

# Fermi $\gamma$ -ray line in an axion mediated DM model

**Myeonghun Park**  
(CERN)

Fermilab theory seminar (Dec. 6th 2012)

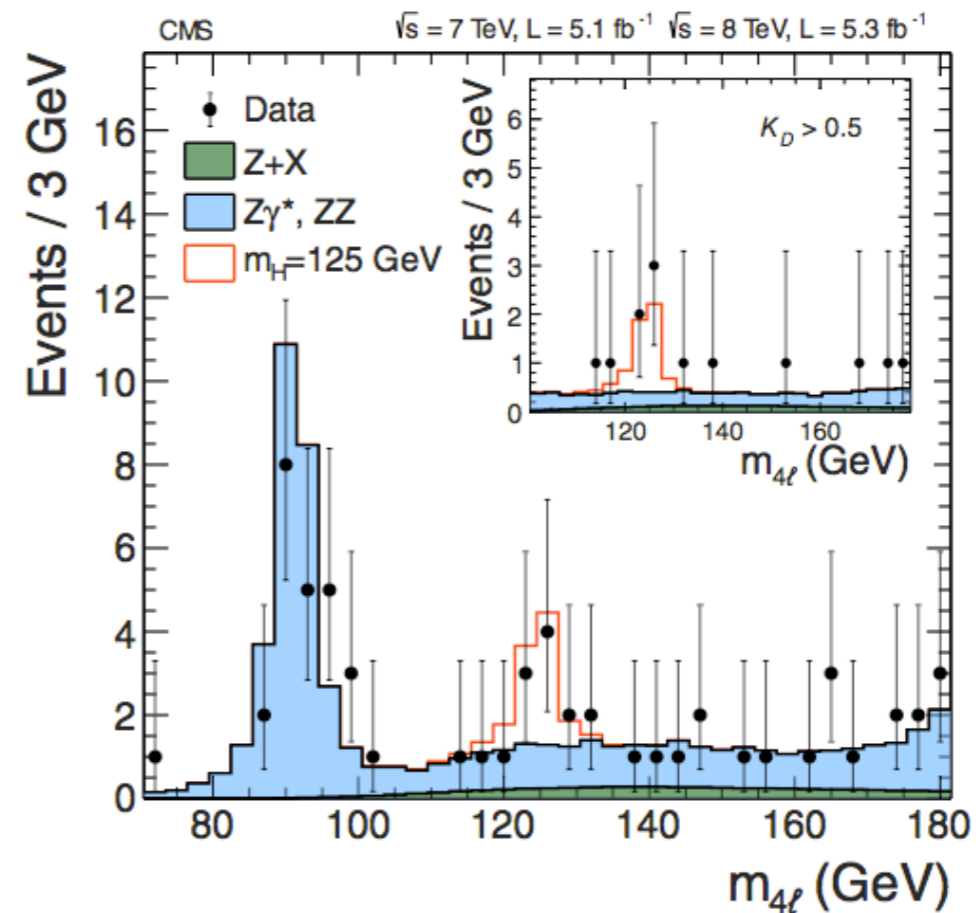
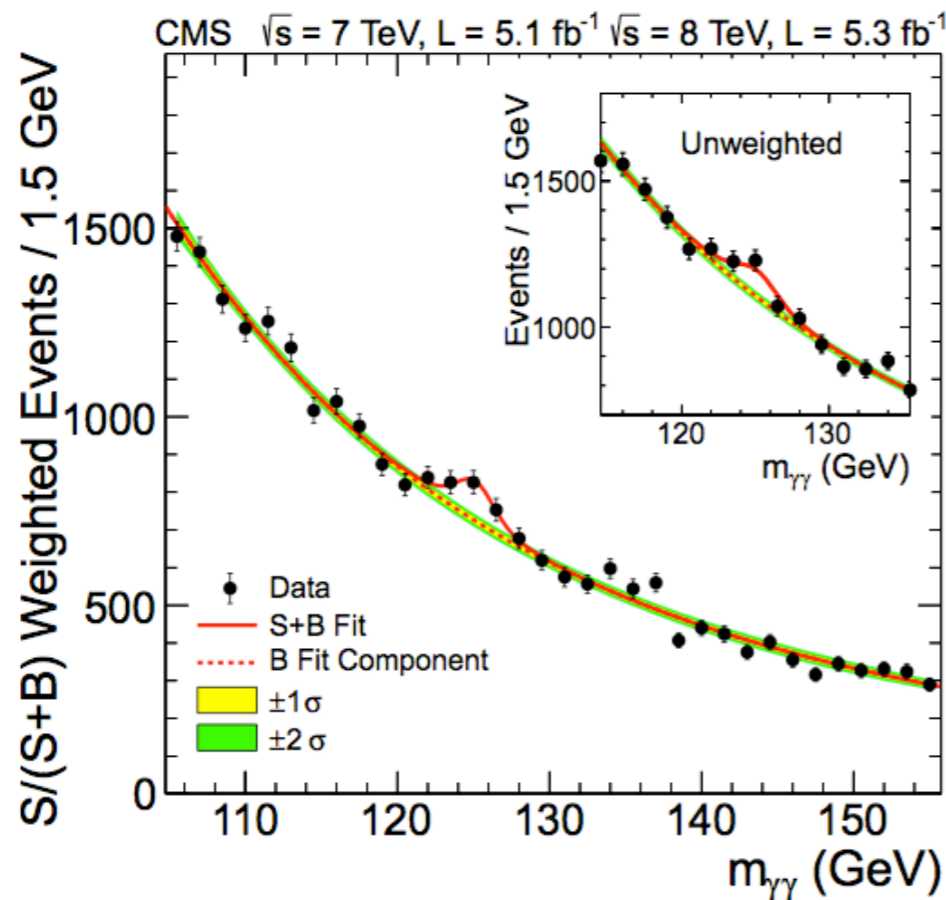
1. Fermi Gamma Ray Line at 130 GeV from Axion-Mediated Dark Matter  
with Hyun Min Lee, Wan-il Park (arXiv:1205.467)
2. Axion-mediated dark matter and Higgs diphoton signal  
with Hyun Min Lee, Wan-il Park (arxiv: 1209.1955)
3. Interplay between Fermi gamma-ray lines and the LHC  
with Hyun Min Lee, Veronica Sanz (arxiv:12XX.XXXX)

# Contents

- A reminder: New boson to two photon
- Fermi  $\gamma$  ray line  $\sim 130\text{GeV}$
- Axion mediated dark matter
- Extra particles in a loop.
- Link to the “higgs  $\rightarrow$  diphoton”

# A bump hunting

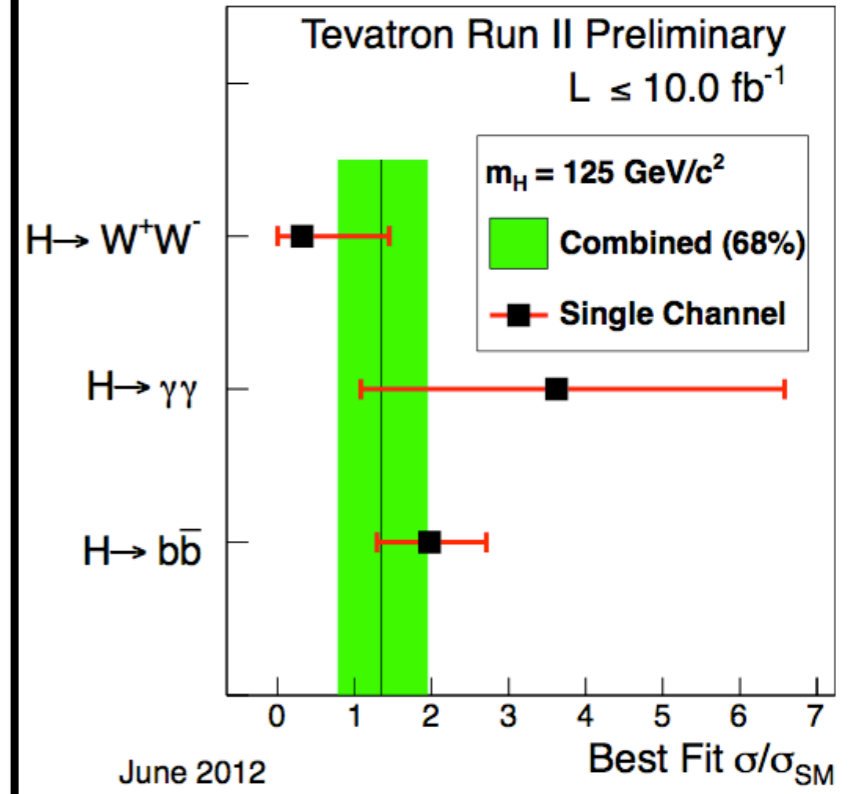
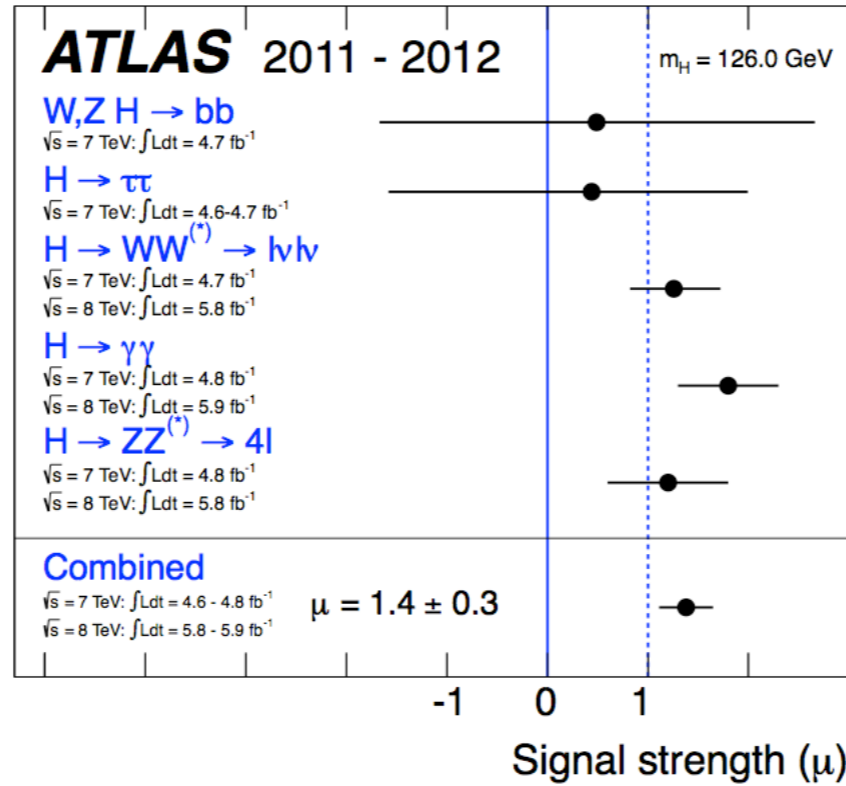
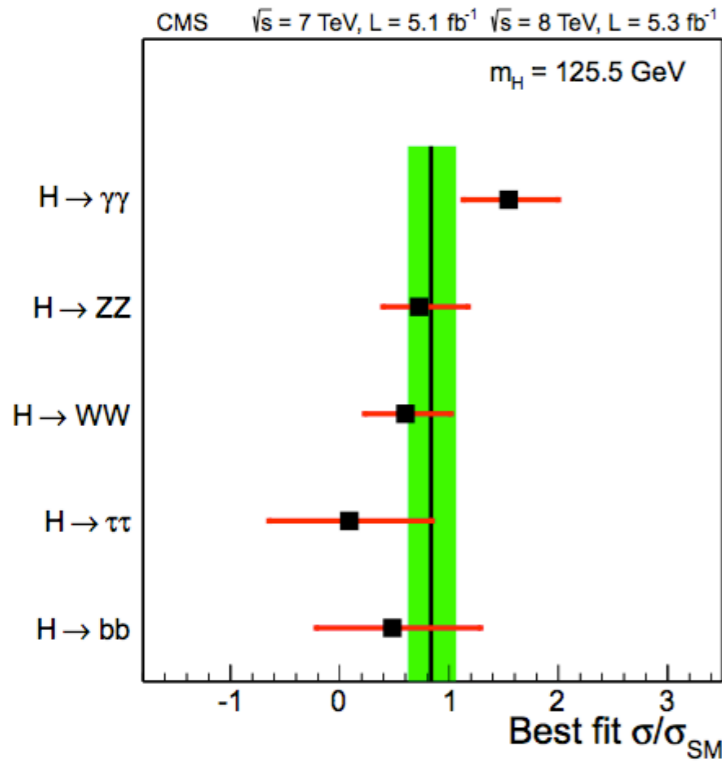
- We got the “bump” at the LHC on the 4th of July



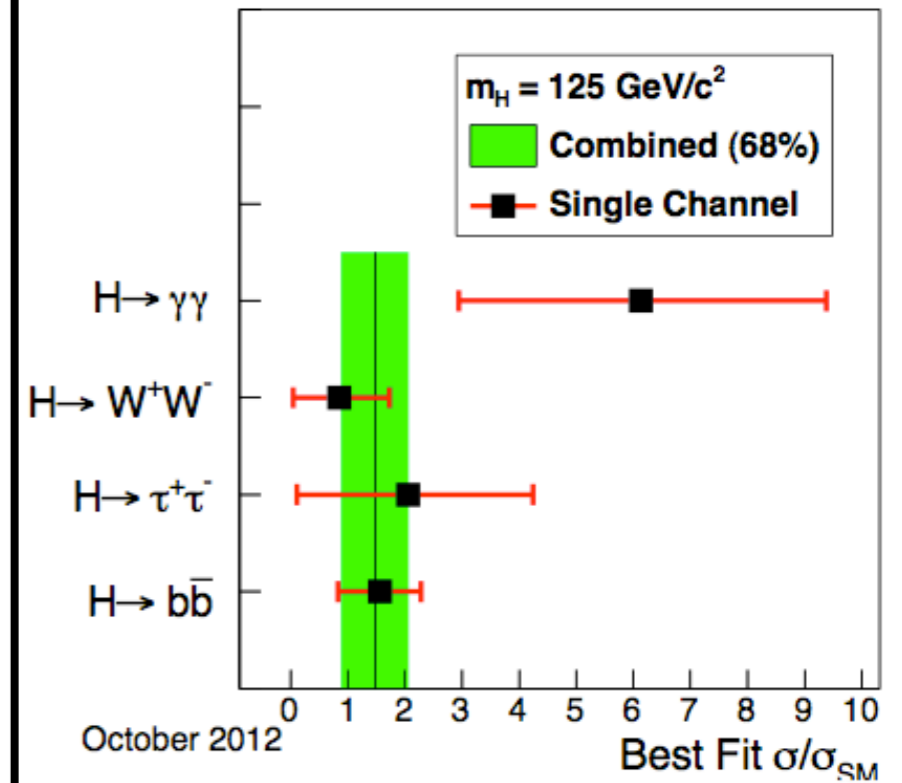
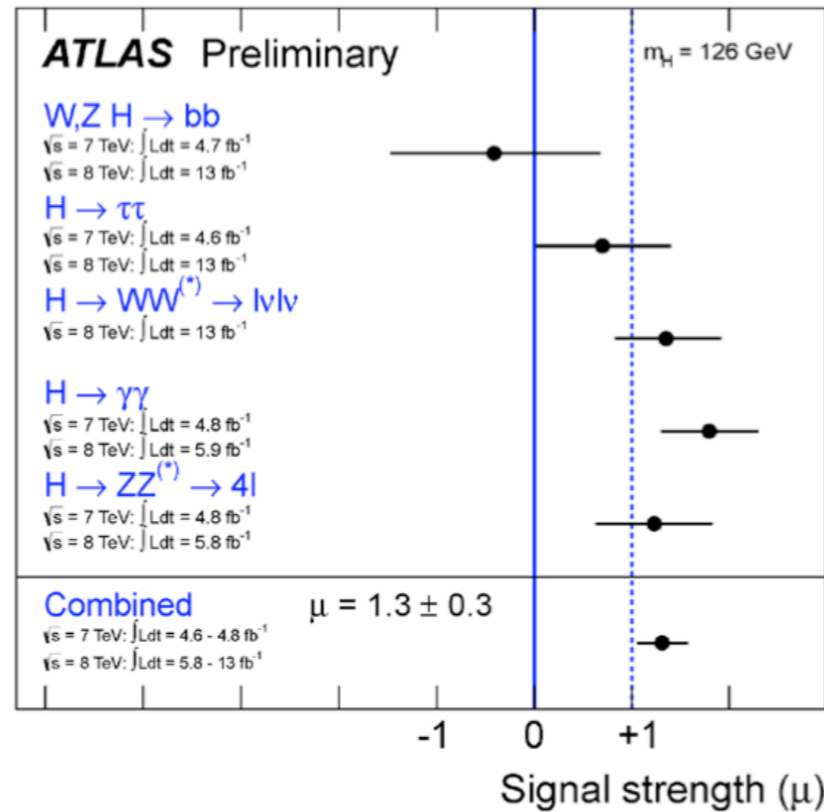
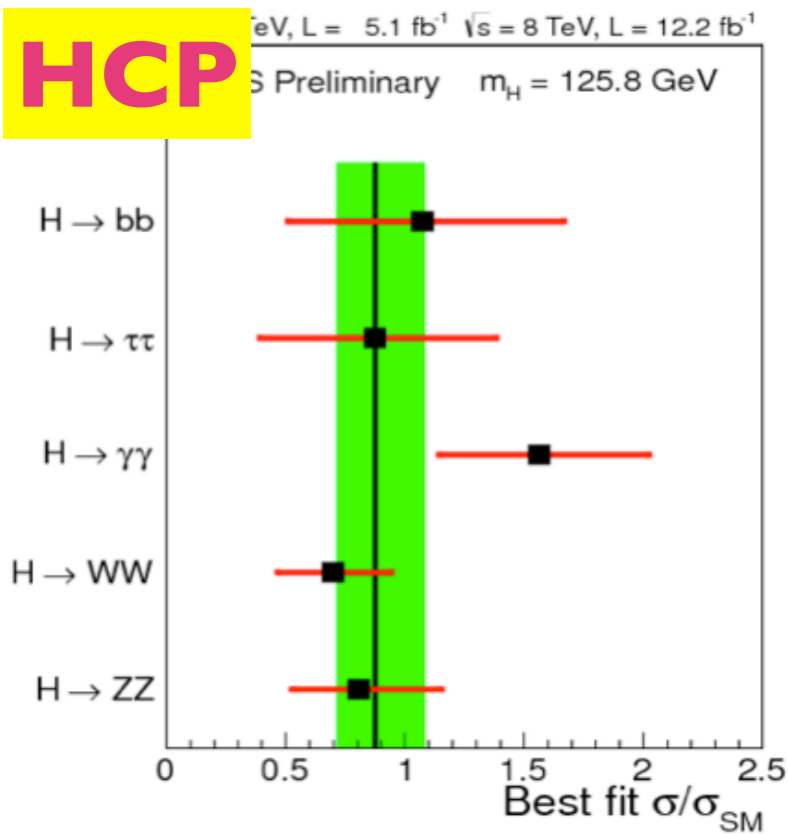
Four lepton channel is the golden channel to study the property of newly discovered boson, but a **di-photon channel** was problematic (compared to “expected higgs boson”)

4th of July

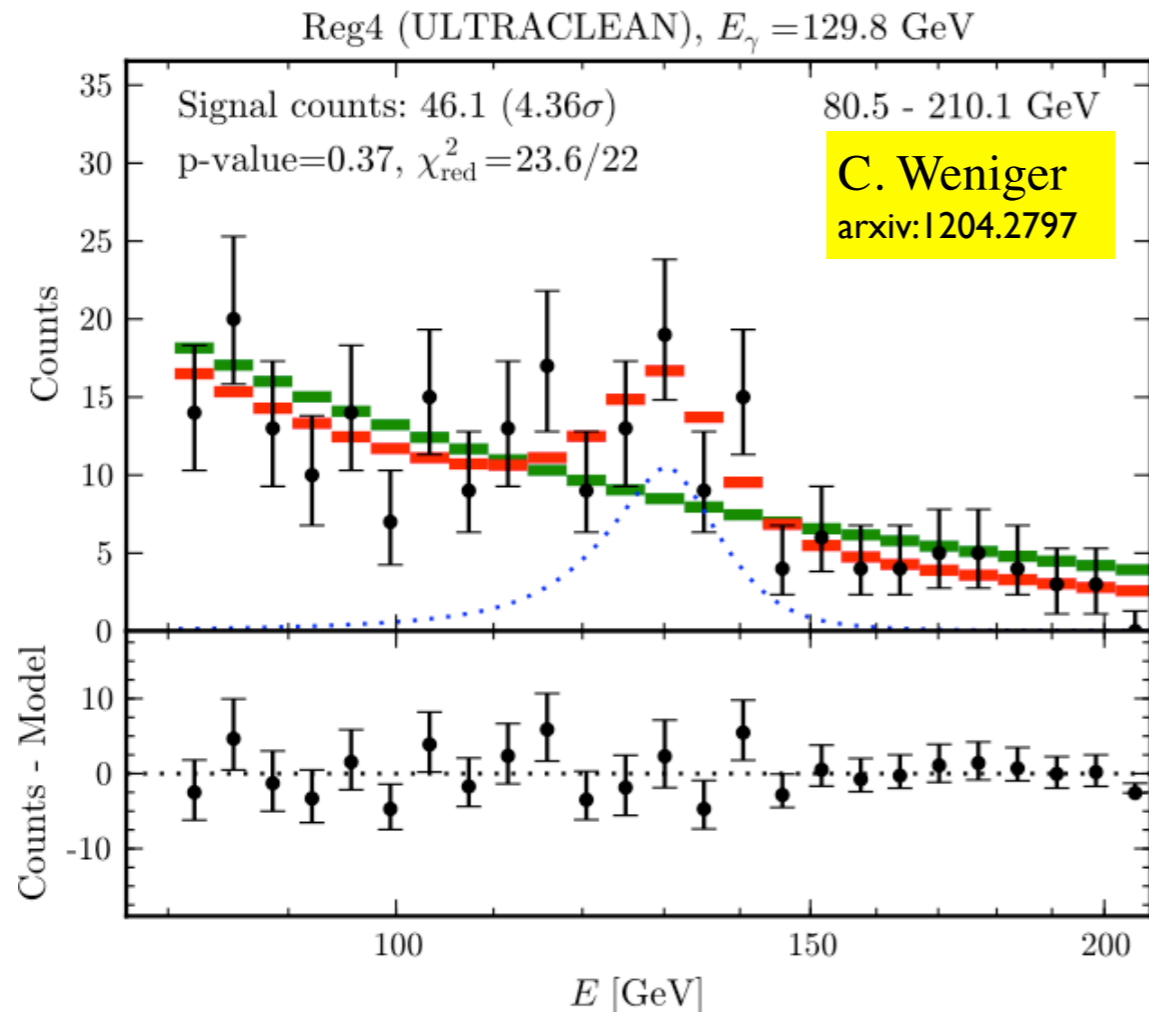
# Di-photon



HCP



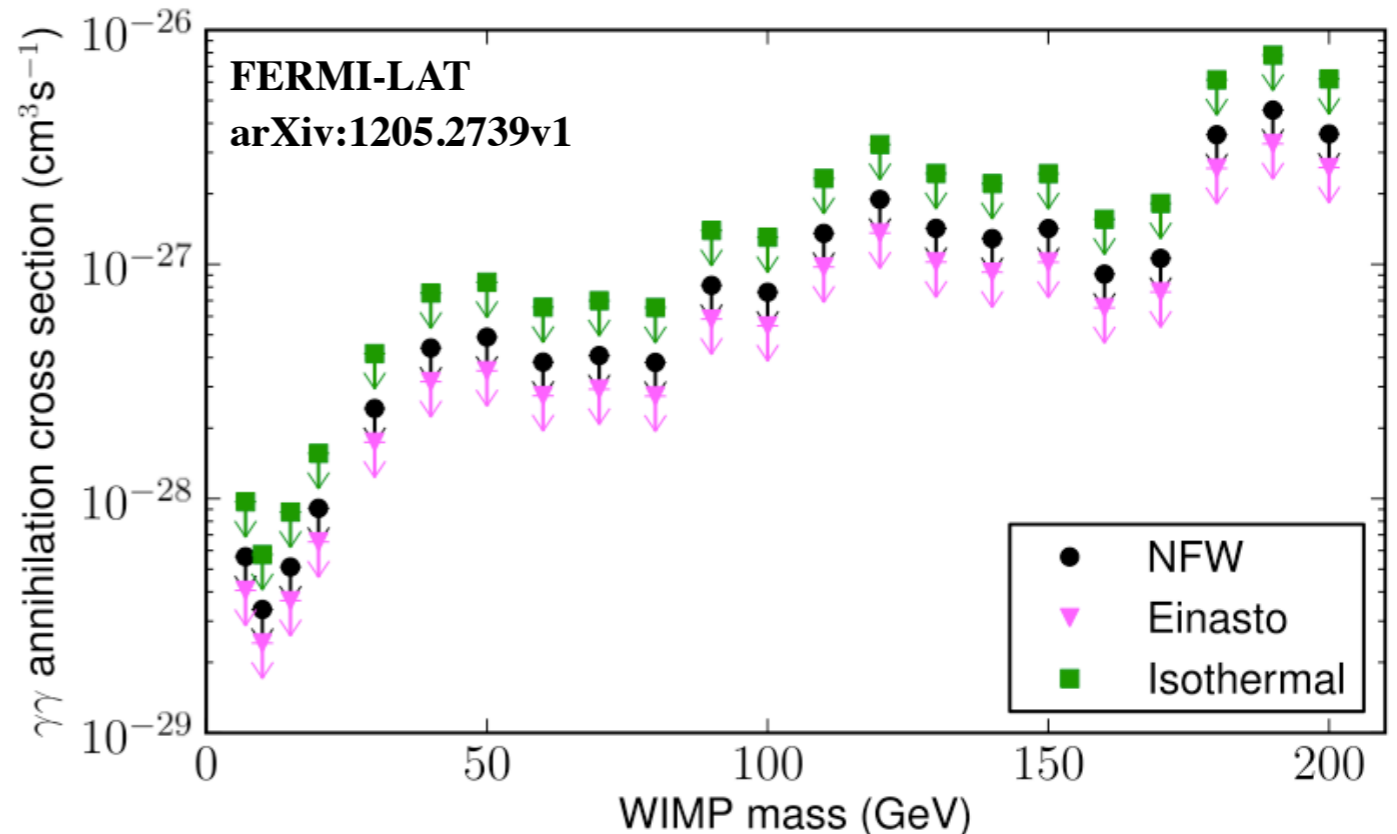
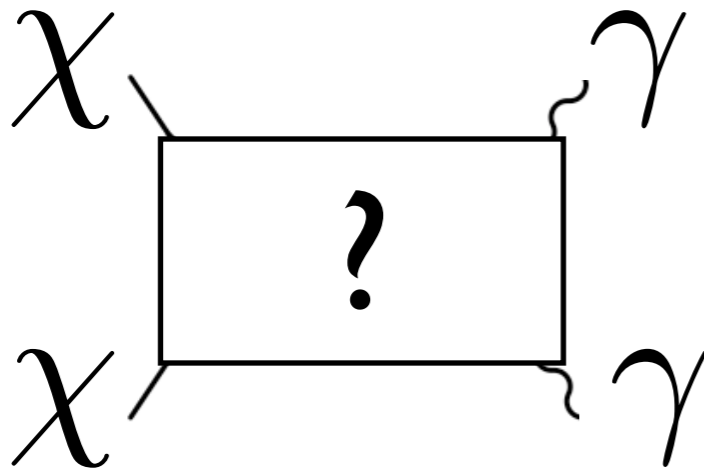
# Fermi $\gamma$ ray line



- An excess of  $E_\gamma$  around 130 GeV have been reported from Fermi Large Area Telescope (Fermi LAT) based on the data; from Aug. 4 2008 to Mar. 8, 2012

- Interpretation with dark matter annihilation provides:

$$\langle \sigma v \rangle_{\gamma\gamma} = (1.27 \sim 2.27) \times 10^{-27} \text{ cm}^3 \text{ s}^{-1}$$



# Set up

- We required a s-channel annihilation to explain a photon line  $\sim 130\text{GeV}$
- Fermion dark matter + pseudo scalar “a” in s-channel (s-wave)
- We used “electro-weak” axion to avoid large annihilation of dark matter pair to gluon.
- We included a CP-even singlet in our consideration.
- An axion gets a weak scale mass  $\rightarrow$   
Cross section enhancement by a resonance effect (e.g. high scale PQ breaking)

# Axion Mediated model

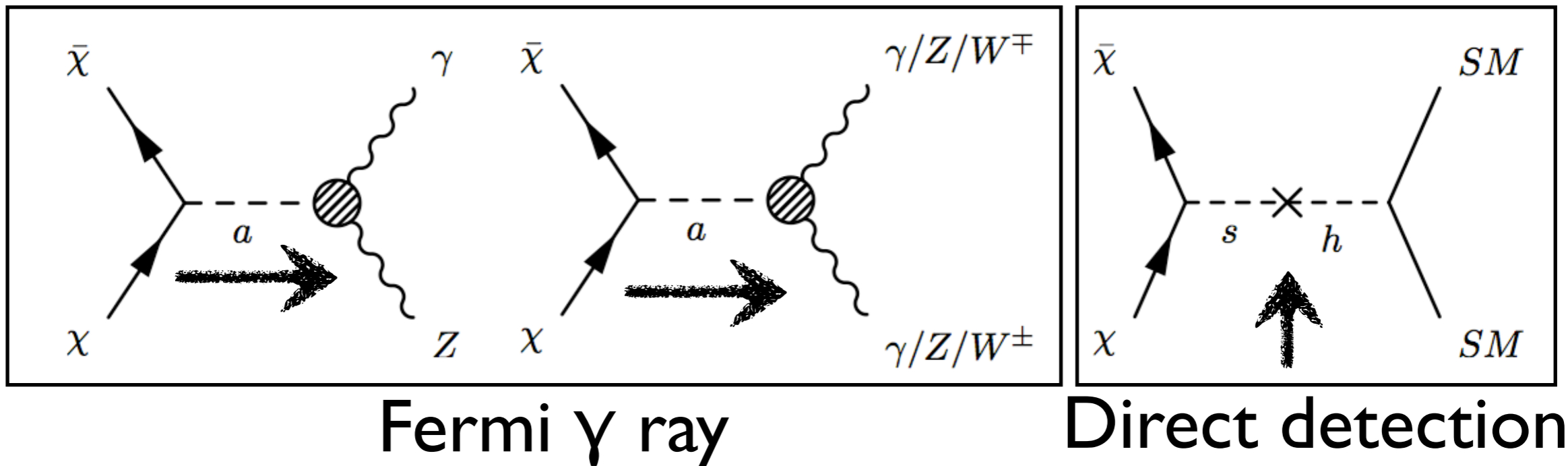
- Dark matters to the photons (SM)

$$\mathcal{L}_{\text{int}} = \lambda_{\chi}(S\bar{\chi}P_L\chi + S^*\bar{\chi}P_R\chi) + \sum_{i=1,2} \frac{c_i\alpha_i}{8\pi v_s} a F_{\mu\nu}^i \tilde{F}^{i\mu\nu}$$

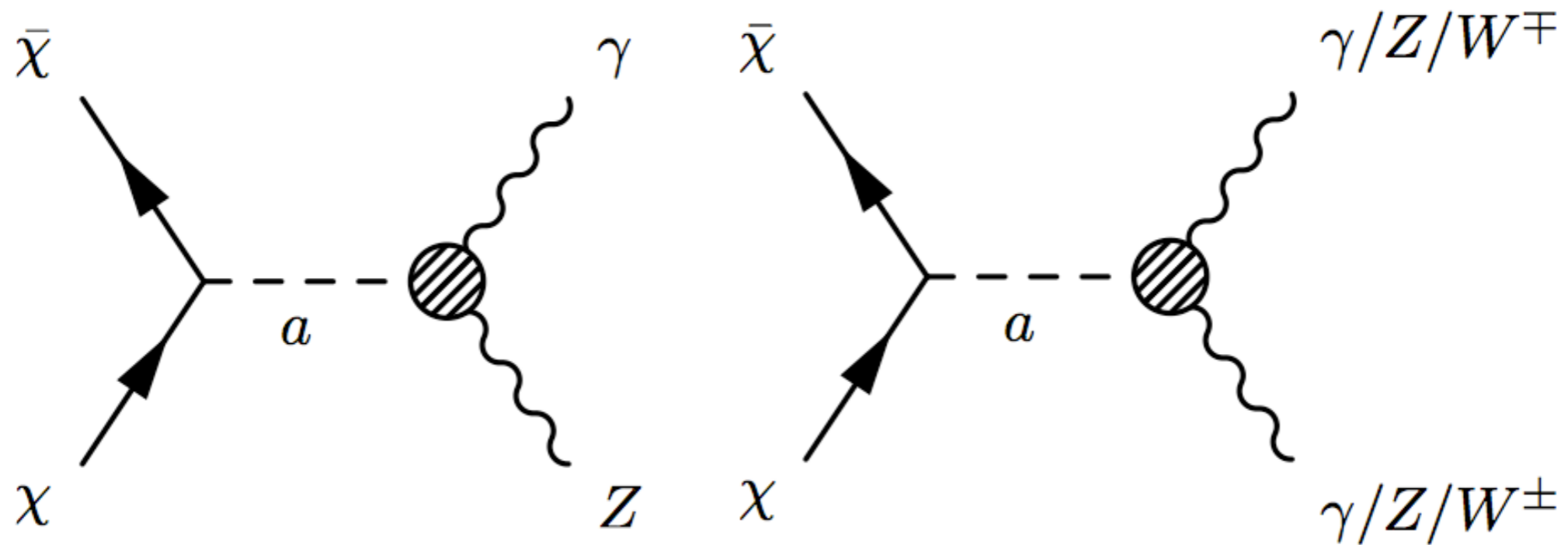
with  $S = (s+ia)/\sqrt{2}$

and potential (+ mixing between s and h)

$$V(H, S) = \lambda_H|H|^4 + \lambda_S|S|^4 + 2\lambda_{HS}|S|^2|H|^2 + m_H^2|H|^2 + m_S^2|S|^2 - \left(\frac{1}{2}m_S'^2 S^2 + \text{h.c.}\right)$$



# Photon line(s)



- We have:

1. DM+DM  $\rightarrow$  photon+photon

$$E_\gamma = M_\chi$$

2. DM+DM  $\rightarrow$  photon+Z

$$E_\gamma = M_\chi \left( 1 - \frac{m_Z^2}{4M_\chi^2} \right)$$

3. DM+DM  $\rightarrow$  (Z, Z) (W, W):

continuum spectrum

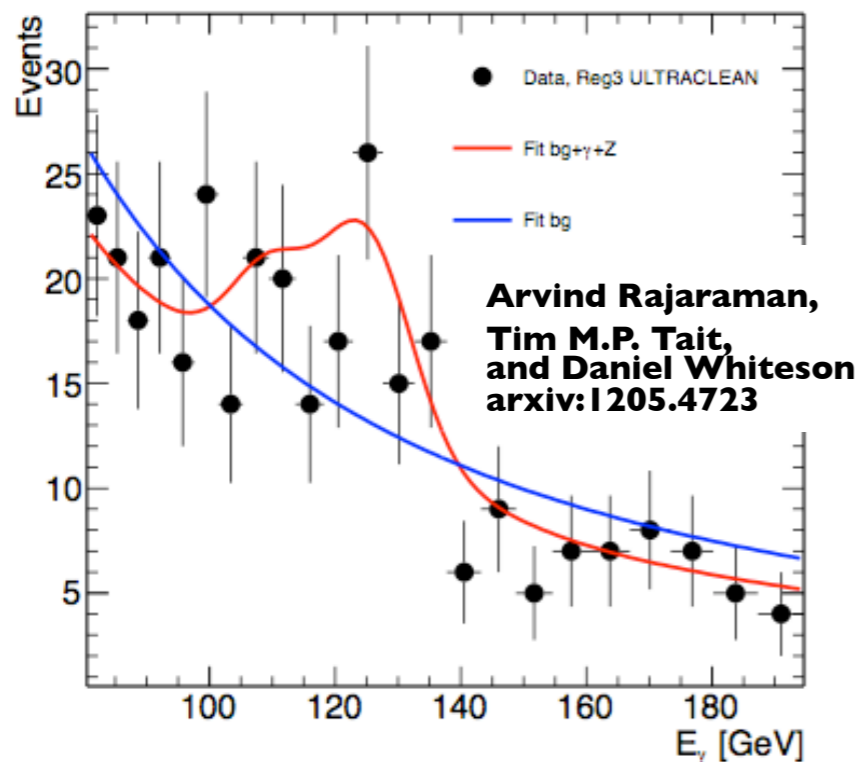


# Photon line(s)

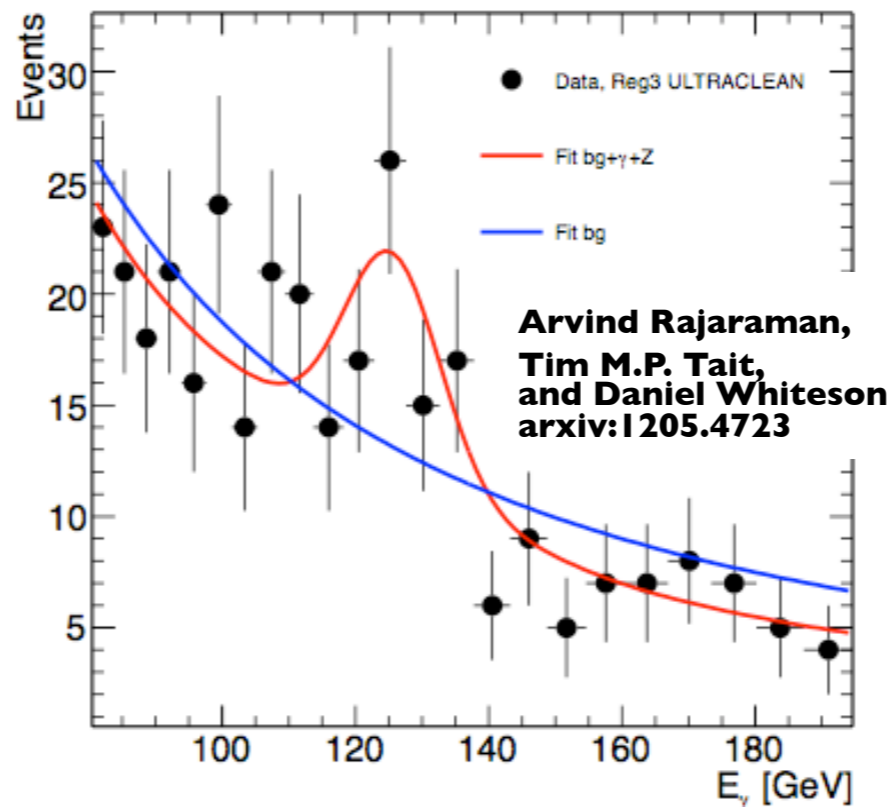
- If  $M_\chi = E_\gamma \sim 130\text{GeV}$ , then

$$E_\gamma = M_\chi \left( 1 - \frac{m_Z^2}{4M_\chi^2} \right) \sim 114\text{GeV}$$

$m_\chi = 130\text{ GeV}$ ,  $N_{\gamma\gamma} = 53.3$ ,  $N_{\gamma Z} = 23.0$ ,  $\text{signif} = 3.47\sigma$



$m_\chi = 145\text{ GeV}$ ,  $N_{\gamma\gamma} = 0.0$ ,  $N_{\gamma Z} = 53.6$ ,  $\text{signif} = 3.60\sigma$



In our case, when  $c_1 = c_2$  the intensity ratio will be

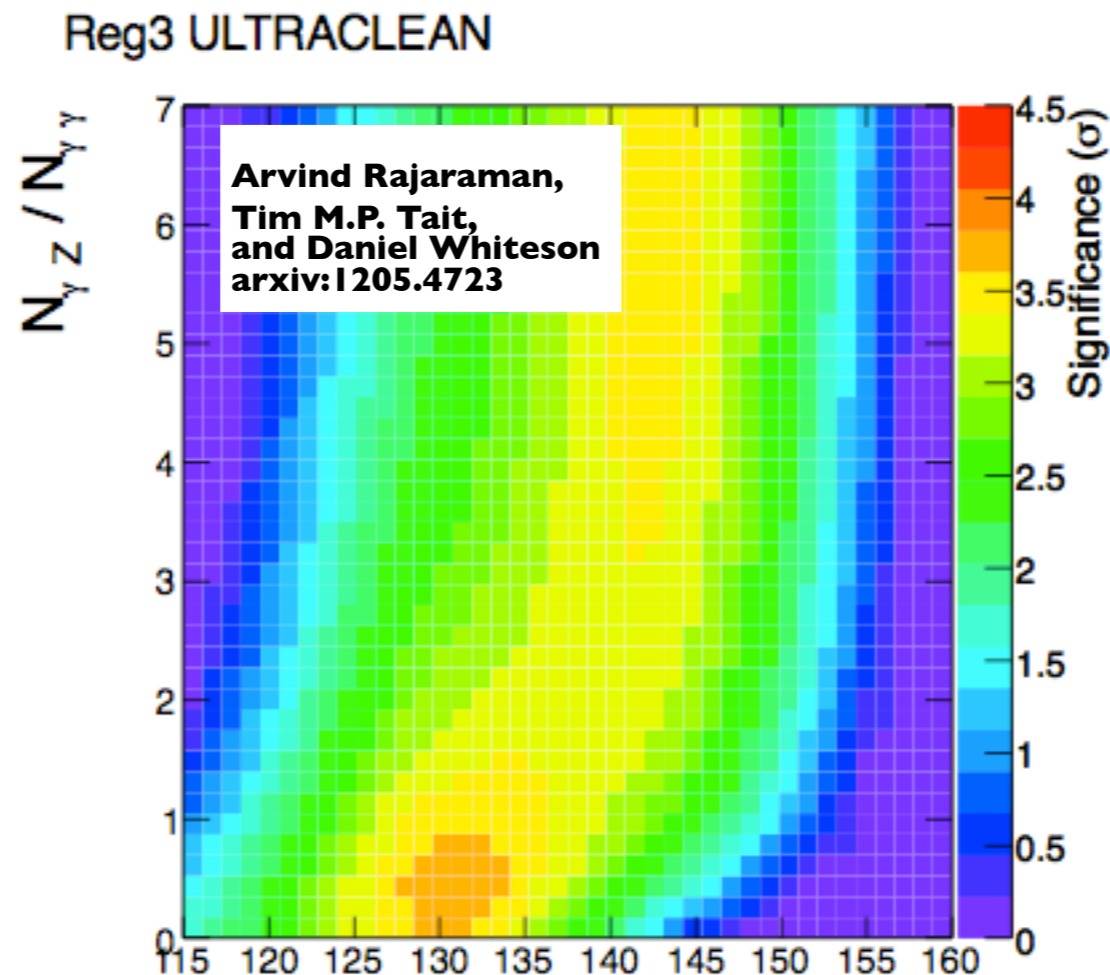
$$\frac{Z\gamma}{\gamma\gamma} = 0.27$$

$$\sum_{i=1,2} \frac{c_i \alpha_i}{8\pi v_s} a F_{\mu\nu}^i \tilde{F}^{i\mu\nu}$$

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# Constraints

- Branch ratio of (DM,DM $\rightarrow\gamma\gamma$ ) are estimated by  $\langle\sigma v\rangle_{\gamma\gamma} = (1.27 \sim 2.27) \times 10^{-27} \text{ cm}^3 \text{ s}^{-1} \sim 4\text{-}8\%$

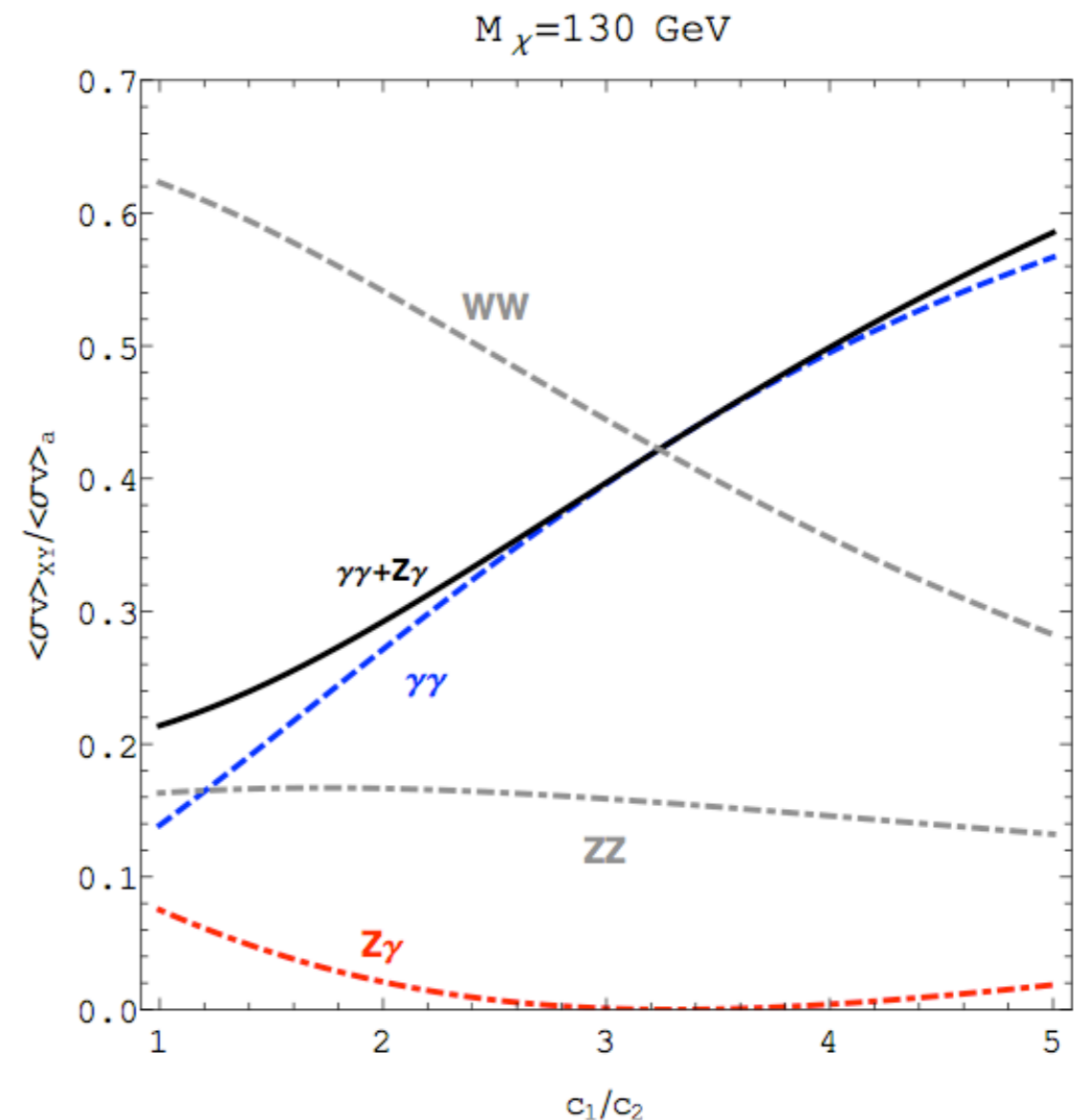
$$\text{Br}(\bar{\chi}\chi \rightarrow XY)_a \equiv \frac{\langle\sigma v\rangle_{XY}}{\langle\sigma v\rangle_a} \sim 14\%$$

$$\langle\sigma v\rangle|_{\text{fr}} \simeq \langle\sigma v\rangle_a|_{\text{fr}} + \langle\sigma v\rangle_s|_{\text{fr}}$$

$$4\sim 8\% = \frac{\langle\sigma v\rangle_{XY}}{\langle\sigma v\rangle_{\text{fr}}} = \frac{\text{Br}(\bar{\chi}\chi \rightarrow XY)_a}{1 + \frac{\langle\sigma v\rangle_s}{\langle\sigma v\rangle_a}|_{\text{fr}}}$$

To be consistent with observation

$$\langle\sigma v\rangle_s / \langle\sigma v\rangle_a = 0.8 - 2.3$$



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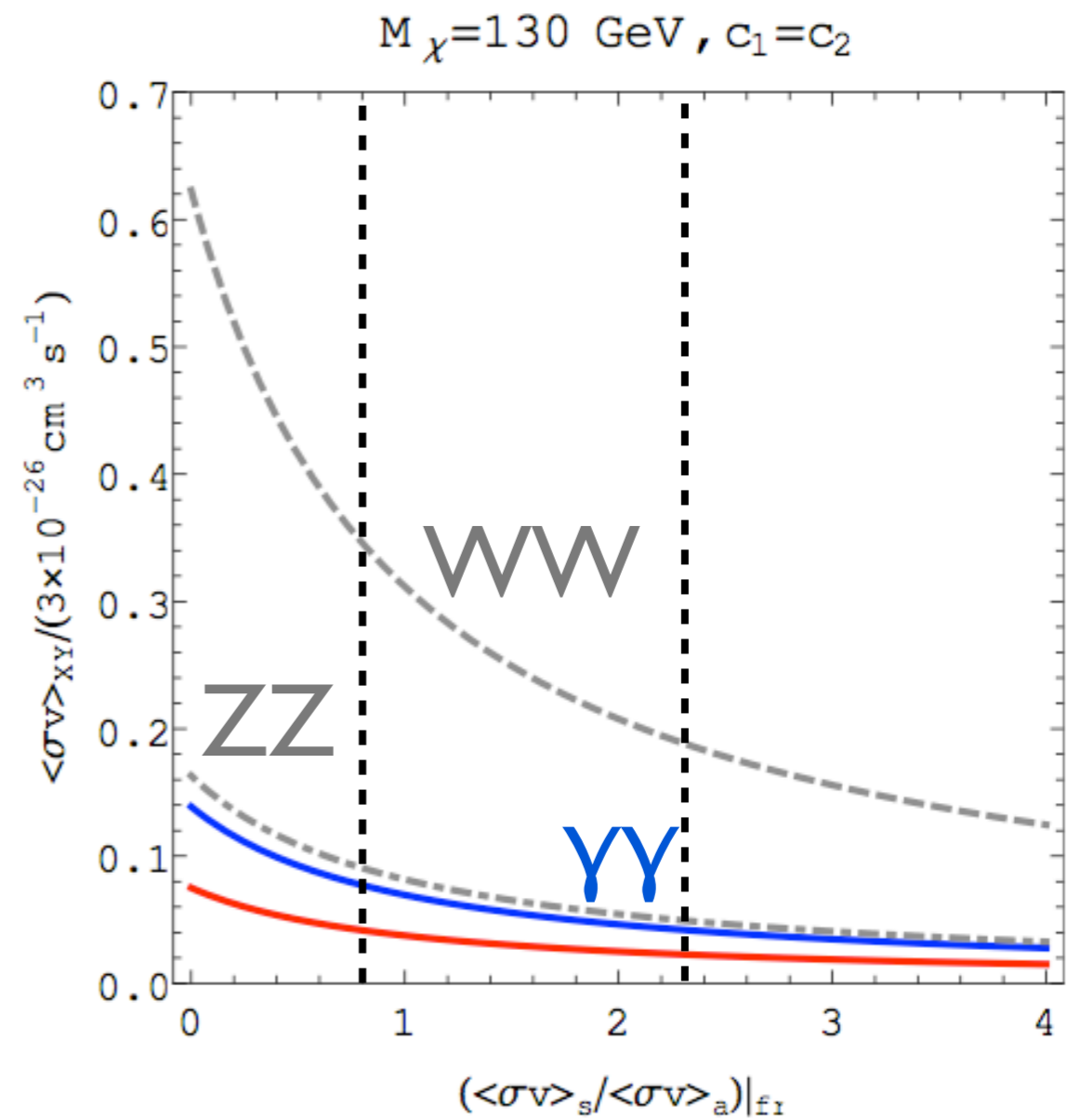
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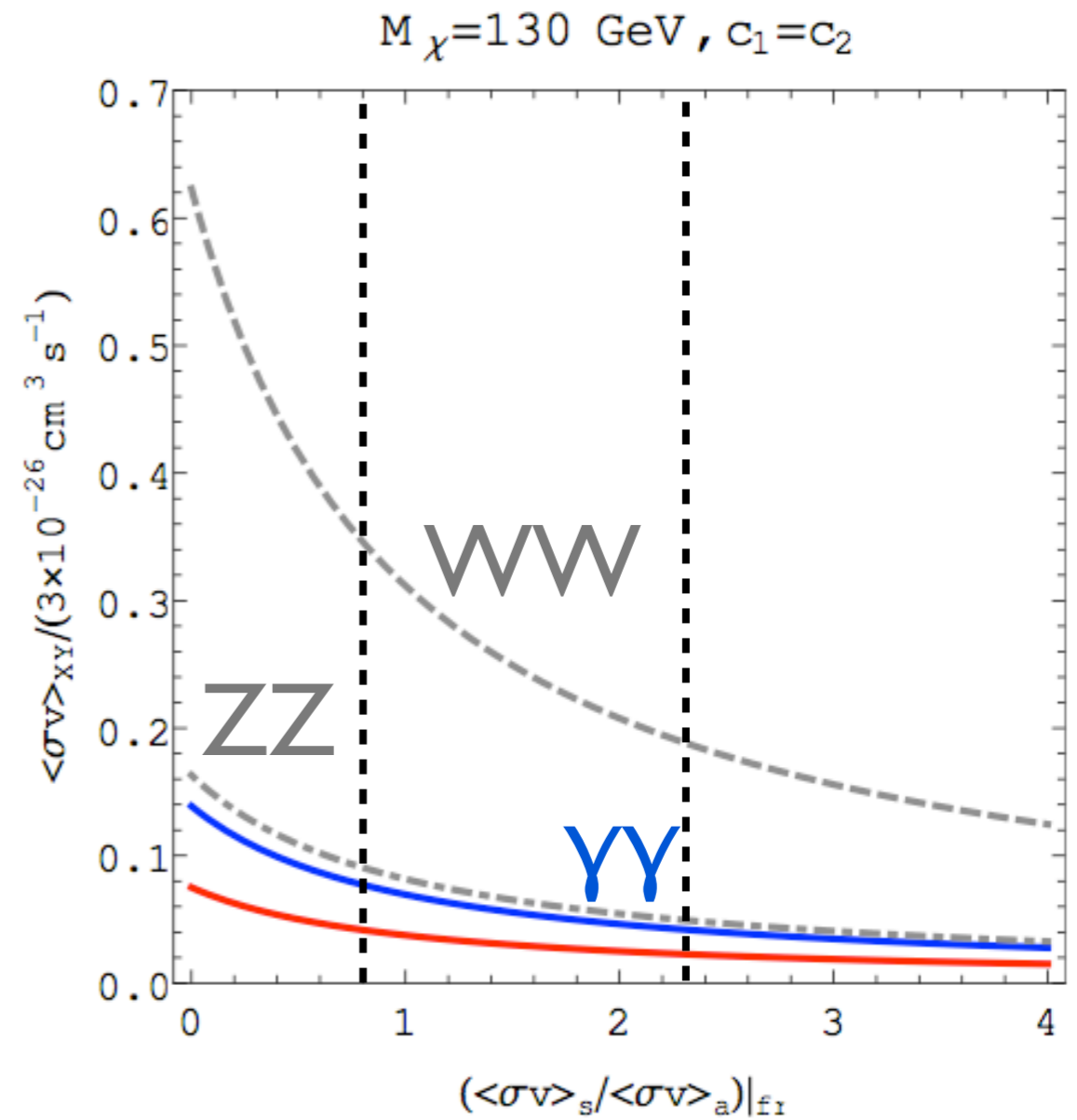
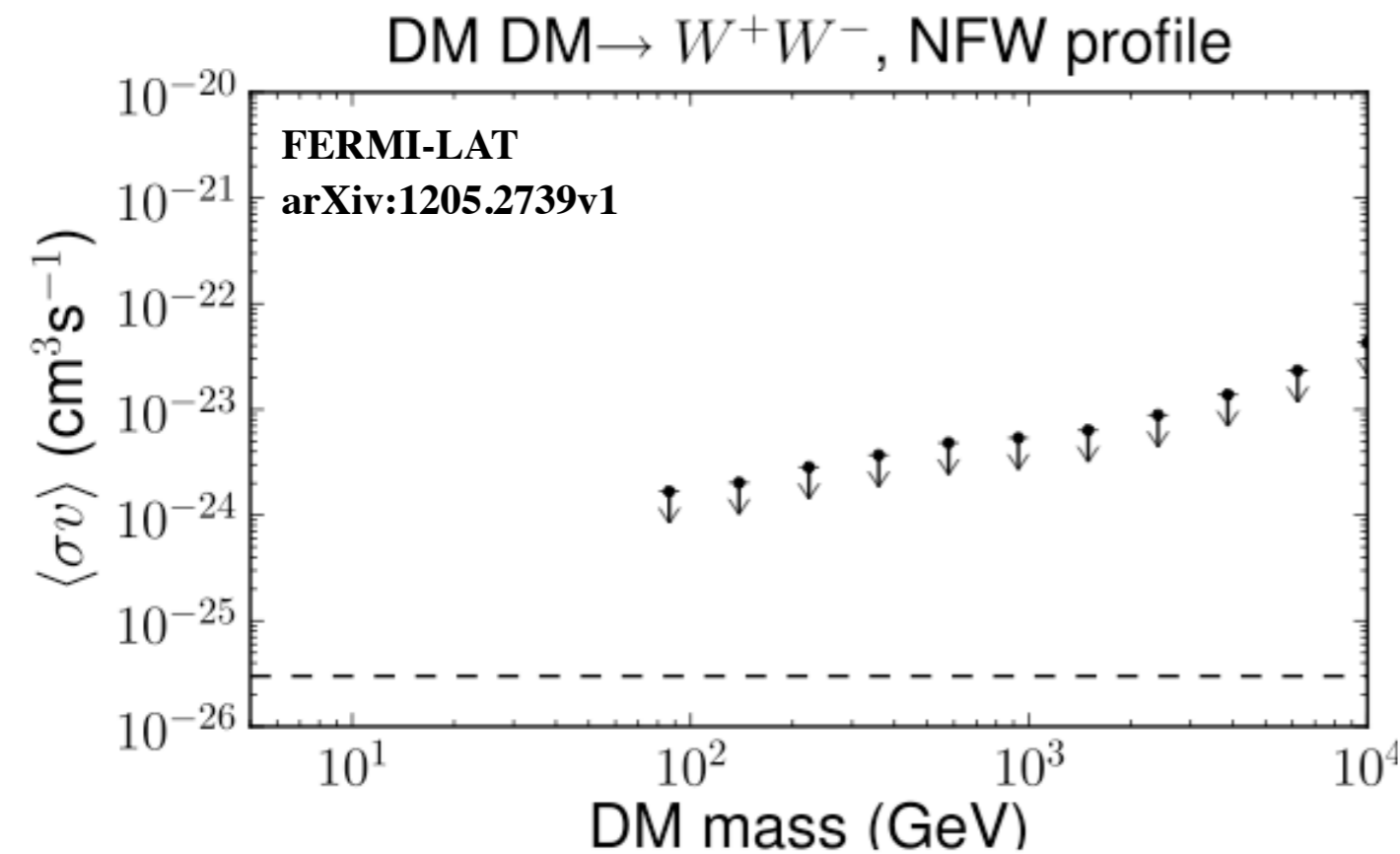
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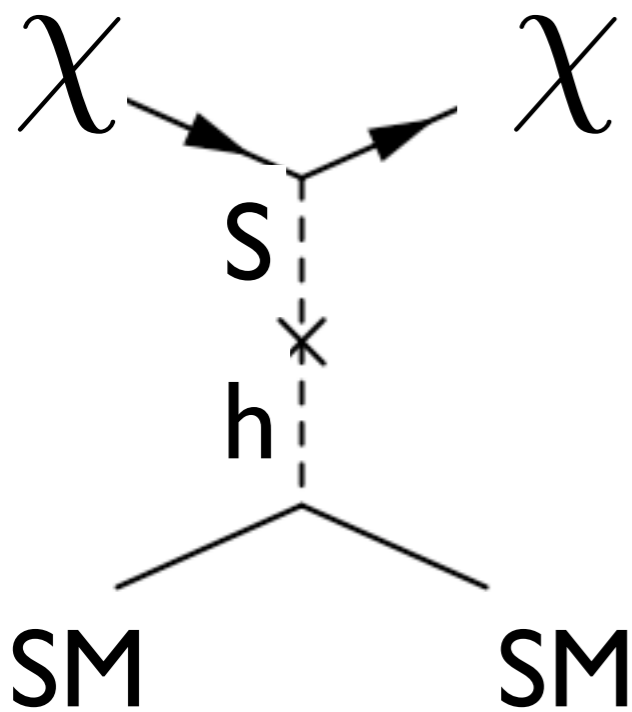
# Constraints

- (DM,DM $\rightarrow$ WW) channel satisfies the current upper limit by Fermi LAT.



# Constraints

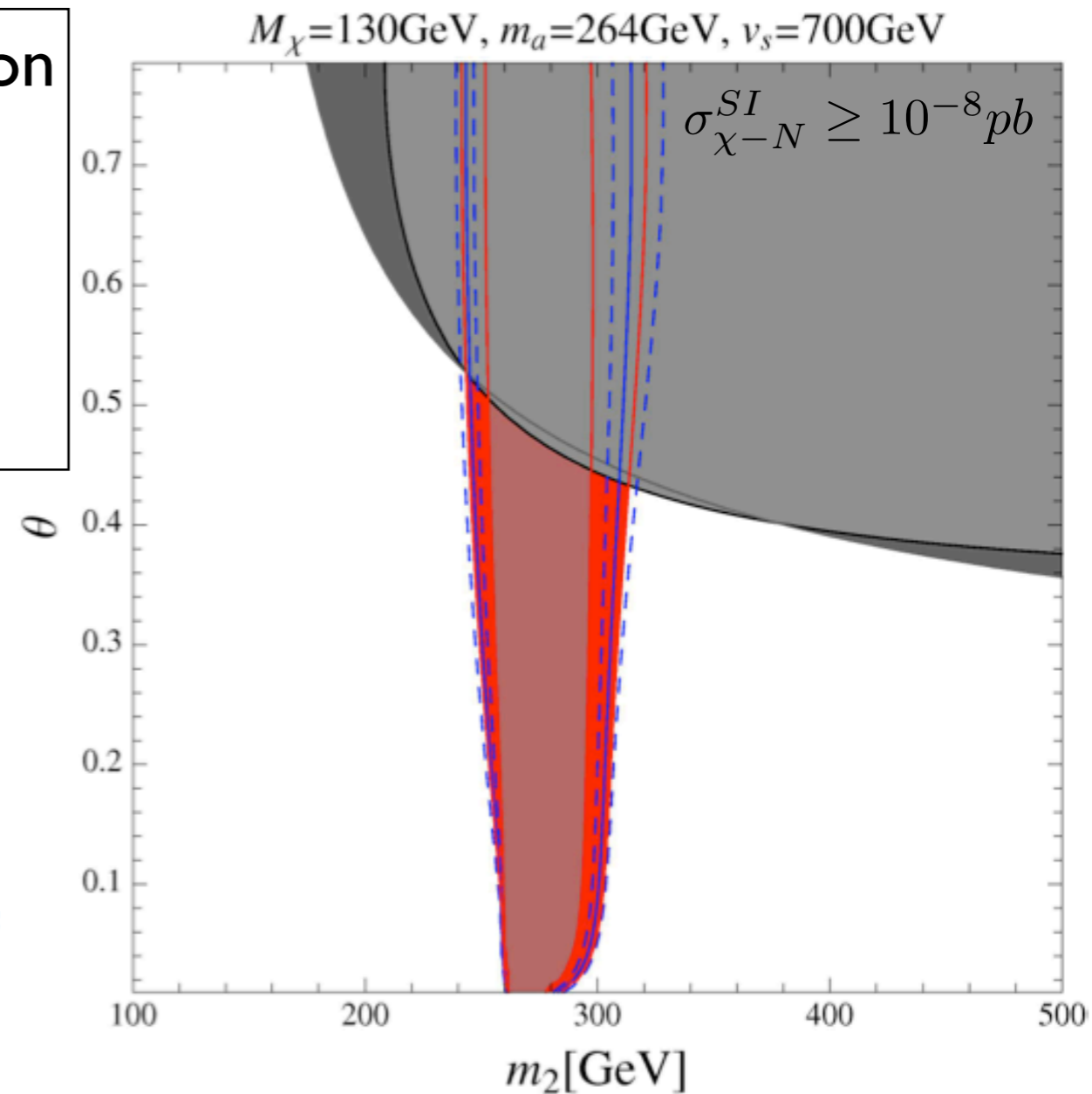
- Direct dark matter detection



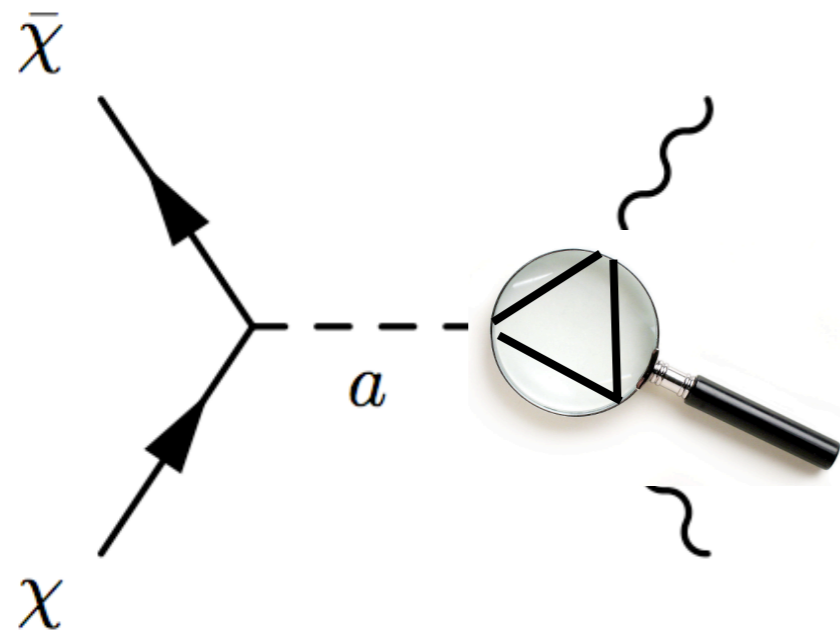
**Grey:** Direct detection  
**DGrey:** EWPD  
**Blue:** WMAP  $3\sigma$   
**Red:**  $\gamma\gamma$  4-8%  
**Pink:**  $\gamma\gamma < 4\%$

$$m_1 \sim 125 \text{ GeV}$$

$$\sigma_{\chi-N}^{\text{SI}} = |\lambda_\chi|^2 \sin^2(2\theta) \cdot \frac{m_r^2 f_N^2}{8\pi} \left(\frac{m_N}{v}\right)^2 \left(\frac{1}{m_1^2} - \frac{1}{m_2^2}\right)^2$$



# There are extra leptons



- We introduce extra vector-like leptons to generate an anomaly.
- We have two higgs doubles to have Yukawa couplings for extra leptons

vector-like lepton doublets,  $l_4, \tilde{l}_4$   
 vector-like lepton singlets,  $e_4^c, \tilde{e}_4^c$

There are three types;

1. vector-like doublet and singlet with PQ charge
2. vector-like doublet and PQ neutral singlet
3. vector-like doublet and triplet with PQ charge

$$\mathcal{L}_{a,\text{eff}} = \sum_{i=1,2} \frac{c_i \alpha_i}{8\pi v_s} a F_{\mu\nu}^i \tilde{F}^{i\mu\nu} \longrightarrow c_1 = \text{Tr}(q_{PQ} Y^2) \text{ and } c_2 = \text{Tr}(q_{PQ} l(r))$$

# Models with extra leptons

Model I:  $-\mathcal{L}_{\text{Yukawa}} = \lambda_\chi S \chi \tilde{\chi} + \lambda_l S l_4 \tilde{l}_4 + \lambda_e S e_4^c \tilde{e}_4^c + y_l H_d l_4 e_4^c - \tilde{y}_l H_u \tilde{l}_4 \tilde{e}_4^c + \text{h.c.}$

Majorana mass  
for neutrino

$$\frac{1}{M} (l_i H_u)(l_j H_u)$$

	$q_i$	$u_i^c$	$d_i^c$	$l_i$	$e_i^c$	$l_4$	$e_4^c$	$\tilde{l}_4$	$\tilde{e}_4^c$	$H_u$	$H_d$	$S$	$\chi$	$\tilde{\chi}$
PQ	1	1	1	2	0	1	1	1	1	-2	-2	-2	1	1
$Z_2$	+	+	-	+	-	-	+	+	+	+	-	-	-	+

Table 1: PQ charges and  $Z_2$  parities for Majorana neutrino case

only dirac mass  
for neutrino

	$q_i$	$u_i^c$	$d_i^c$	$l_i$	$e_i^c$	$N_i^c$	$l_4$	$e_4^c$	$\tilde{l}_4$	$\tilde{e}_4^c$	$H_u$	$H_d$	$S$	$\chi$	$\tilde{\chi}$
PQ	1	1	1	1	1	1	1	1	1	1	-2	-2	-2	1	1
$Z_2$	+	+	-	+	-	+	-	+	+	+	+	-	-	-	+

Table 2: PQ charges and  $Z_2$  parities for Dirac neutrino case

Possible mixing between leptons by  $H_d l_i e_4^c$  exist. We introduced  $Z_2$  to protect mixings.

$$\mathcal{L}_{a,\text{eff}} = \sum_{i=1,2} \frac{c_i \alpha_i}{8\pi v_s} a F_{\mu\nu}^i \tilde{F}^{i\mu\nu} \longrightarrow c_1 : c_2 = 3, 1$$



# Models with extra lepton

Model 2: vector-like doublet and PQ neutral singlet

$$- \mathcal{L} = \lambda_\chi S \chi \tilde{\chi} + \lambda_l S l_4 \tilde{l}_4 + m_e e_4^c \tilde{e}_4^c + y_l H_d l_4 e_4^c - \tilde{y}_l H_u \tilde{l}_4 \tilde{e}_4^c + \text{h.c.}$$

$$\mathcal{L}_{a,\text{eff}} = \sum_{i=1,2} \frac{c_i \alpha_i}{8\pi v_s} a F_{\mu\nu}^i \tilde{F}^{i\mu\nu} \longrightarrow c_1 : c_2 = 1, 1$$

Anomaly only comes from vector-like doublet lepton

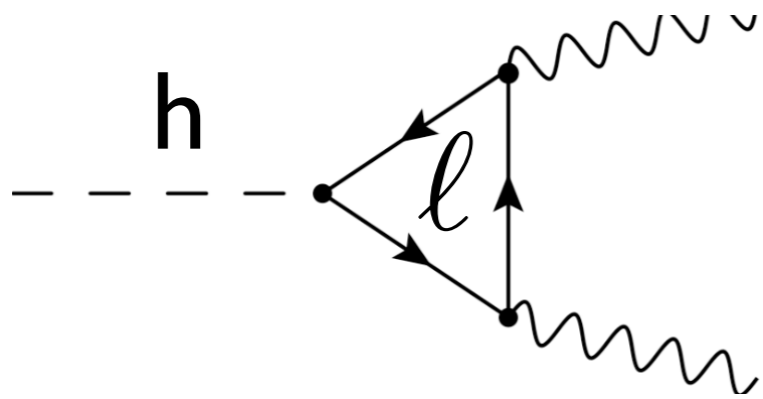
Model 3: vector-like doublet and triplet with PQ charge

$$- \mathcal{L}_{\text{Yukawa}} = \lambda_\chi S \chi \tilde{\chi} + \lambda_l S l_4 \tilde{l}_4 + \lambda_e S e_4^c \tilde{e}_4^c + y_l H_d l_4 T - \tilde{y}_l H_u \tilde{l}_4 T + \text{h.c.}$$

$$\mathcal{L}_{a,\text{eff}} = \sum_{i=1,2} \frac{c_i \alpha_i}{8\pi v_s} a F_{\mu\nu}^i \tilde{F}^{i\mu\nu} \longrightarrow c_1 : c_2 = 3, 1$$

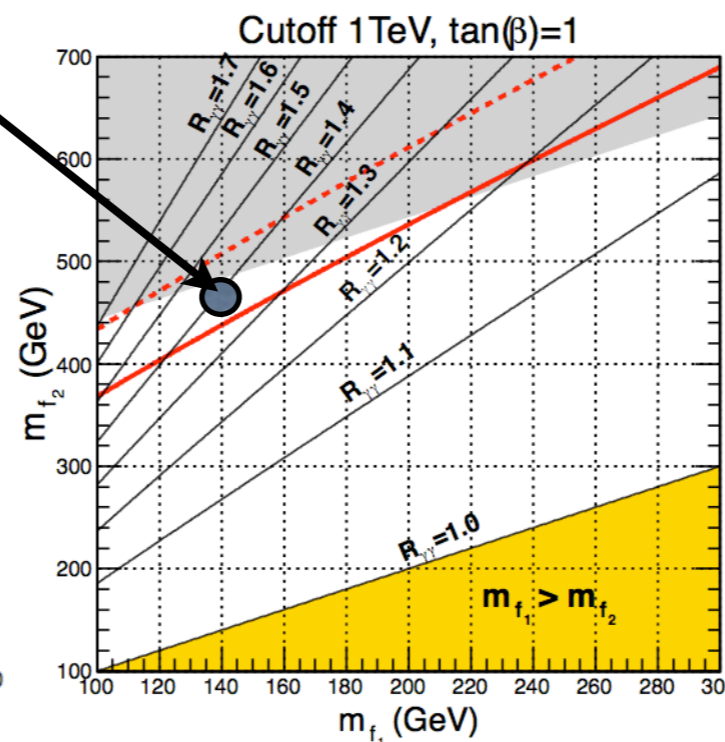
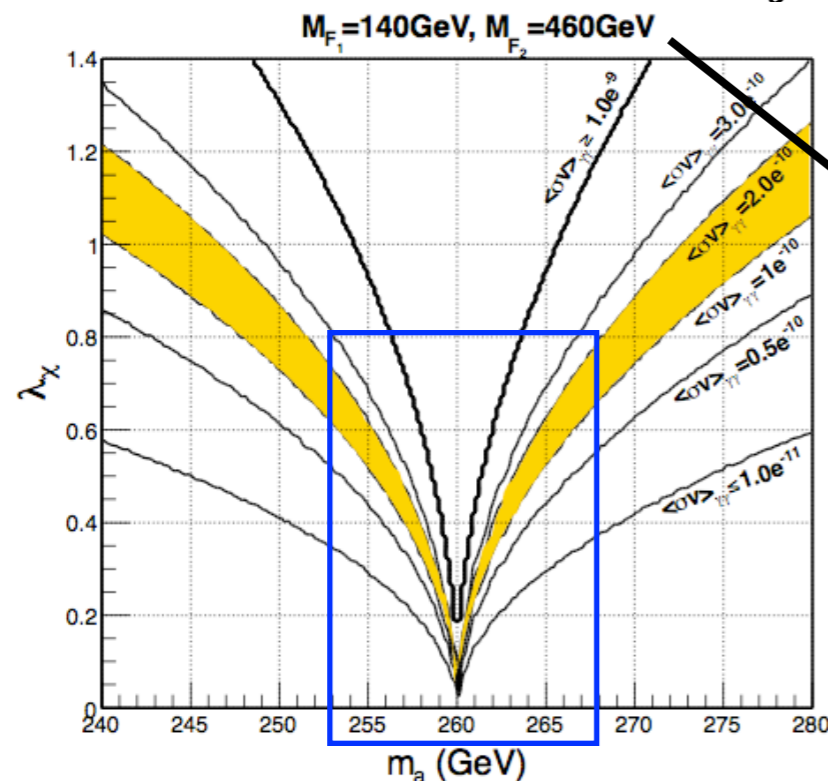
# Link to higgs

- New charged leptons are inside the higgs to two photon loop.



$$R_{\gamma\gamma} = \left| 1 + y_f \frac{v}{m_f} \frac{A_f(\tau_f)}{A_V(\tau_W) + N_c Q_t^2 A_f(\tau_t)} \right|^2$$

is the ratio of the higgs to diphoton ratio compared to SM value



model I

**Grey:** vacuum metastability

**Red:** EWPD ( $1\sigma, 2\sigma$ )

we see the resonant enhancement for the two photon production.

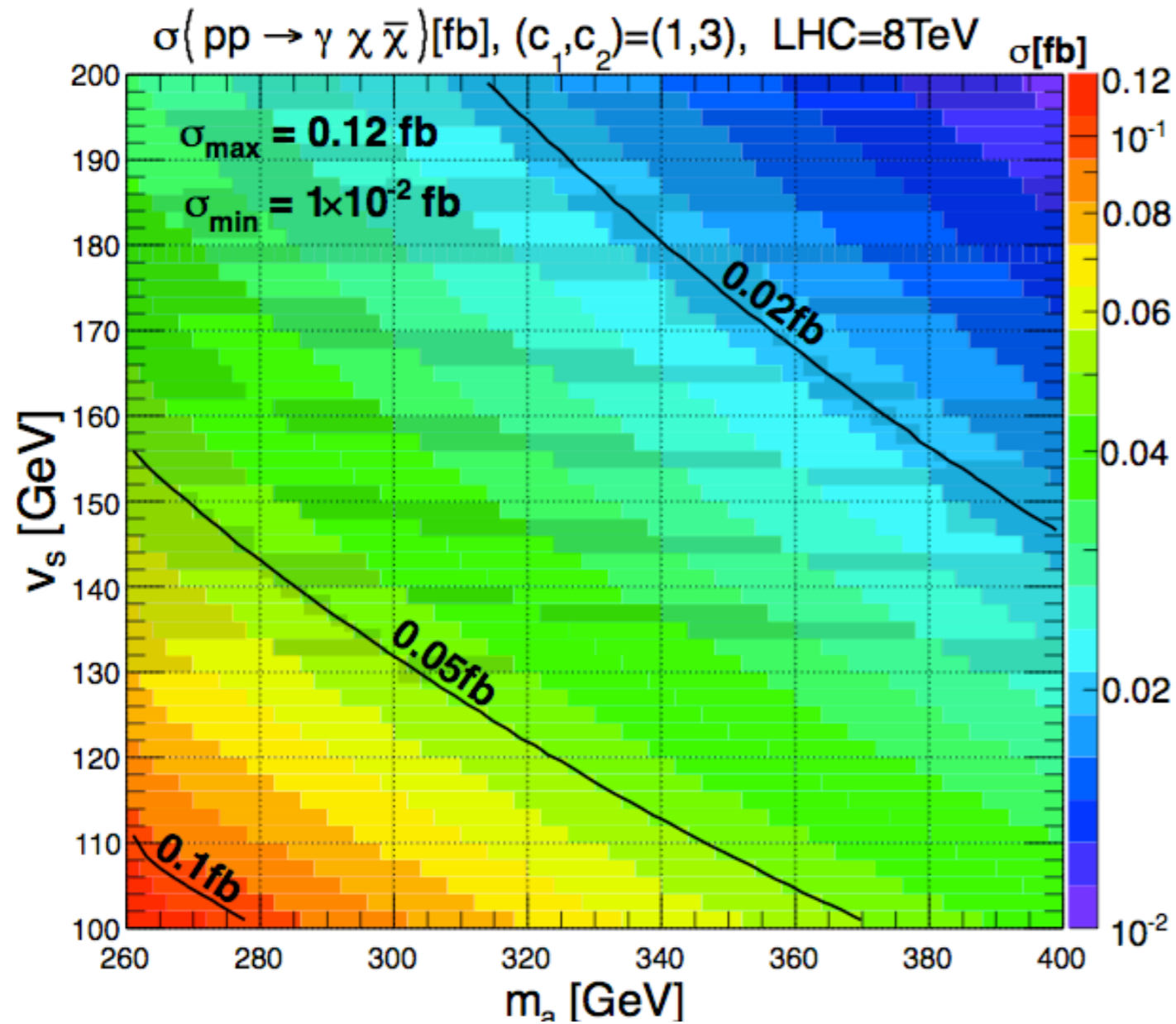
# Summary of model prediction

	Model I	Model II	Model III
$(c_1, c_2)$	$(3, 1)$	$(1, 1)$	$(1, 3)$
$\text{Br}(\bar{\chi}\chi \rightarrow \gamma\gamma)$	$\gtrsim 40\%$	$\gtrsim 14\%$	$\gtrsim 6\%$
$R_{\gamma\gamma}$	$\lesssim 1.5$	$\lesssim 1.5$	$\lesssim 1.4$

- Singlet model (I, II) have larger higgs to two photon enhancement. But for the Fermi gamma ray, they require extra annihilation channels.
- Triplet model (III) does not require extra annihilation channels. It has smaller higgs to two photon enhancement compared to other models.

# Conclusions

- We proposed a singlet fermion dark matter to explain recent Fermi LAT data.
- Extra vector-like leptons can explain Fermi LAT data and higgs to two photon enhancement.
- But we can not see the trace of axion at the LHC yet.



Thank you.

$$\begin{aligned}
 \mathcal{L}_{\text{axion}} = & i\bar{\chi}\gamma^\mu\partial_\mu\chi - m_\chi\bar{\chi}\chi + \frac{1}{2}(\partial_\mu a)^2 - \frac{1}{2}m_a^2 a^2 - \frac{1}{4}F_{\mu\nu}^i F^{i\mu\nu} \\
 & + \frac{1}{\sqrt{2}}\lambda_\chi i a \bar{\chi}\gamma^5\chi + \sum_{i=1,2} \frac{c_i\alpha_i}{8\pi v_s} a F_{\mu\nu}^i \tilde{F}^{i\mu\nu}
 \end{aligned}$$