



Pressure Dependence of Fragile-to-Strong Transition and a Possible Second Critical Point in Supercooled Confined Water

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The anomalous behaviors of thermal dynamic response functions and transport properties of supercooled water imply the existence of a Liquid-to-Liquid (L-L) transition line and the associated second low-temperature critical point. However, both the L-L transition line and second critical point are predicted to lie inside the inaccessible "No man's land." (See Fig. 1)

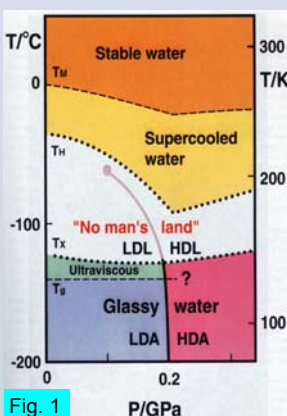


Fig. 1

We investigated, using quasi-elastic and inelastic neutron scattering, the slow single-particle dynamics of water confined in lab synthesized nanoporous silica matrices, MCM-41-S, with pore diameters ranging from 10 Å to 18 Å (see Fig. 3). Inside the pores of these matrices, the freezing process of water is strongly inhibited.

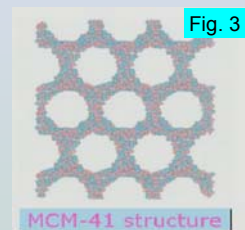


Fig. 3

We analyzed QENS spectra with a relaxing-cage model (RCM) and determined the temperature and pressure dependences of the Q-dependent translational relaxation time. (See Fig. 4)

Relaxing-Cage Model (RCM)

Experiment Measures Double Differential Cross Section $\rightarrow \frac{d^2\sigma_H}{d\Omega d\omega} = 2N \frac{\sigma_H k_L}{4\pi k_i} S_H(Q, \omega)$

Dynamic Structure Factor

Dynamics Structure Factor $S_H(Q, \omega) = pR(Q_0, \omega) + (1-p)FT \{F_H(Q, t)R(Q_0, t)\} = FT [\text{Intermediate Scattering Function } F_H(Q, t)]$

$F_H(Q, t) = F_T(Q, t) \cdot F_R(Q, t)$ For $Q < 1.1 \text{ \AA}^{-1}$, $F_R(Q, t) \approx 1$, so $F_H(Q, t) \approx F_T(Q, t)$.

$F_T(Q, t) = F_T^+(Q, t) \exp[-(t/\tau_T)^\beta]$, $\tau_T = \tau_0 (aQ)^{-\gamma}$; $\langle \tau_T \rangle = \frac{\tau_0}{\beta} \Gamma\left(\frac{1}{\beta}\right)$

$F_T^+(Q, t)$ is calculated from known proton density of states.

Fragile $\tau = \tau_0 \exp\left[\frac{DT_0}{T-T_0}\right]$ Fragile-to-Strong (F-S) transition is defined as a temperature T_L where:
e.g. VFT law
Strong $\tau = \tau_0 \exp\left[\frac{E_A/k_B}{T}\right]$ Arrhenius law

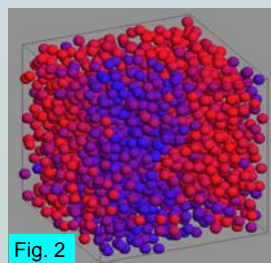


Fig. 1: Schematic illustration indicating the various phases of liquid water (color-coded). Courtesy of Dr. O. Mishima. The figure is taken from H. E. Stanley, "Mysteries of Water" Les Houches Lecture, May 1998.

Computer simulation image of the liquid-liquid phase separation in ST2 water, generated by P. Poole. (See Fig. 2)

The calculated Q-independent average relaxation time ($\langle \tau_T \rangle$) shows a fragile-to-strong (F-S) dynamic cross-over transition for pressures less than 1500 bar. Above this pressure, it is no longer possible to discern the characteristic feature of the F-S transition. (See Fig. 5~8)

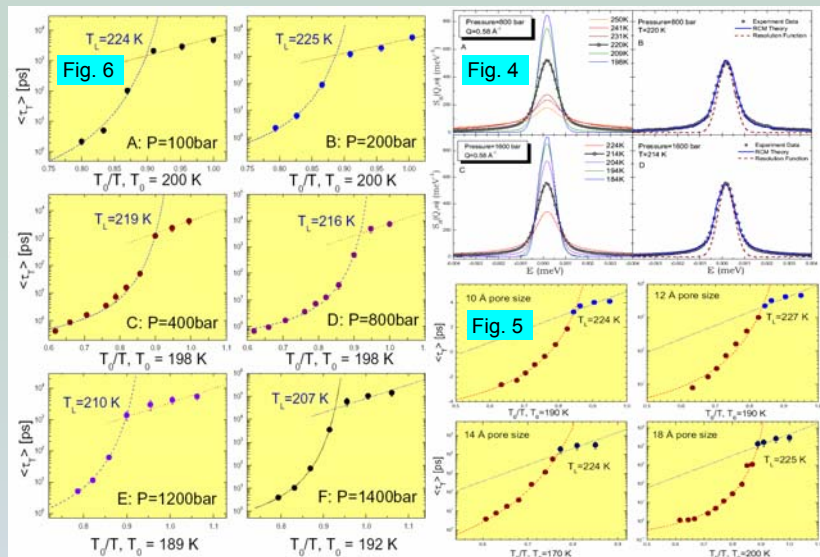


Fig. 4

Fig. 5

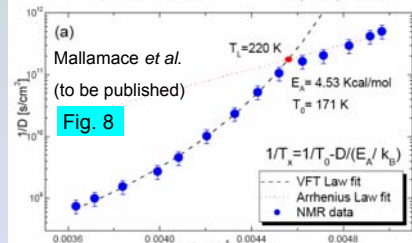


Fig. 8

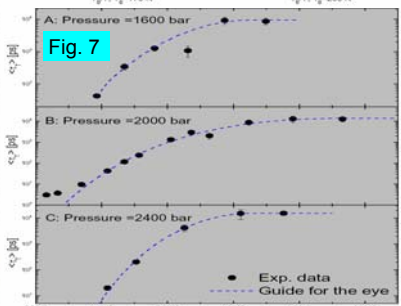


Fig. 7

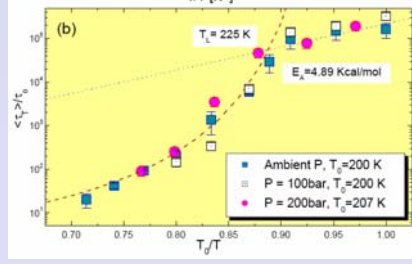


Fig. 9

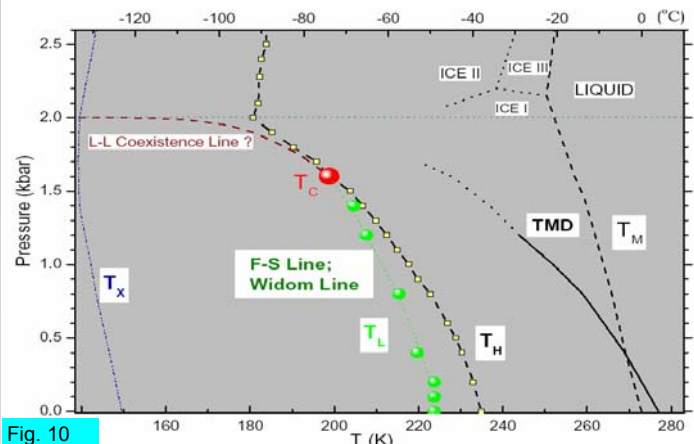


Fig. 10

New Phase Diagram (L Liu, SH Chen, et al, PRL 95, 117802 (2005)).

Discussions and Conclusion

We detected a new kind of dynamic cross-over phenomena called fragile-to-strong (F-S) transition in deeply supercooled water experimentally, for the first time.

The high-temperature (fragile) liquid corresponds to a higher density liquid (HDL), and the low-temperature (strong) liquid corresponds to a lower density liquid (LDL) (see Fig. 9).

The F-S transition line is the so-called Widom line which is the extension of the L-L coexistence line above the critical point.

We propose that the lower end point of the L-L coexistence line, or the upper end point of the Widom line, is the second low-temperature critical point ($P_C = 1500 \pm 100$ bar, and $T_C = 200 \pm 8$ K).