

***Relaxation of shear-induced complex fluid states
using time resolved SANS***

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&

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Two examples

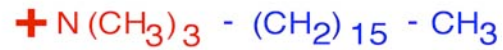
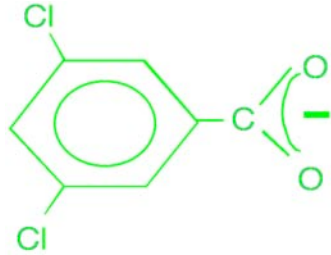
(1) Relaxation of a Poiseuille shear-induced surface state

*Melting of a crystalline state
in a threadlike micellar system*

(2) Relaxation of a Couette shear-induced bulk state

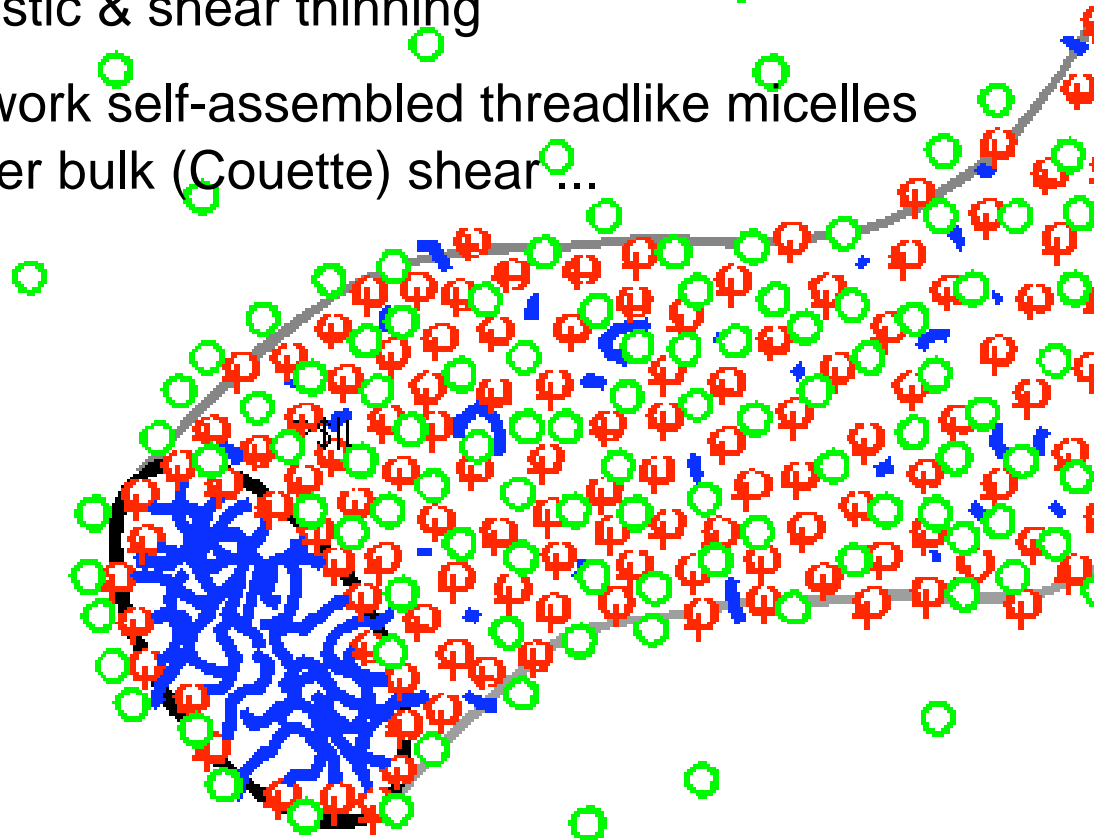
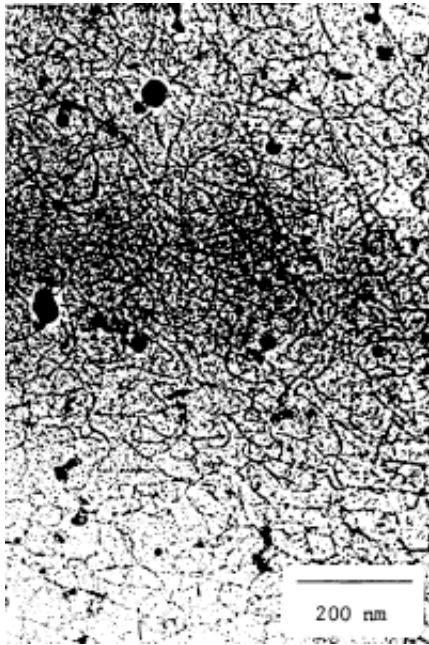
*The kinetics and energetics
of
passage formation in membrane phases*

Threadlike micelles (CTA35CIBz)



Macroscopically: Even quite dilute solutions are thick, but viscoelastic & shear thinning

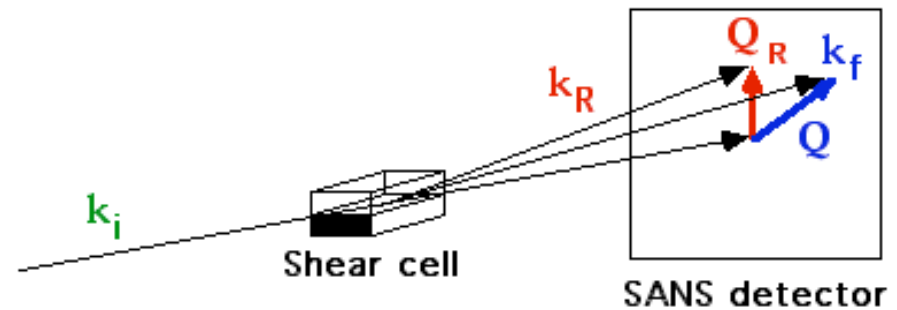
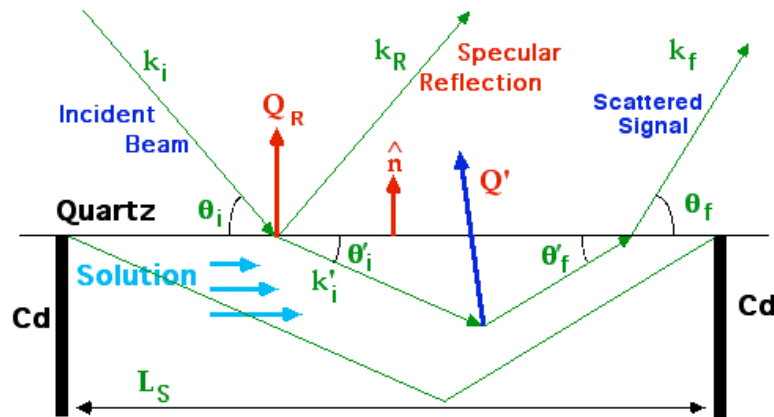
Microscopically: Entangled network self-assembled threadlike micelles which SANS shows line up under bulk (Couette) shear ...



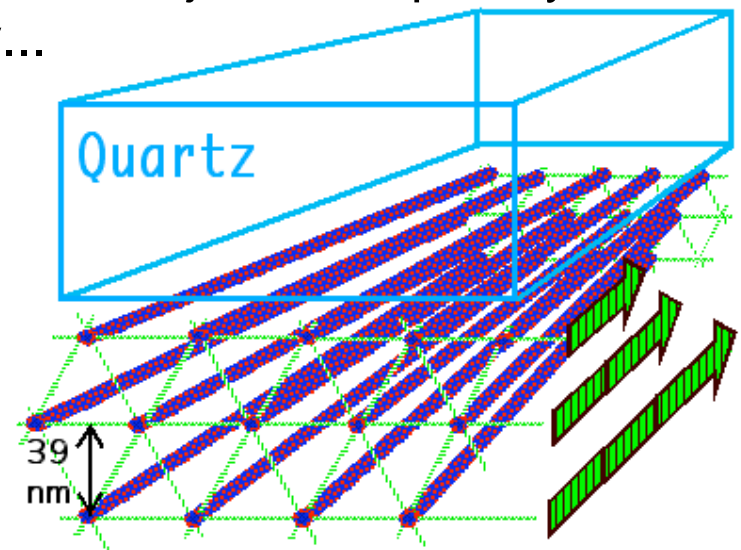
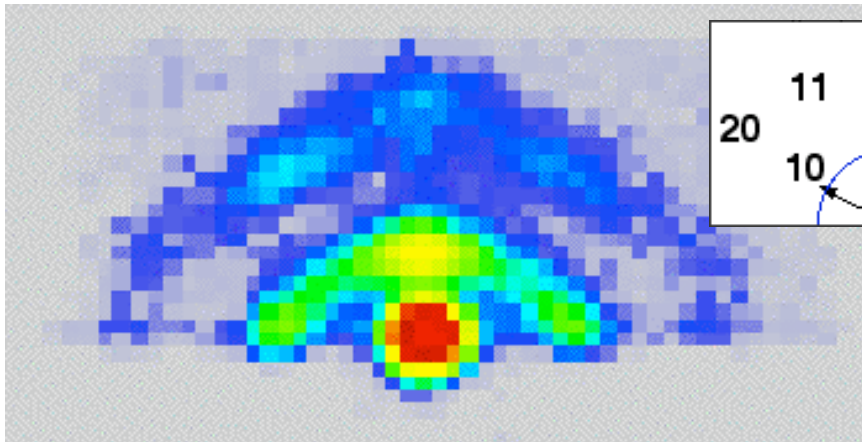
30Å diameter, 1000's Å long ...

CryoTEM micrograph of pure 3,5-dichlorobenzoate micelles
L.J. Magid, J.C. Gee & Y. Talmon, *Langmuir* **6**, 1609 (1990)

Poiseuille shear response of mixed counterion 20mM 70% CTA35CIBz & 30% CTABr threadlike micelles



In Poiseuille shear past a surface the micelles don't just line up, they form a strongly oriented crystalline hexagonal array...



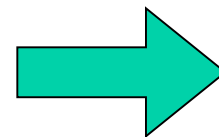
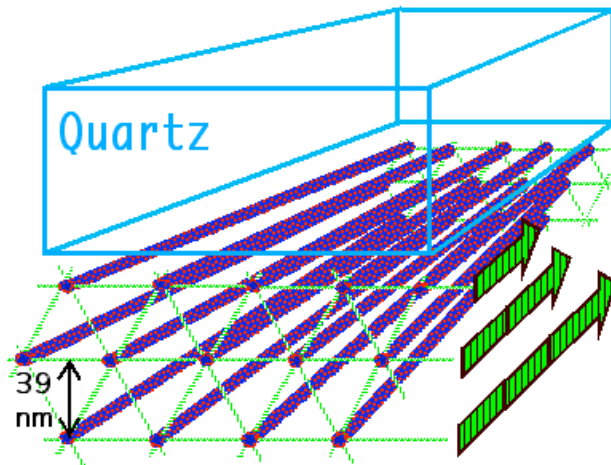
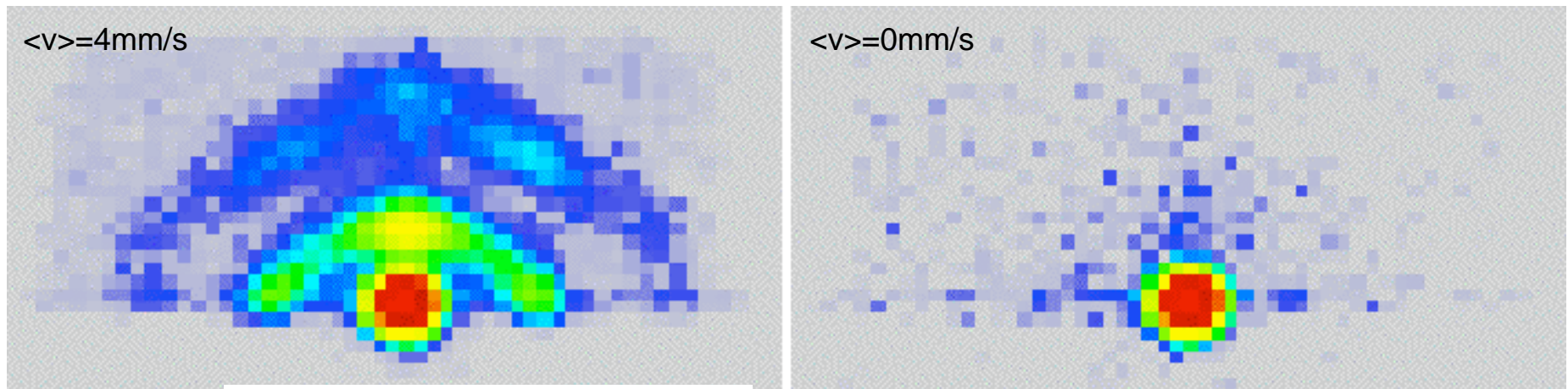
Grazing incidence “Near Surface” SANS data ($< \sim 100$ micron) from surface

W.A. Hamilton, P.D. Butler, S.M. Baker, G.S. Smith, J.B. Hayter, L.J. Magid and R. Pynn, *Physical Review Letters* **72**, 2219 (1994)

W.A. Hamilton, P. D. Butler, John B. Hayter, L. J. Magid and P. J. Kreke, *Physica B* **221**, 309 (1996)

*A question of relaxation when shearing flow stopped:
Couette measurements => slow bulk micellar reentanglement
~30min*

*But surface Xtal state apparently disappeared “immediately”
~seconds*

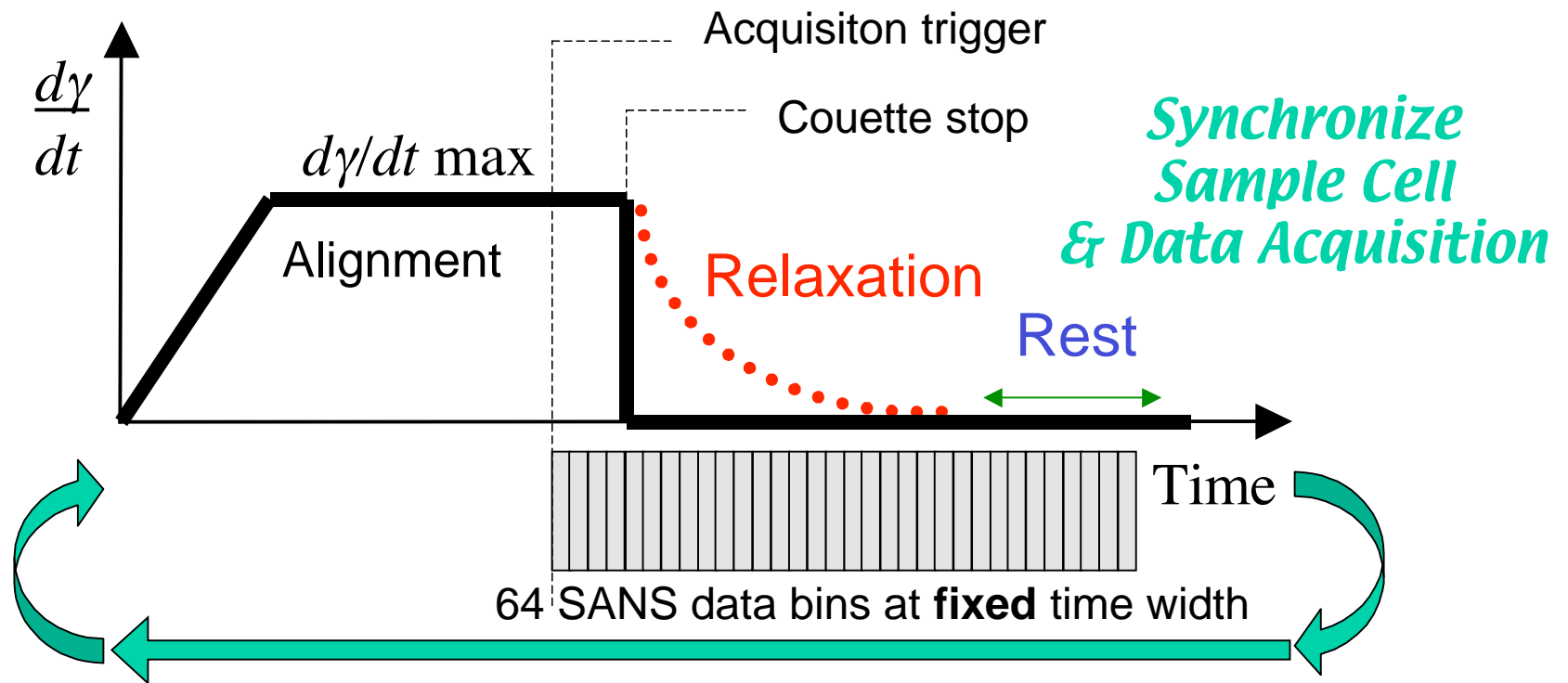


Where does it go ...
how does it relax ?

When things are (a little) too fast for SANS - “t-SANS”

Cycled multiplexed synchronized data collection

Very fast (for SANS) \Rightarrow “time sliced” cycled SANS (NIST-NG7)



Cycle alignment-relaxation-rest to build up statistics in SANS time bins

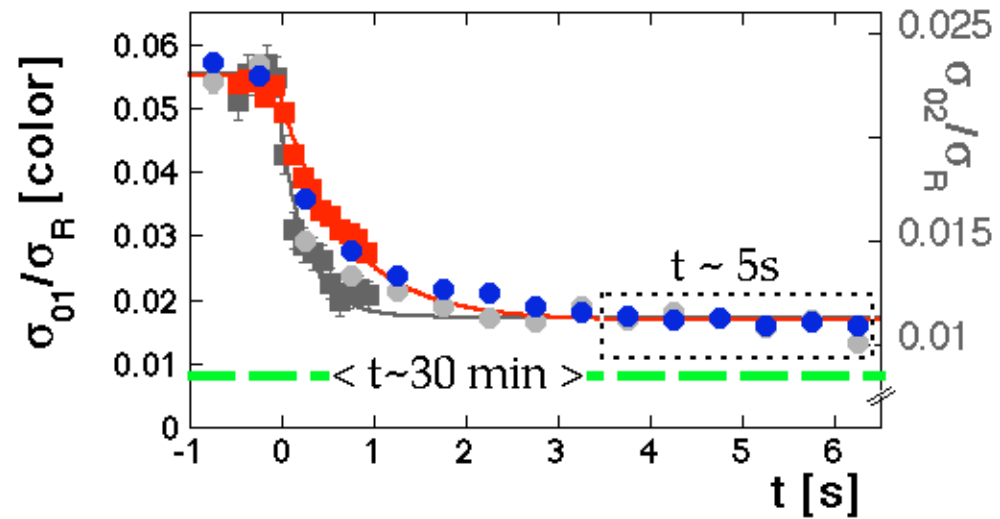
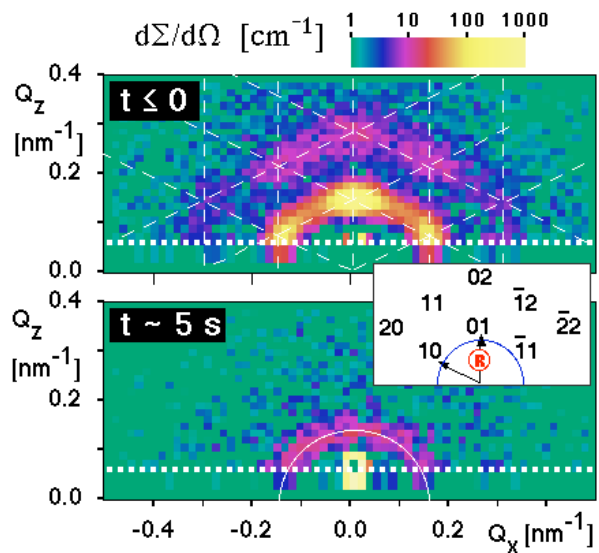
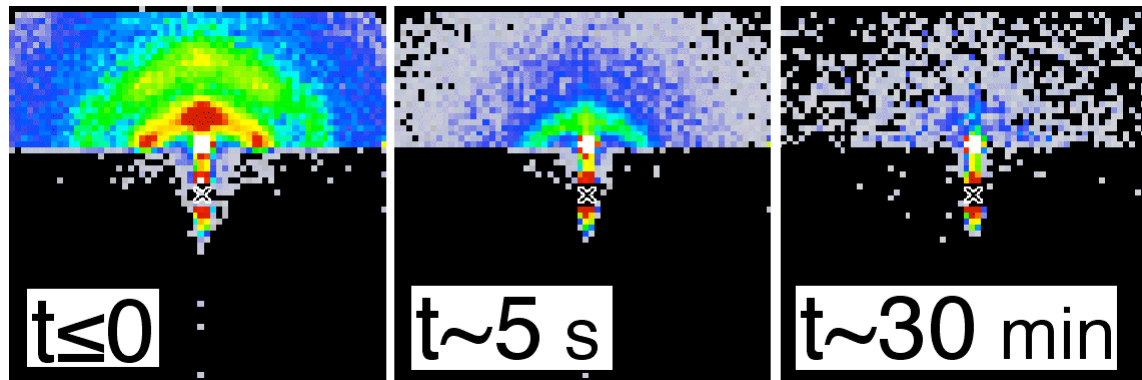
“Relaxation of a shear-induced lamellar phase measured with time resolved small angle neutron scattering”,

L. Porcar, W.A. Hamilton, P.D. Butler and G.G. Warr, *Physica B* **350**, e963 (2004)

Once upon a time: “Fast Relaxation of a Hexagonal Poiseuille Shear-induced Near-Surface Phase in a Threadlike Micellar Solution”,
W.A. Hamilton, P.D. Butler, L.J. Magid, Z. Han and T.M. Slawcki, *Physical Review E (Rapid Communications)* **60**, 1146 (1999)

What we saw in threadlike micelle relaxation

0.1s and 0.5s multiplexed data reveals a two stage process



Xtal phase 01 peak fast decay time 0.7 ± 0.2 s

Nature of the relaxation

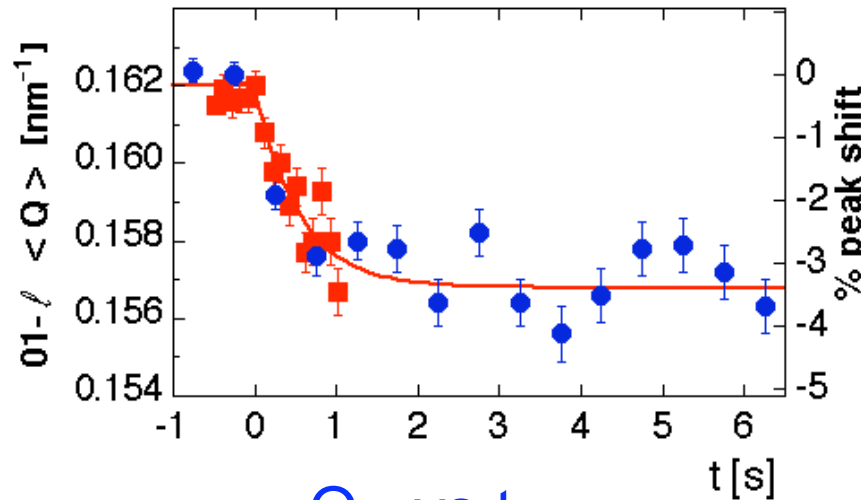
nearest neighbor $\langle Q \rangle$
2D liquid vs hex lattice

$$\begin{aligned}
 I[Q_R] &= \int_{R=0}^{\infty} \int_{\phi=0}^{2\pi} (\beta[R] - \beta_o) e^{iQ_R R \cos\phi} R dR d\phi \\
 &= \int_{R=0}^{\infty} R (\beta[R] - \beta_o) \int_{\phi=0}^{2\pi} e^{iQ_R R \cos\phi} d\phi dR \\
 &= \int_{R=0}^{\infty} 2\pi R (\beta[R] - \beta_o) J_0[Q_R R] dR
 \end{aligned}$$

$$\begin{aligned}
 \frac{\partial}{\partial Q_R} I = 0 &\Rightarrow \frac{\partial}{\partial Q_R} J_0[Q_R R] = 0 \\
 &\Rightarrow R J_1[Q_R R] = 0 \\
 &\Rightarrow Q_R R = 0, 3.832, \mathbf{7.0156}, 10.173, \dots
 \end{aligned}$$

$$Q_{01} = \frac{2\pi}{\sqrt{3} R_m / 2} = \frac{7.2552}{R_m}$$

$$Q_\ell = \frac{7.0156}{R_m} = (1 - 3.30\%) \frac{7.2552}{R_m}$$



$\langle Q \rangle$ vs t
in 01 peak region

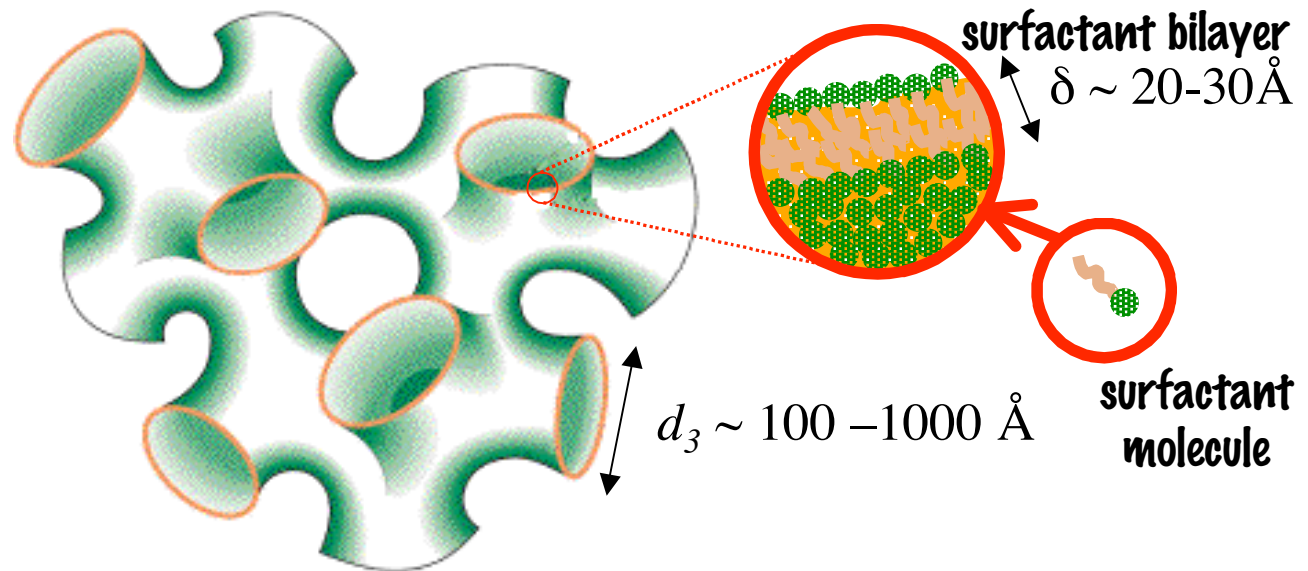
-3.4±0.3% shift between first order diffraction spot Q and that of the arc on same time scale as decay. This is consistent with melting of the hexagonal phase (predicted 2D xtal shift -3.3%) rather than simple loss of crystalline alignment in which case the peak would not shift.

The crystal is melting before micellar reentanglement

W.A. Hamilton, P.D. Butler, L.J. Magid, Z. Han & T.M. Slaweki, Phys. Rev. E (Rapid Comm.) **60**, R1146 (1999)

L_3 “sponge” phases

Self assembled
isotropic
membrane
phases



*A convoluted solution spanning labyrinth of membrane passages
(topologically “handles”)*

Scaling of dilution $d_3 \phi \sim \text{const}$, volume fraction $\phi \sim 40\% \dots < 1\%$

Typically very fluid

*Not much of a response to applied shear \sim Newtonian
rapid realignment of passages relieves stress*

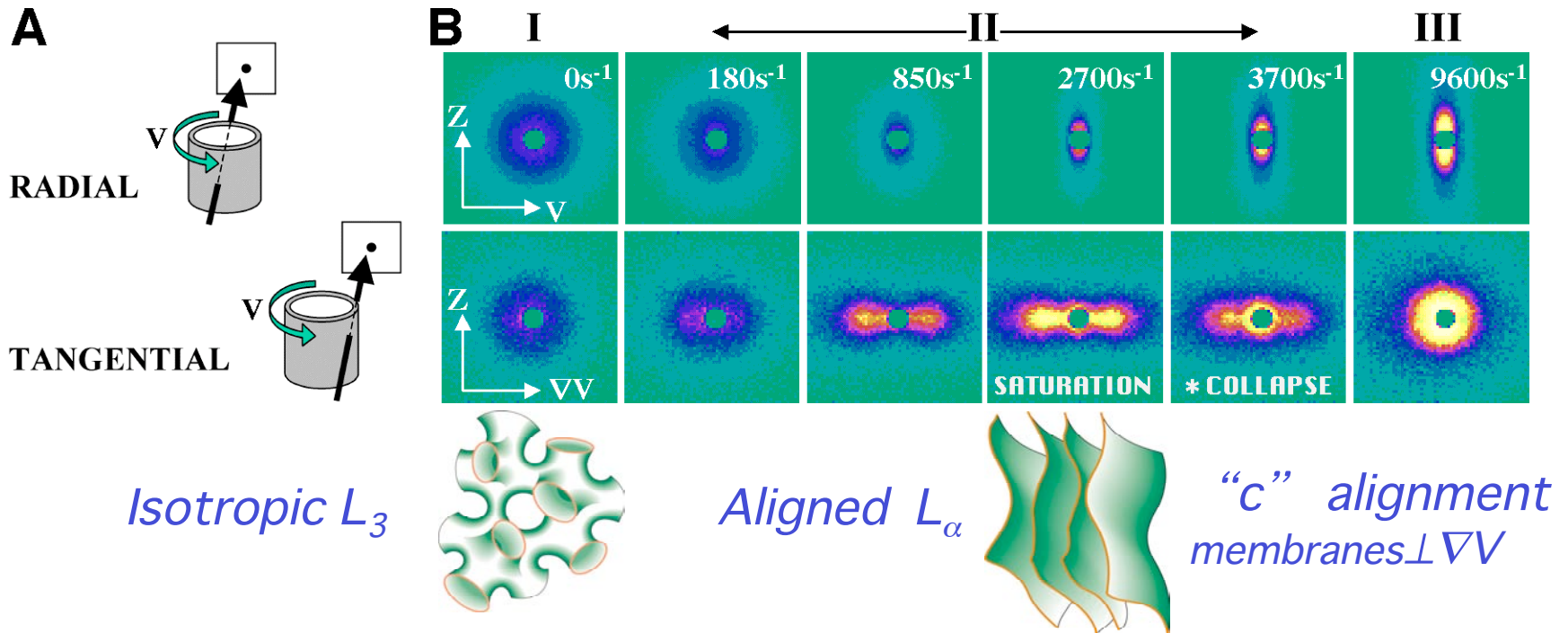
but:

Stir \Rightarrow Transient birefringence for dilute samples $d_3 > 500 \text{ \AA}$

Structural response of the “sweetened” sponges

Small Angle Neutron Scattering from $\phi=5\text{vol\% CPCI-hexanol in } 40\text{vol\% dextrose-brine } (\eta_s=16.3\text{cP})$

Equilibrium L_3 to Couette Shear-induced L_α transformation



Correlation peak $Q_3 = 0.008 \text{ \AA}^{-1}$

passage size $d_3 = 780 \text{ \AA}$

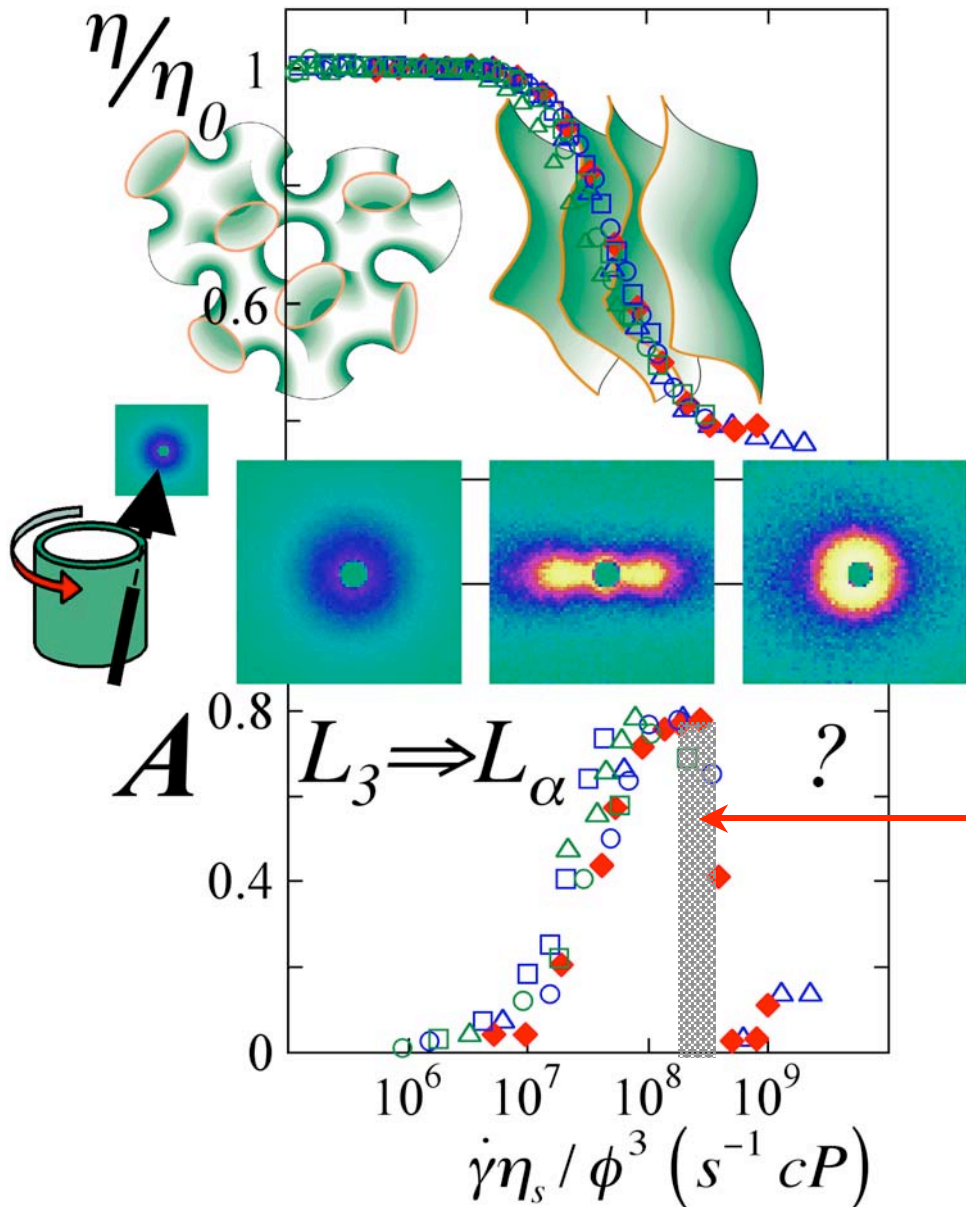
$d_\alpha / d_3 \sim 80\%$ - same ratio as for equilibrium L_α & L_3 at same ϕ

Bragg peak $Q_\alpha = 0.010 \text{ \AA}^{-1}$

separation $d_\alpha = 630 \text{ \AA}$

1st shear $L_3 \rightarrow L_\alpha$ J. Yamamoto and H. Tanaka, PRL 77, 4390(1996) $C_{12}E_5$ $\phi < 2\%$, our system $\phi < 3-7\%$

Well defined scaling of shear-induced $L_3 \rightarrow L_\alpha$



Tunable
shear-induced
 L_3 to L_α
transformation

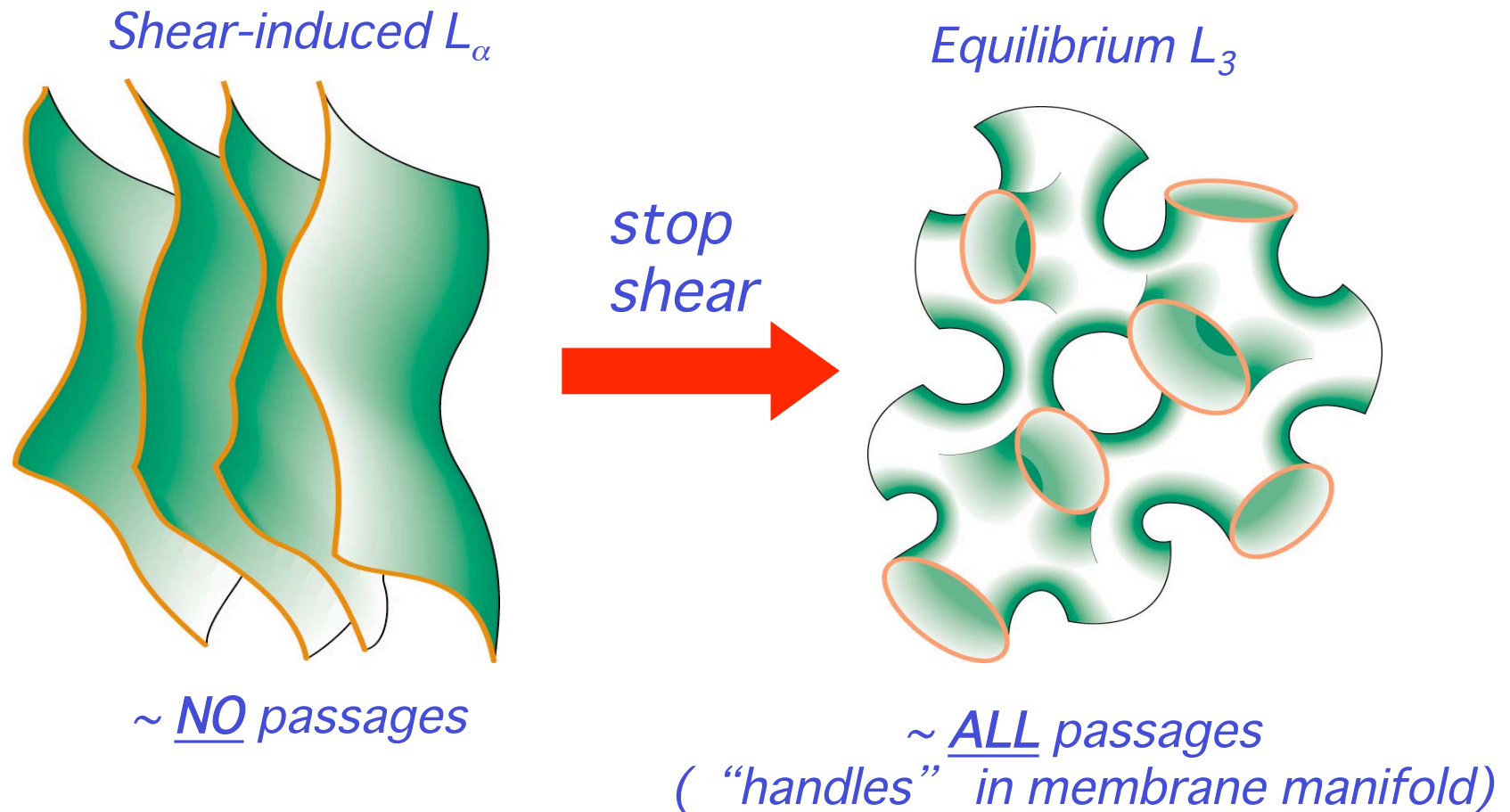
A well characterized system

L. Porcar, W.A. Hamilton, P.D. Butler and G.G. Warr,
Physical Review Letters **89**,168301 (2002)
Langmuir **19**, 10779 (2003)

L_α anisotropy saturates for
 $\dot{\gamma}\eta_s/\phi^3 \sim 2 - 4 \times 10^8 \text{ cP} / \text{s}$
Easily accessible for high η_s

So what can we do with it?

Relaxation is now a topological transformation

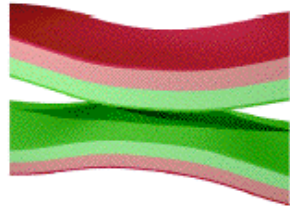


*Offers a strong clear signal
of membrane solution passage creation. So ... ?*

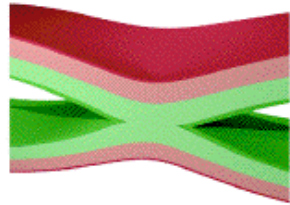
*Solution passage/pore formation ...
an important chemical and (especially) biological process*

The “stalk” hypothesis of membrane fusion

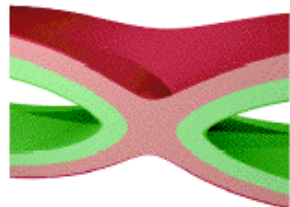
MD simulations



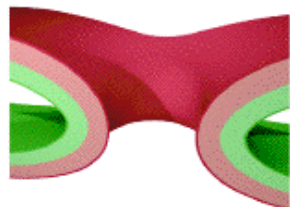
(a) Approach to small distances



(b) Merger of proximal monolayers
“Stalk” formation

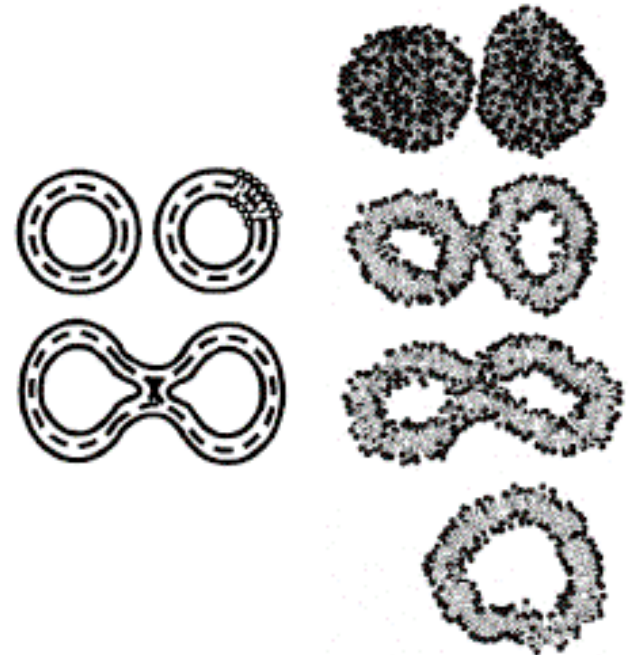


(c) Merger of distal monolayers
Stalk to TMC evolution



(d) Expansion to pore formation

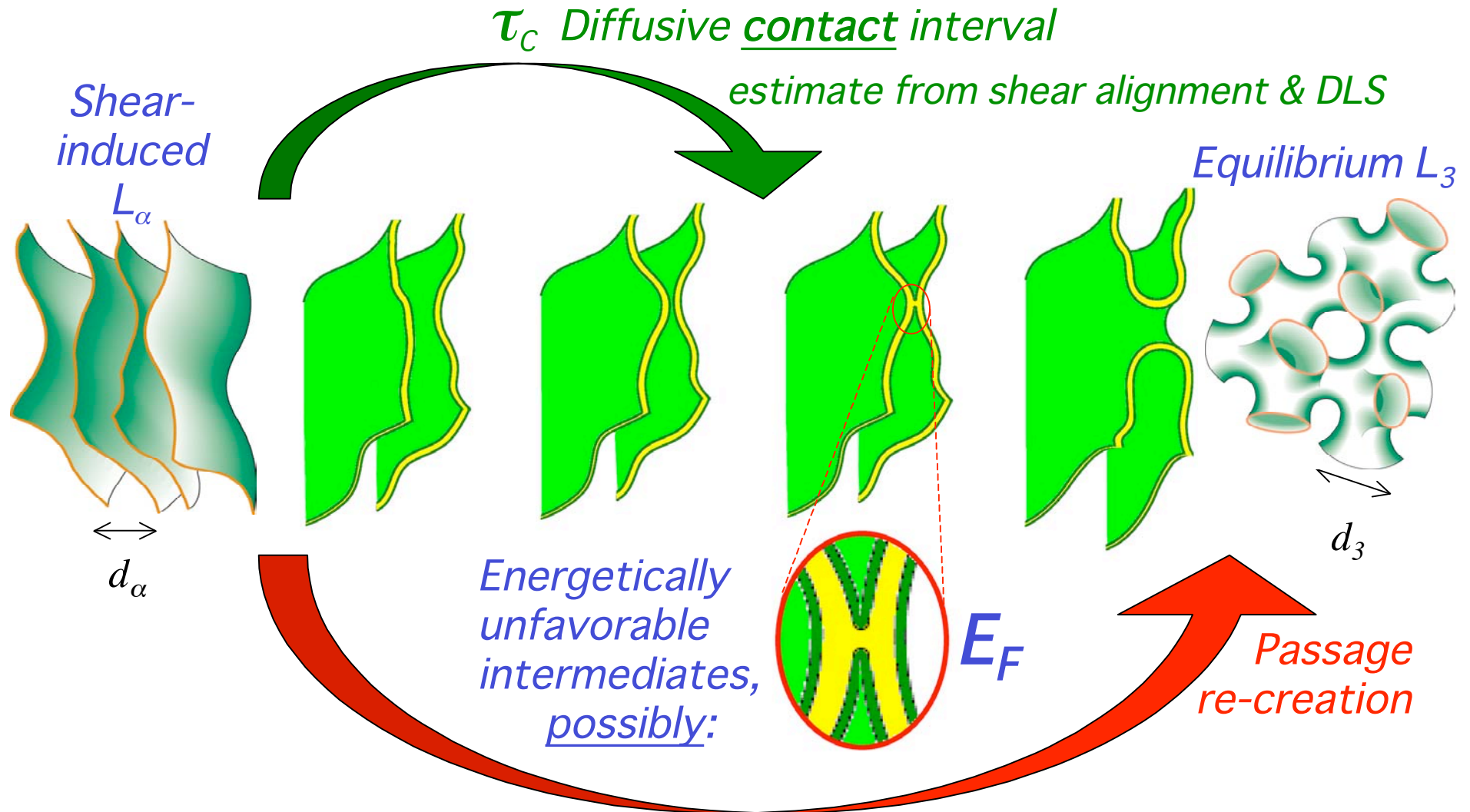
from R. Jahn and H. Grubmüller,
Current Opinion in Cell Biology **14**, 488 (2002)



from H. Noguchi and M. Takasu,
J. Chem. Phys. **115**, 9547 (2001)

*Stalk, TMC, or some other intermediate structure
⇒ energy barrier to solution passage formation
(NB Our results can benchmark biophysical theories)*

Energetics of passage formation

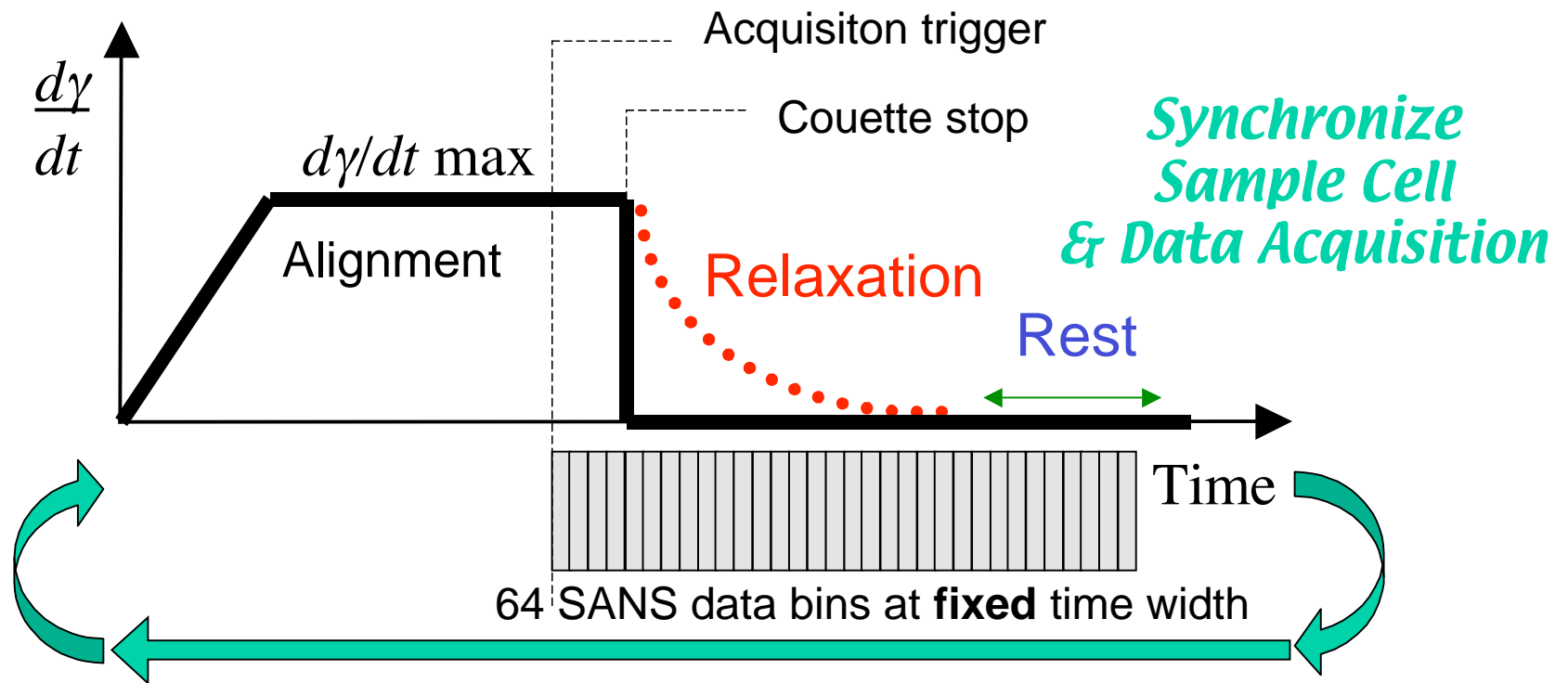


Topological relaxation time $\tau_R = \tau_C \exp[-E_F/k_B T]$

When things are (a little) too fast for SANS - “t-SANS”

Even at highest viscosity full topological relaxation $\tau_R \sim$ seconds

Very fast (for SANS) \Rightarrow “time sliced” cycled SANS (NIST-NG7)



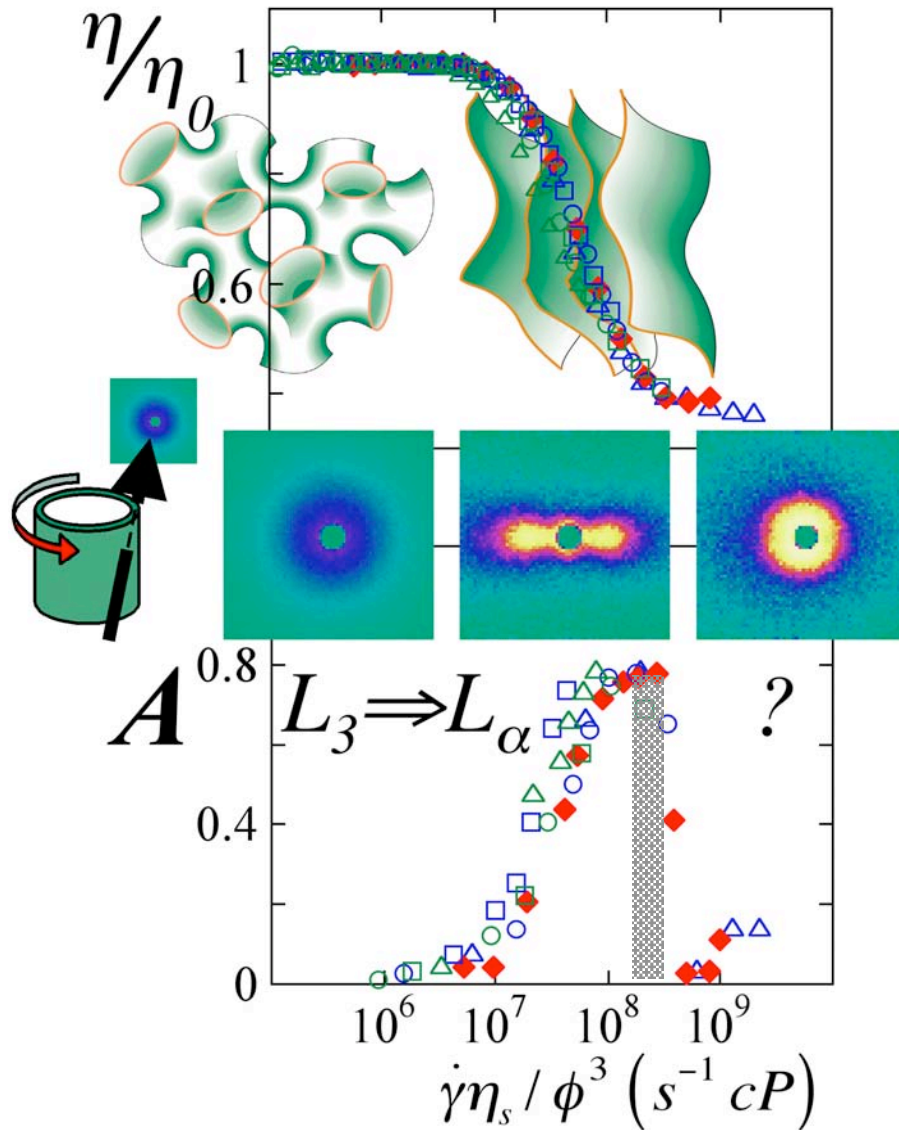
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Can estimate τ_c from earlier SANS vs shear



We already know:
 L_α signal saturates for

$$\dot{\gamma}\eta_s/\phi^3 \sim 2 - 4 \times 10^8 \text{ cP} / s$$

(shaded)

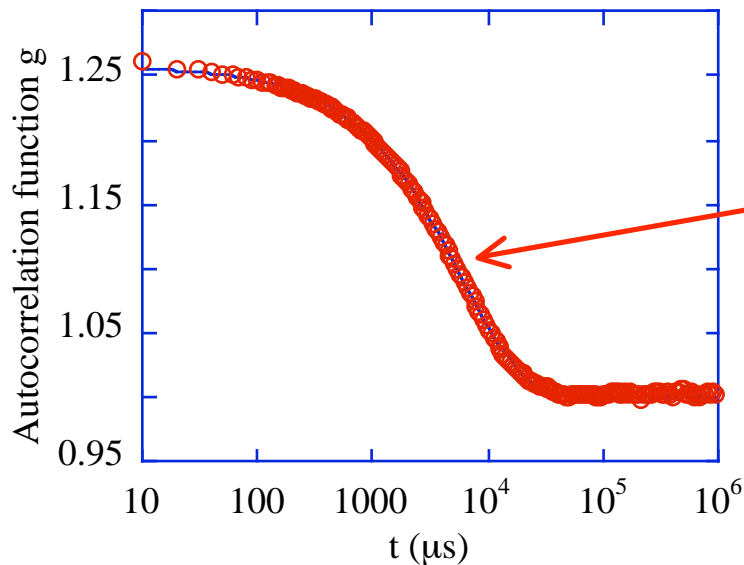
Applied shear rate (s^{-1}) represents
 1/time which totally frustrates
 (re)formation of disrupted
 membrane passage ...
 because they never meet

So can expect

$$\tau_c \sim \frac{1}{\dot{\gamma}_{Saturation}} \sim \frac{\eta_s/\phi^3}{2 - 4 \times 10^8 \text{ cP} / s}$$

Or ... we can use Dynamic Light Scattering (DLS) to measure membrane diffusion rates
 $\tau_c \Rightarrow$ time to bring membranes separated by an average separation d_α into contact

Determination of contact time τ_c (#2) - DLS



Homodyne autocorrelation function:

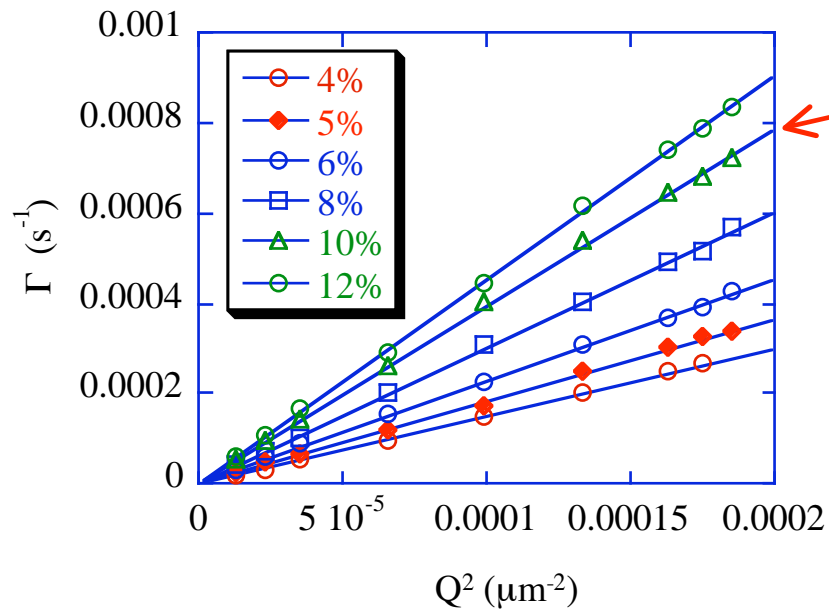
$$1 + A \exp[-2(\Gamma t)^\beta]$$

Stretched exponential $\beta=0.7-0.8$

$$\Gamma[Q] = DQ^2$$

\Rightarrow simple cooperative diffusion process

$$D \propto \phi$$



Shear-induced L_α separation $d_\alpha \propto 1/\phi$

$$\Rightarrow \tau_c \approx 1/\Gamma[Q_\alpha] = 1/DQ_\alpha^2$$

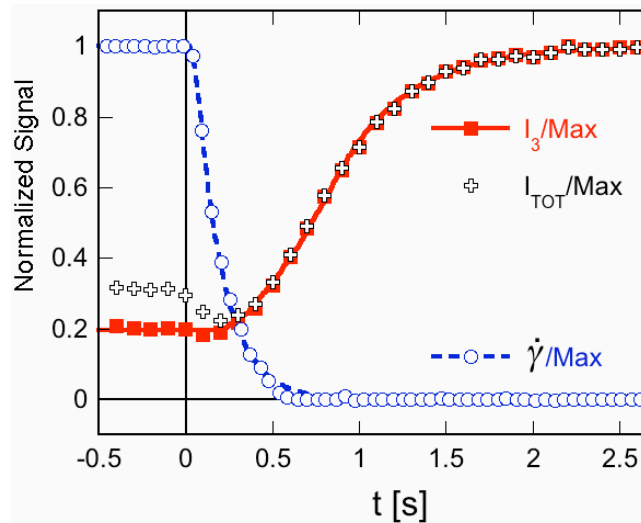
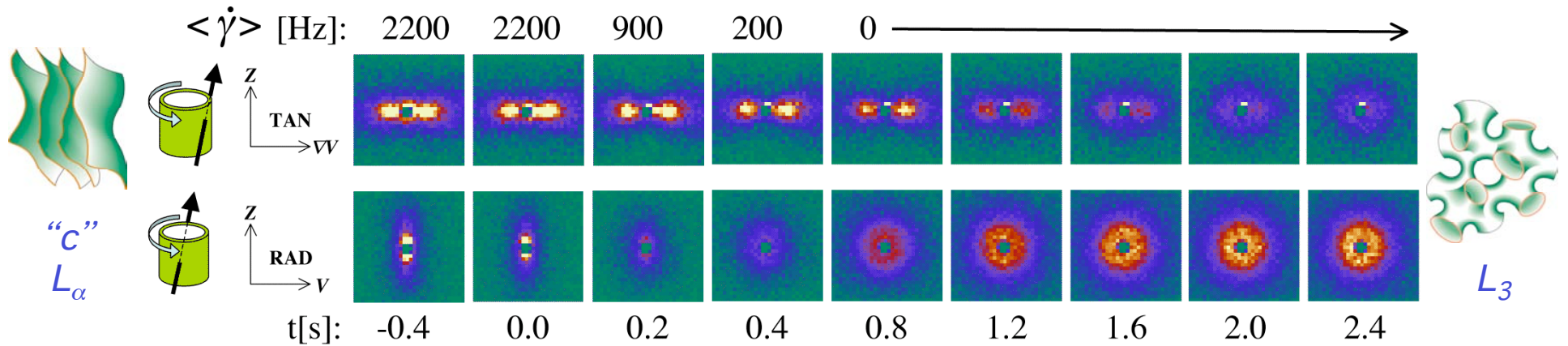
where $Q_\alpha = 2\pi/d_\alpha \propto \phi$

$$\Rightarrow \tau_c \propto 1/\phi^3$$

agrees with Shear-induced L_α plateau estimate

Determination of topological relaxation time τ_R - *t*-SANS

t-SANS Shear-induced L_α to equilibrium L_3 relaxation
 $\phi=5\text{vol\%}$ CPCI-hexanol in 40vol% dextrose-brine ($\eta_s=16.3\text{cP}$)



Shear aligned at

$$\dot{\gamma} \eta_s / \phi^3 \sim 3 \times 10^8 \text{ cP} / \text{s}$$

\sim center L_α signal plateau

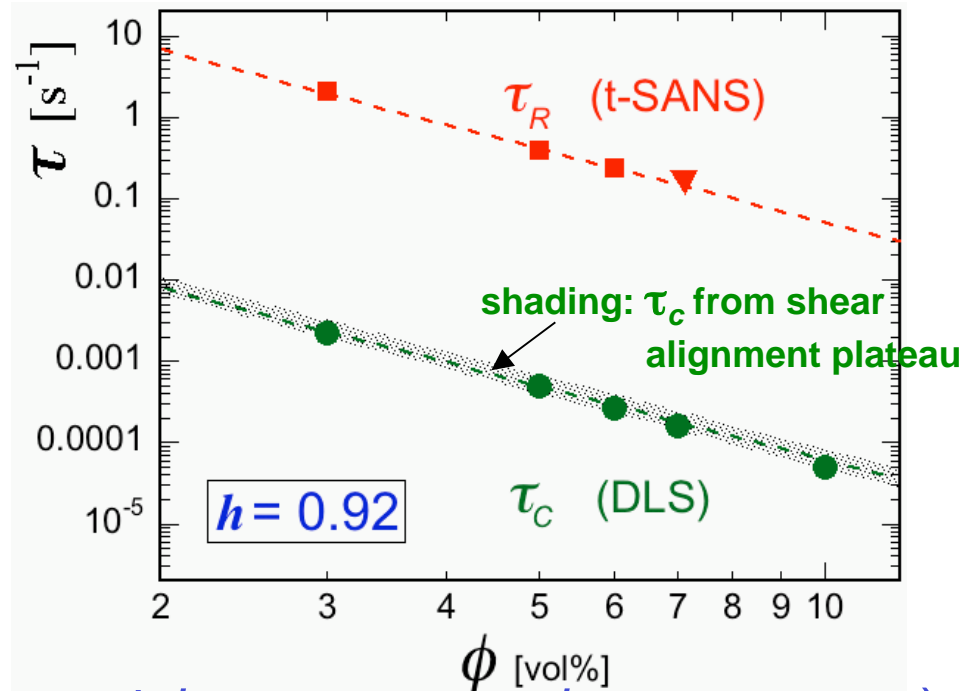
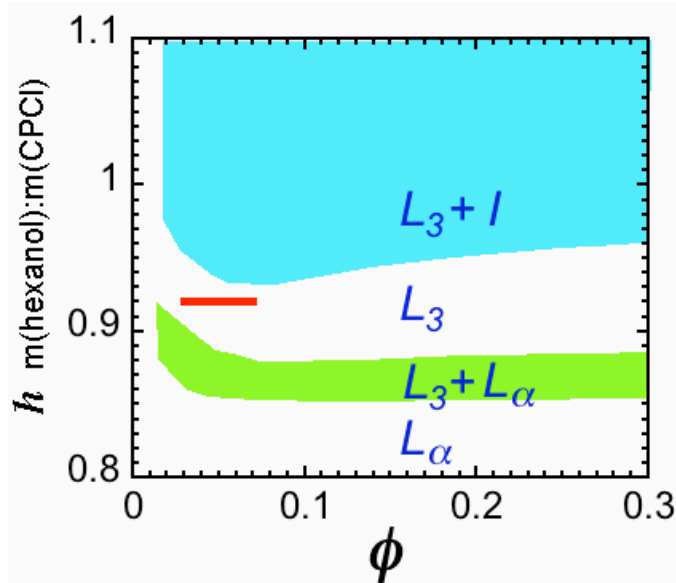
When Couette cell is stopped
 L_3 signal (passages) re-established

$$\tau_R = 0.40 \pm 0.08 \text{ s}$$

τ_R and τ_C versus membrane volume fraction ϕ

Constant membrane composition, i.e. properties

Dilution series: $d_3\phi \approx \text{const}$, $d_\alpha\phi \approx \text{const}$



τ_R and τ_C scale as ϕ^{-3} (as you might expect per shear response)

Constant Arrhenius relationship (despite $>2X$ scaling)

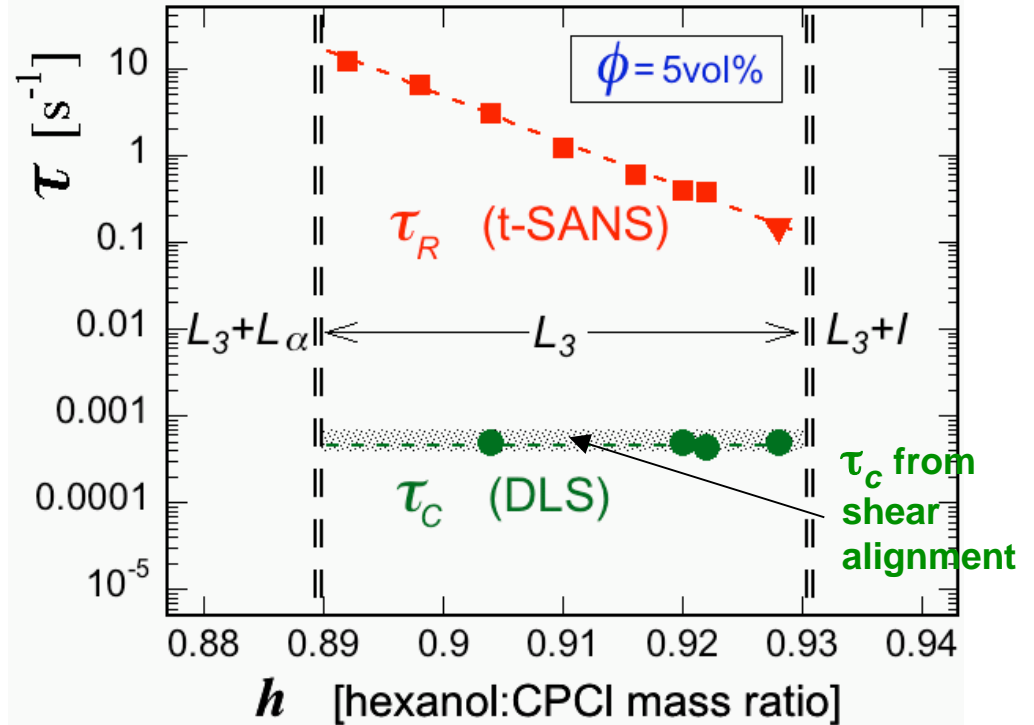
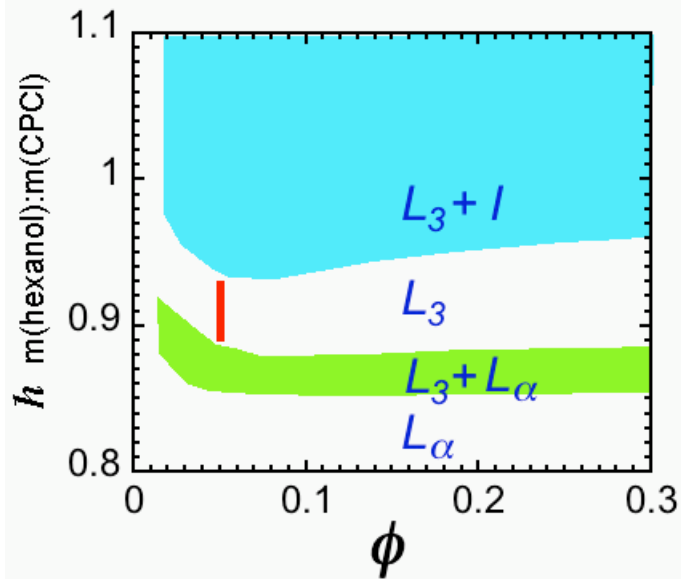
$$\tau_R = \tau_C \exp[-E_F/k_B T] \Rightarrow E_F = 6.7k_B T \quad (170 \text{ meV})$$

Topological relaxation of a shear-induced lamellar phase to sponge equilibrium and the energetics of membrane fusion",

L. Porcar, W.A. Hamilton, P.D. Butler and G.G. Warr, *Physical Review Letters* **93**, 198301(2004)

τ_R and τ_C versus membrane composition h

Change membrane composition, i.e. properties across L_3 phase region for constant ϕ



Increasing hexanol to CPCI mass ratio h

⇒ Increasing Gaussian curvature modulus of membranes

⇒ Decreasing energy cost of passages (and stalk structures)

4% increase in h $E_F = 10.3k_B T$ (260 meV) down to $5.8k_B T$ (150 meV)

Summary/Conclusions/Future

Nice information from relatively simple multiplexed time-resolved SANS

(1) Threadlike micellar system

Second order peaks decay 4 times faster 0.1-0.2 s than first order peaks

\Rightarrow Debye-Waller factor $\sim \exp[-Q^2 \langle \Delta r^2(t) \rangle / 2]$

2D translational diffusion constant: $\langle \Delta r^2(t) \rangle = \langle \Delta r^2(0) \rangle + Dt$

$\tau = 0.7 \pm 0.2$ s for $Q_0 = 0.16 \text{ nm}^{-1}$ $D = 2 / (Q^2 \tau) = 110 \pm 30 \text{ nm}^2/\text{s}$

Fast relaxation of the hexagonal threadlike lattice is melting

(2) “Sweetened” Cetylpyridinium-Hexanol/dextrose-brine L_3 sponges

Viscosity tuning \Rightarrow accessible shear-induced L_3 to L_α transition

Activation energy for membrane fusion (handle creation)

$$\tau_R = \tau_C \exp[-E_F/k_B T] \Rightarrow E_F \sim 5 - 10 k_B T$$

E_F constant wrt ϕ (\Rightarrow constant barrier state - stalk/TMC ?)

E_F linear decrease wrt h across L_3 phase region \Leftrightarrow Gaussian curvature modulus

**Relaxation measurements to “0.01s” easy enough to do in
data acquisition hardware/software ...**

we will do this on new HFIR SANS in 2006 ...

with a little more synchronization possible for SNS SANS

So start thinking about measurements ...

References

Threadlike micellar shear response and relaxation

"Fast Relaxation of a Hexagonal Poiseuille Shear-induced Near-Surface Phase in a Threadlike Micellar Solution",
W.A. Hamilton, P.D. Butler, L.J. Magid, Z. Han and T.M. Slaweki,
Physical Review E (Rapid Communications) **60**, 1146 (1999)

"Effect of a solid/liquid interface on bulk solution structures under flow",
P.D. Butler, W.A. Hamilton and L.J. Magid et al.,
Physica B **241**, 1074 (1997)

"Over the Horizon" SANS: Measurements on Near-Surface Poiseuille Shear-Induced Ordering of Dilute Solutions of Threadlike Micelles",
W.A. Hamilton, P. D. Butler, John B. Hayter, L. J. Magid and P. J. Kreke,
Proceedings of the Fourth International Conference on Surface X-Ray and Neutron Scattering , *Physica B* **221**, 309 (1996).

"Kinetics of Decay and Alignment in a Highly Entangled Transient Threadlike Micellar Network Studied by Small Angle Neutron Scattering",
P.D. Butler, L.J. Magid, W.A. Hamilton, J.B. Hayter, B. Hammouda and P.J. Kreke, *Journal of Physical Chemistry* **100**, 442-445 (1996)*.

"Shear Induced Hexagonal Ordering Observed in an Ionic Viscoelastic Fluid in Flow past a Surface",
W.A. Hamilton, P.D. Butler, S.M. Baker, G.S. Smith, J.B. Hayter, L.J. Magid and R. Pynn,
Physical Review Letters **72**, 2219 (1994).

"Shear cell for the study of liquid_solid interfaces by neutron scattering",
S.M. Baker, G.S. Smith, P.D. Butler, J.B. Hayter, W.A. Hamilton, R. Pynn and L.J. Magid,
Review of Scientific Instruments **65**, 412 (1994).

Sponge phase shear response and relaxation

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L. Porcar, W.A. Hamilton, P.D. Butler and G.G. Warr,
Physical Review Letters **93**, 198301(2004)*
[and *Virtual Journal of Biological Physics Research* **8(10)** (2004) <http://www.vjbio.org>].

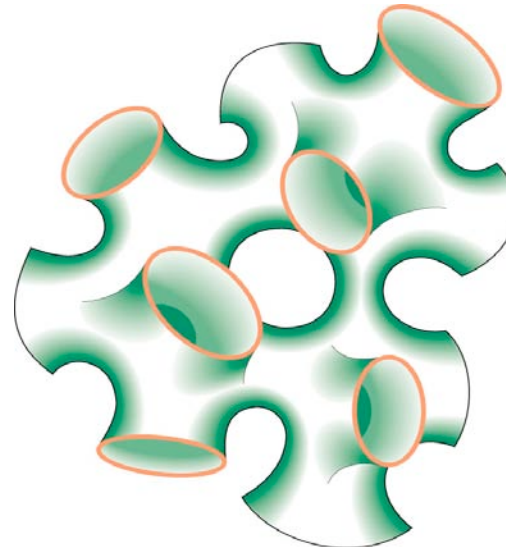
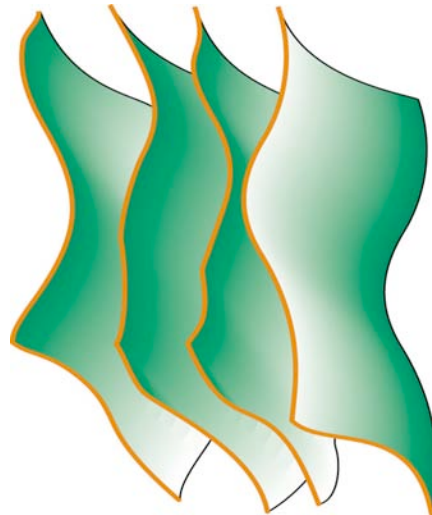
"Relaxation of a shear-induced lamellar phase measured with time resolved small angle neutron scattering",
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Physica B **350**, e963 (2004)*

"Scaling of Structural and Rheological Response of L₃ Sponge phases
in the "Sweetened" Cetylpyridinium-Hexanol/Dextrose-Brine System",
L. Porcar, W.A. Hamilton, P.D. Butler and G.G. Warr,
Langmuir **19**, 10779-10794 (2003)

"Scaling behavior of shear-induced sponge to lamellar transformations",
L. Porcar, W.A. Hamilton, P.D. Butler and G.G. Warr, *Physical Review Letters* **89**,168301 (2002)
[and *Virtual Journal of Biological Physics Research* **4** (2002) <http://www.vjbio.org>].

Bonus: Membrane phase energetics

L_α
 \sim NO
 passages



L_3
ALL
 passages
 (“handles”
 in membrane
 manifold)

Helfrich membrane Hamiltonian: $dE = \kappa(1/r_1 + 1/r_2)/2 + \bar{\kappa}(1/r_1 r_2)$

$1/r_1, 1/r_2$ curvatures κ bending modulus $\bar{\kappa}$ Gaussian curvature modulus

Gauss-Bonnet Theorem: $\iint dA(1/r_1 r_2) = 4\pi(n_p - n_h)$

\Rightarrow *Gaussian Curvature Energy per passage: $-4\pi\bar{\kappa}$*

$\Rightarrow \bar{\kappa} > 0$ *Favors passage formation i.e. $\Rightarrow L_3$*