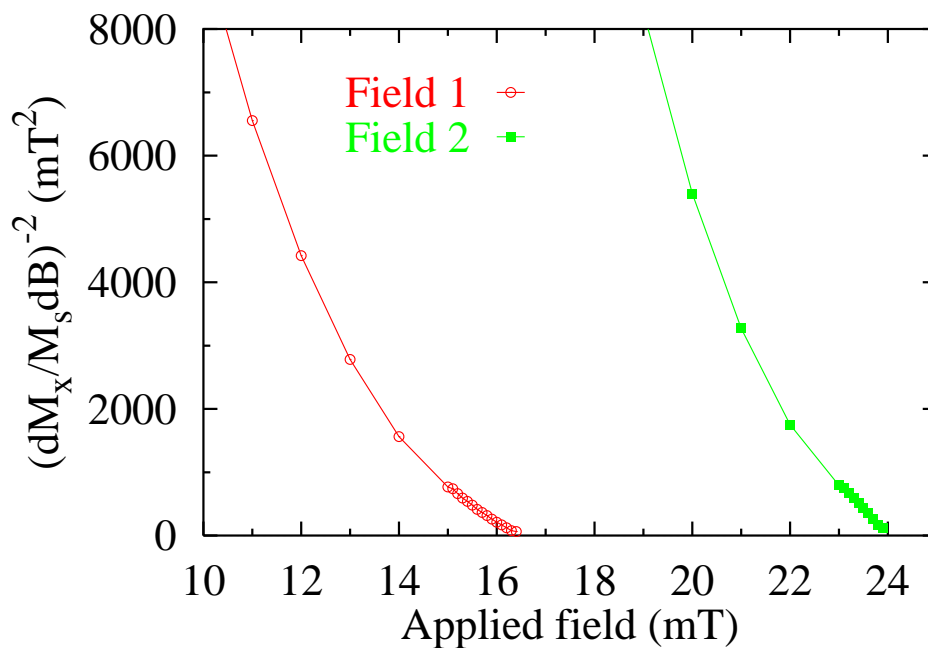
# Critical Field Susceptibility

Divergent susceptibility  $\chi_{\eta,\eta}$  is given by

$$\chi_{\eta,\eta} = M \frac{\partial \eta_{\text{eq}}}{\partial h_{\eta}} = M \left[ h_{\zeta}^2 - \frac{12U}{\mu_0 M} h_{\eta} \right]^{-1/2}.$$

If  $h_{\zeta}^2 \ll 12U h_{\eta} / \mu_0 M$ , then

$$\chi_{\eta,\eta} \propto h_{\eta}^{-1/2} \quad \text{as } |\mathbf{h}| \downarrow 0.$$



See also: A. Hubert and W. Rave, *Physica Status Solidi B*, **211**, pp 815–829 (1999).



# Ring Down Frequency

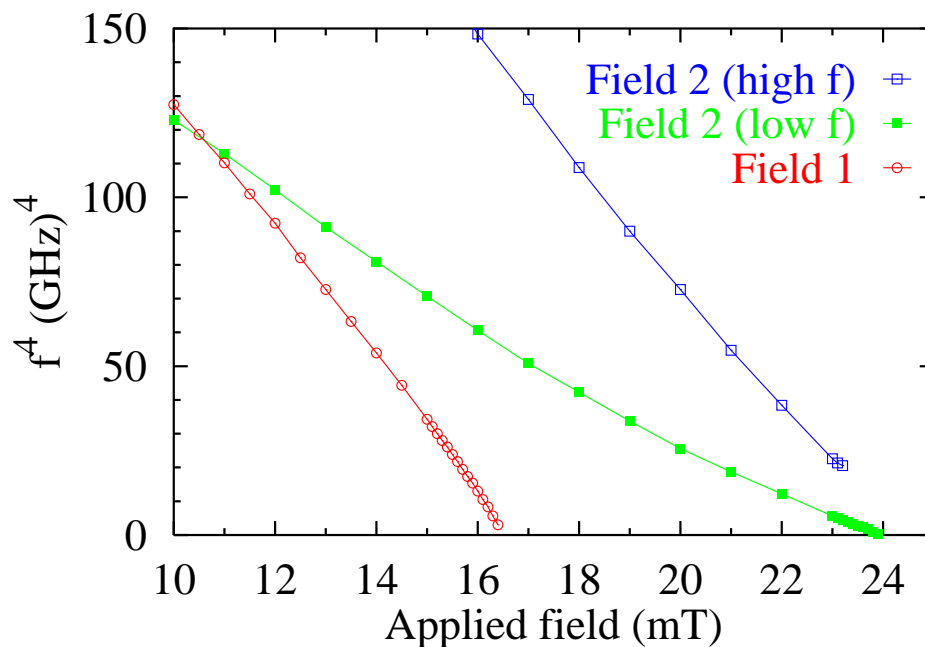
Precession frequency  $f$  is given by

$$f = \frac{\gamma}{2\pi\mu_0 M} \left[ \frac{\partial^2 E}{\partial \eta^2} \frac{\partial^2 E}{\partial \nu^2} - \left( \frac{\partial^2 E}{\partial \eta \partial \nu} \right)^2 \right]^{1/2} \Bigg|_{\eta_{\text{eq}} \nu_{\text{eq}}}$$

$$= \frac{\gamma}{2\pi} \left[ h_\zeta + \frac{2K}{\mu_0 M} \right]^{1/2} \left[ h_\zeta^2 - \frac{12U}{\mu_0 M} h_\eta \right]^{1/4} .$$

If  $h_\zeta^2 \ll 12U h_\eta / \mu_0 M$ , then

$$f \propto h_\eta^{1/4} \quad \text{as } |\mathbf{h}| \downarrow 0.$$



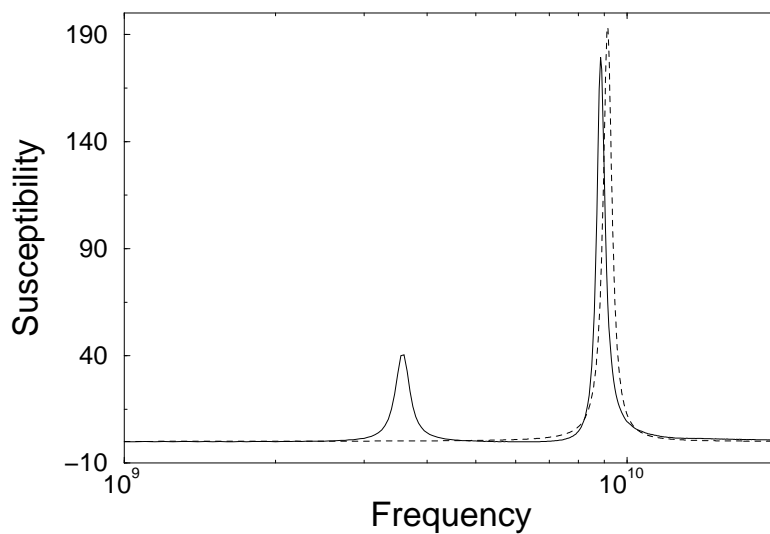
# High Frequency Susceptibility

Define susceptibility  $\chi$  by

$$\langle \mathbf{M}(t) \rangle \cdot \mathbf{u} = \int_{-\infty}^{+\infty} \chi(t - t') [\mathbf{h}(t') \cdot \mathbf{u}] dt',$$

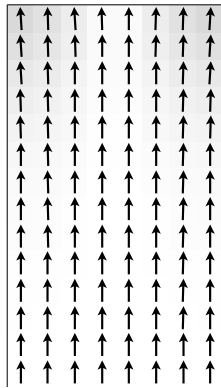
or in Fourier domain

$$M(t) = \chi(t) \star h(t) \leftrightarrow M(\omega) = \chi(\omega) \cdot h(\omega).$$

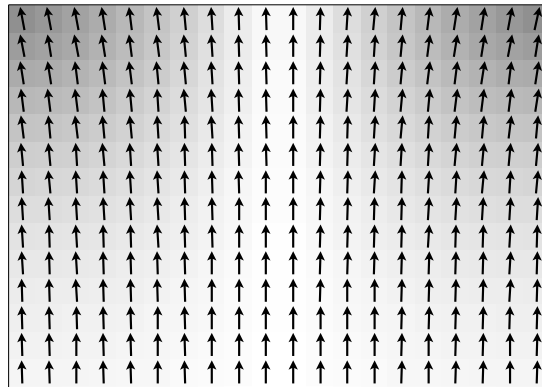


Permalloy stripe,  $1 \mu\text{m} \times 50 \text{ nm} \times 5 \text{ nm}$ .  
Solid line is micromagnetic simulation.

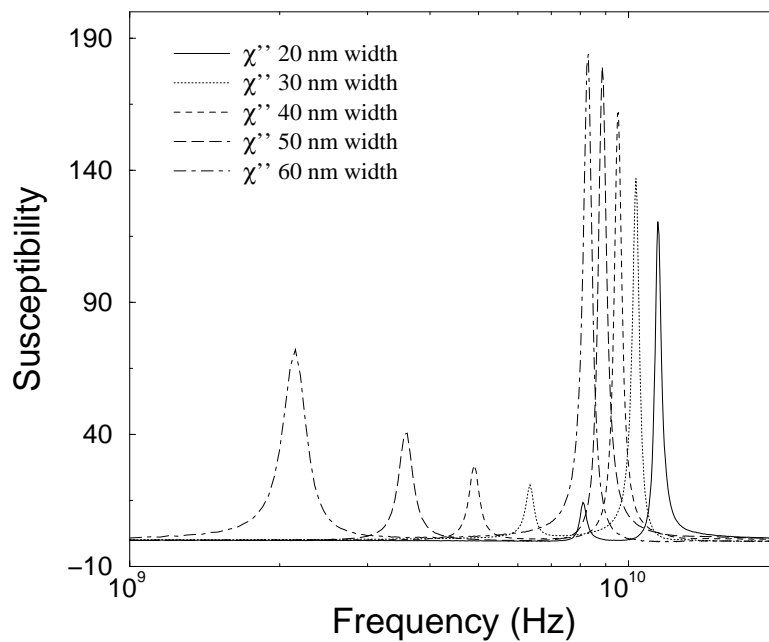
# Stripe Width Effects



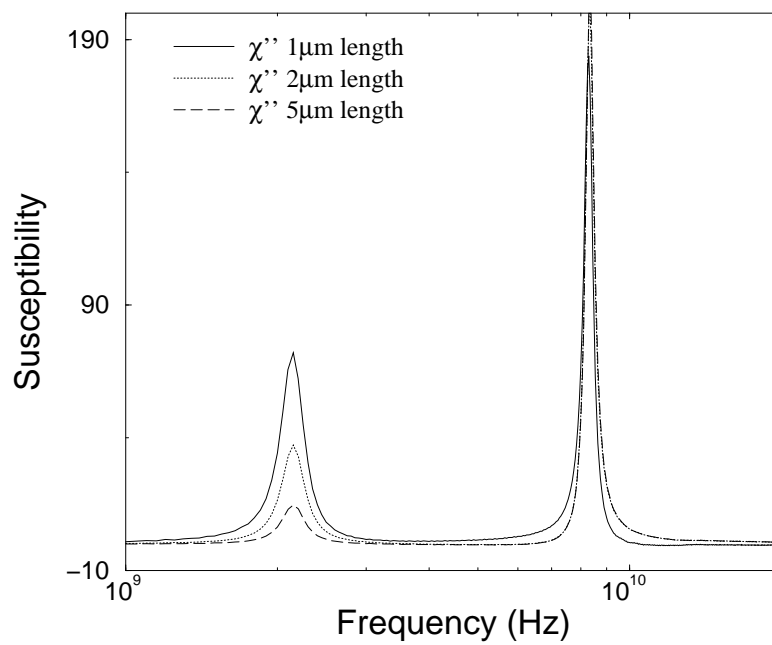
20 nm



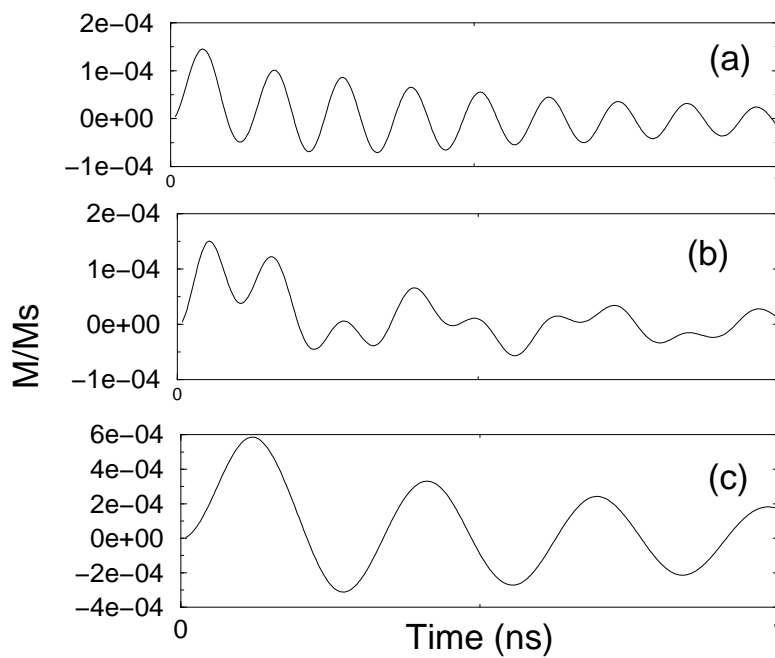
50 nm



# Stripe Length Effects

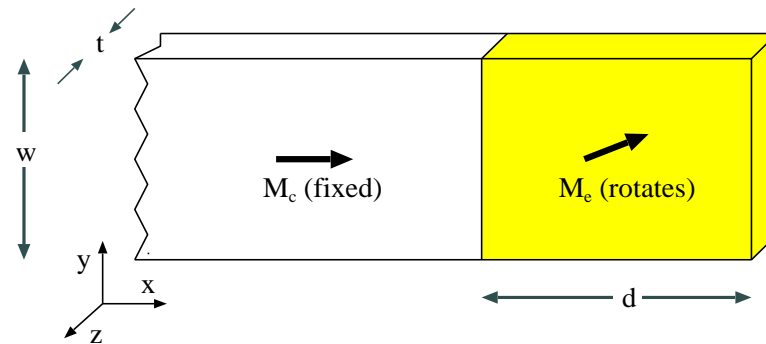


# Individual Spin Dynamics



- (a) Center spin
- (b) Intermediate spin
- (c) End spin

# Simple Edge Spin Model



$$H_{\text{exch}} = (2A/(\mu_0 M_s d^2), 0, 0)$$

$$H_{\text{demag}} = (-N_x M_x/2, -N_y M_y, -N_z M_z)$$

Then

$$\text{Freq} = \frac{\gamma M_s}{2\pi} \sqrt{\left( \frac{2A}{\mu_0 M_s^2 d^2} - \frac{N_x}{2} + N_y \right) \left( \frac{2A}{\mu_0 M_s^2 d^2} - \frac{N_x}{2} + N_z \right)}$$

where

$$d \approx 4 l_{\text{ex}} = 4 \sqrt{2A/\mu_0 M_s^2}$$

## Simple Model Results

Part width	Frequency (GHz)
20	9.6
30	7.8
40	6.5
50	5.4
60	4.5

# Conclusions

- $\mu$ MAG Problem 4 exhibits 2 reversal modes:
  - $H_{\text{app},170^\circ}$ : End domains propagate towards center.
  - $H_{\text{app},190^\circ}$ : Counter-rotating domains,  $360^\circ$  walls, complex dynamics.
- Near critical field  $H_c$ :
  - $\chi \propto (H_c - H)^{-1/2}$  for small range of  $H$ .
  - $f \propto (H_c - H)^{1/4}$  for wider range of  $H$ .
- Remanence dynamics:
  - Distinct central core and end domain behavior.
  - End domain frequency especially dependent on stripe width.



## References

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- $\mu$ MAG:

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