Testing gravity on cosmological scales with weak lensing

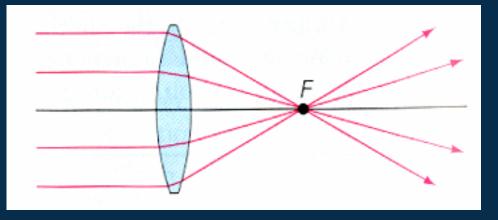
JPL Postdoc Seminar Series August 13, 2009 Ali Vanderveld (JPL, Caltech)

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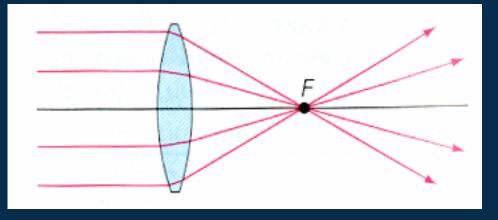
Abell 2218



(1) Lensing via refraction

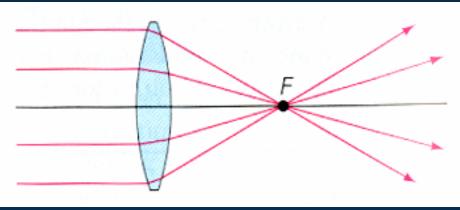


(1) Lensing via refraction



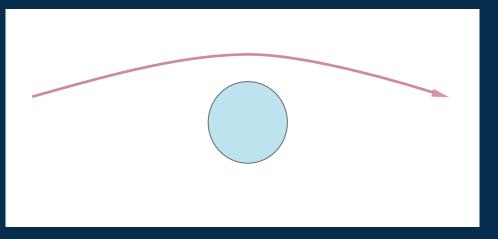
Light rays move at a different speed in the lens medium

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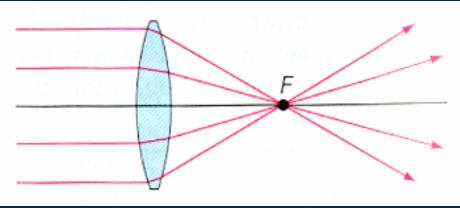


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(2) Lensing via gravitation

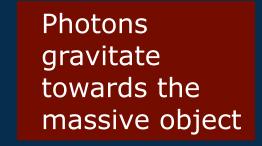


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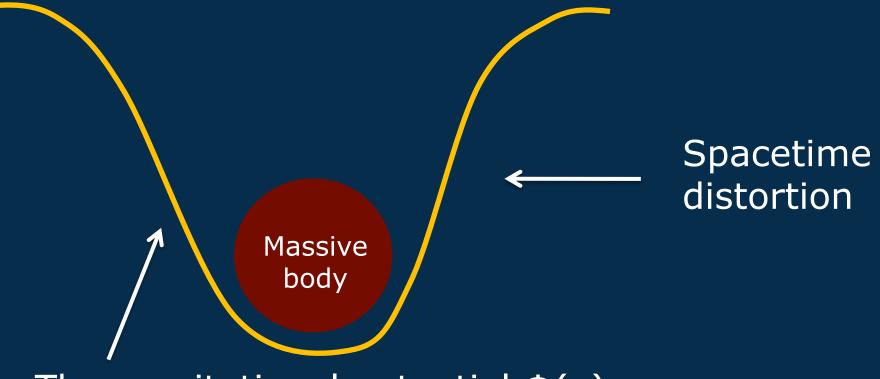


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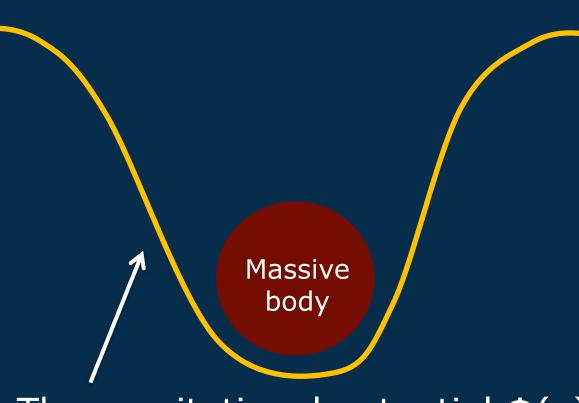
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Massive body



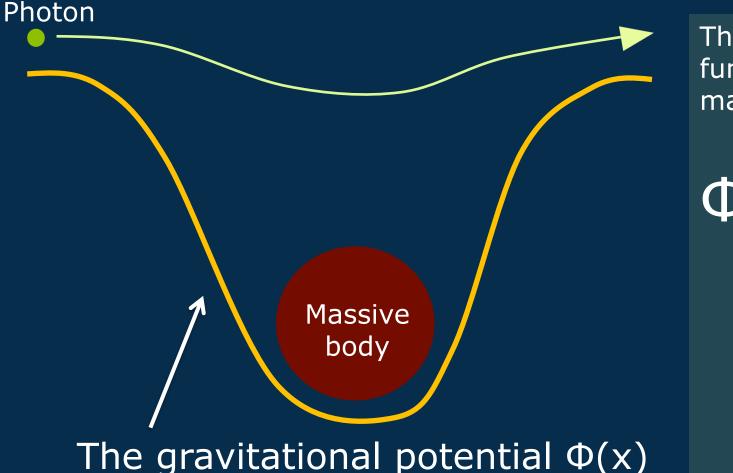
The gravitational potential $\Phi(x)$



The potential is a function of the matter distribution:

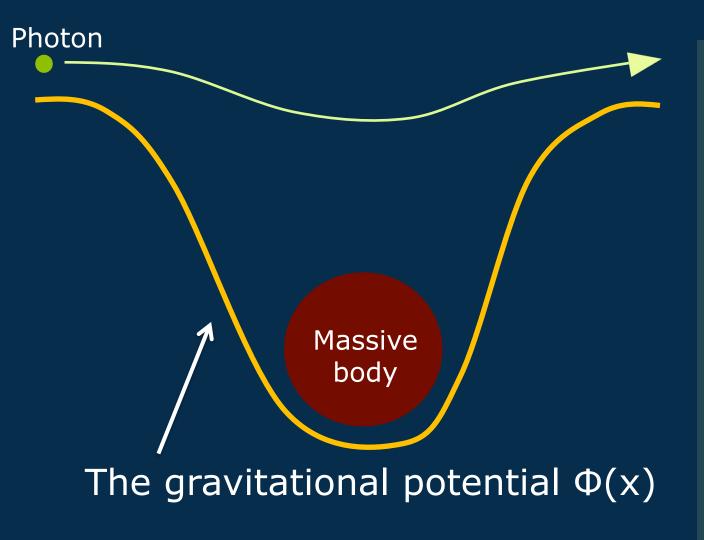
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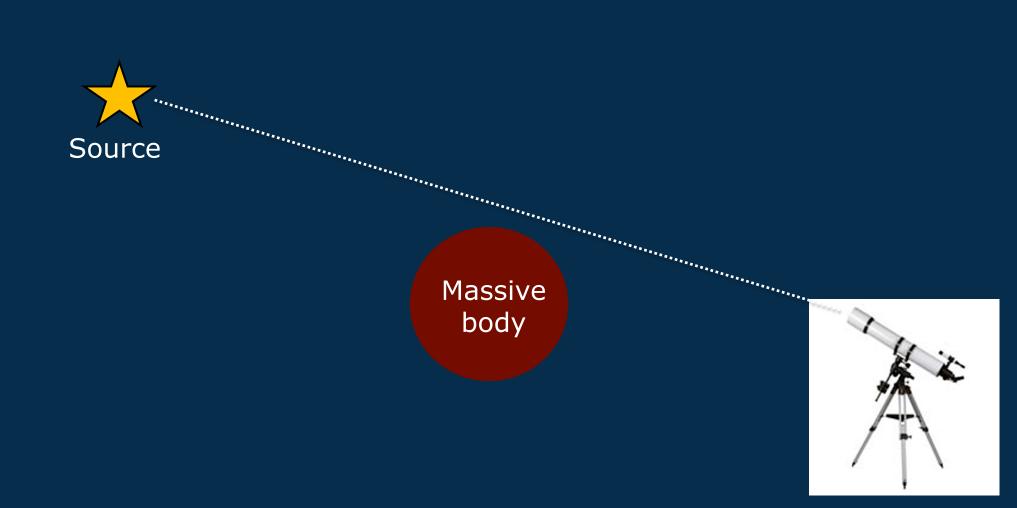
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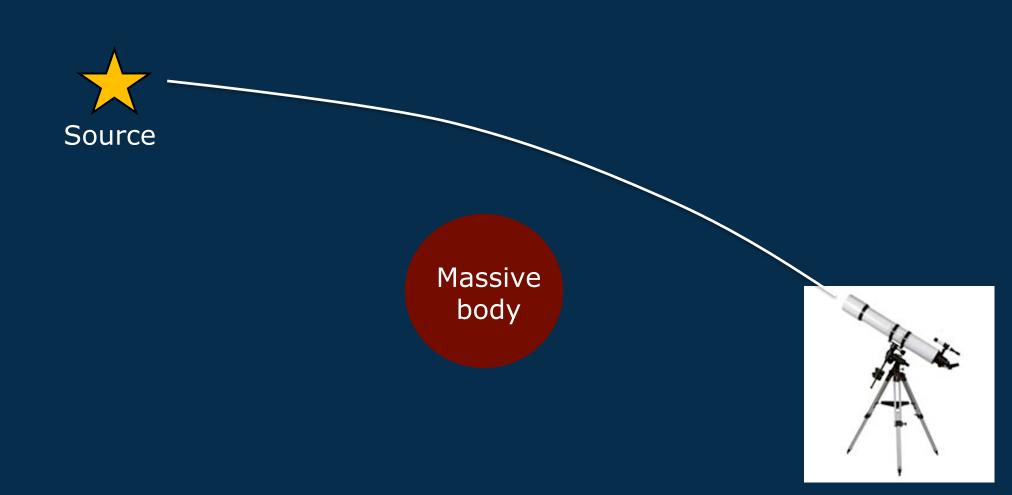


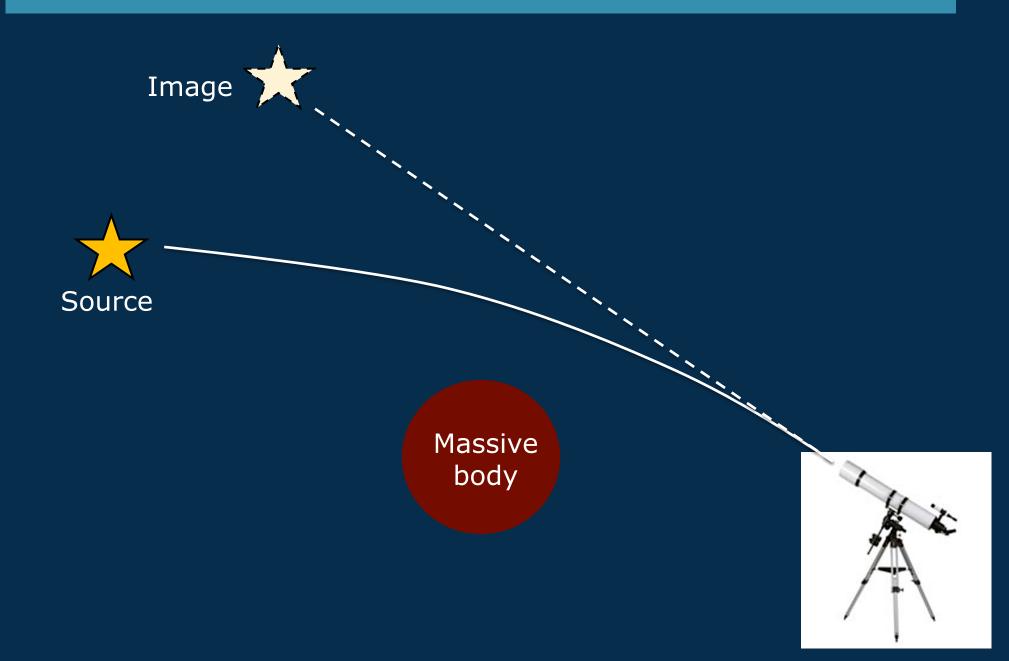


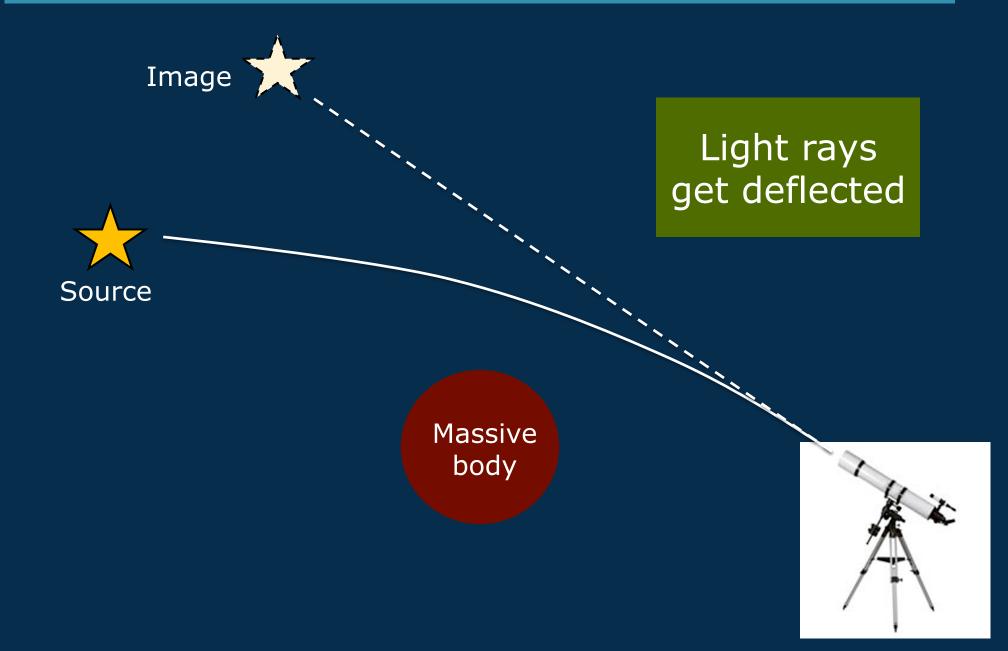




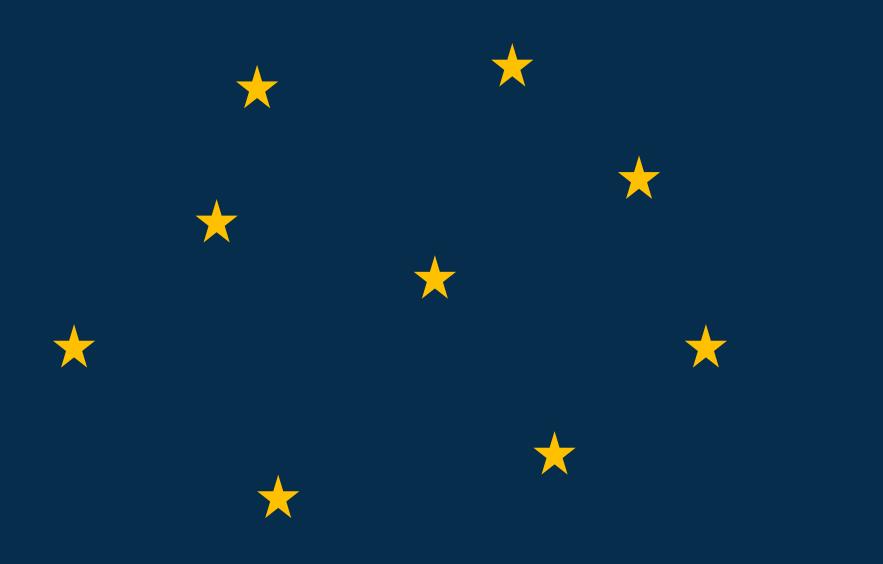




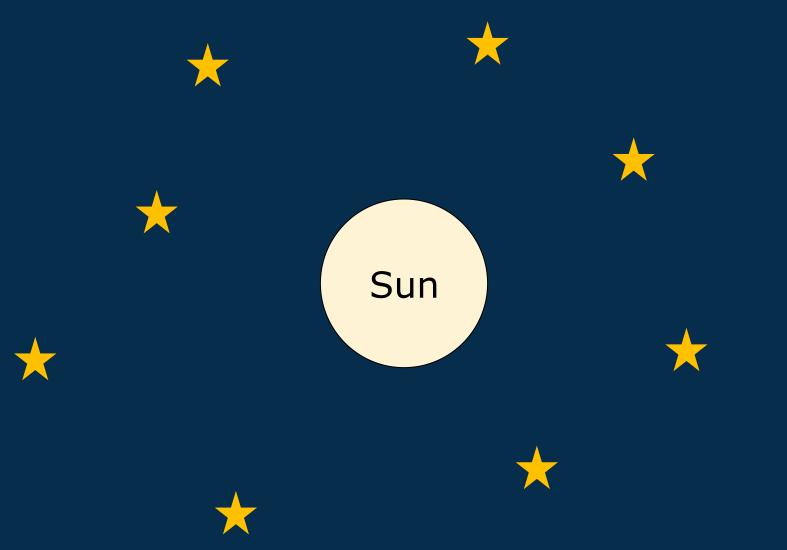




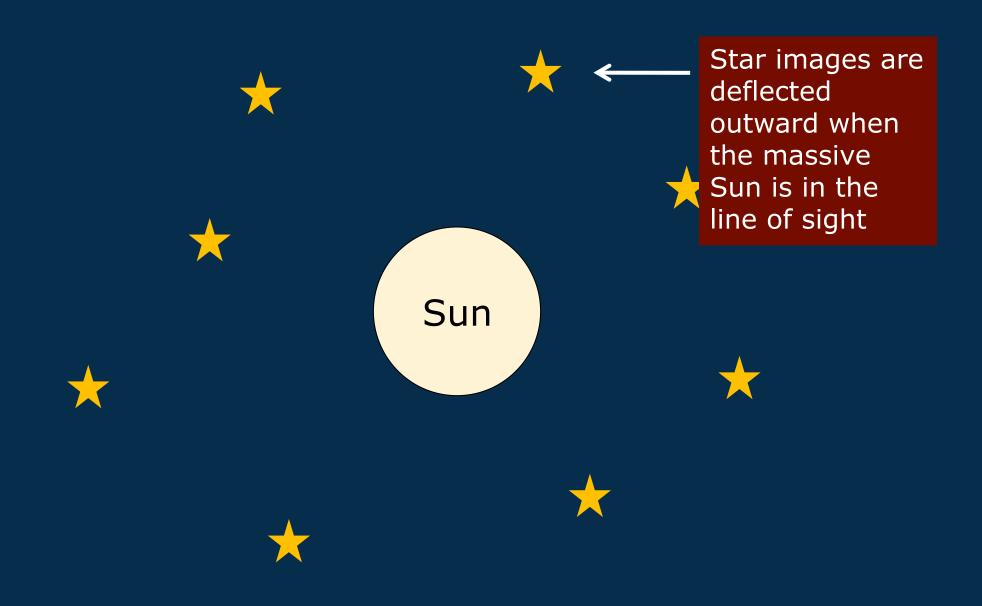
The Eddington experiment



The Eddington experiment

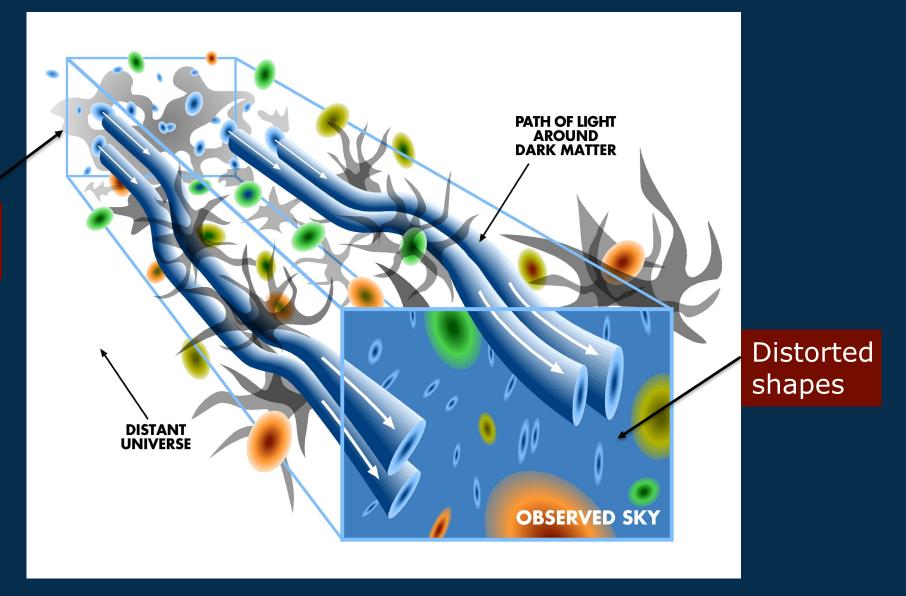


The Eddington experiment



Lensing changes image shapes

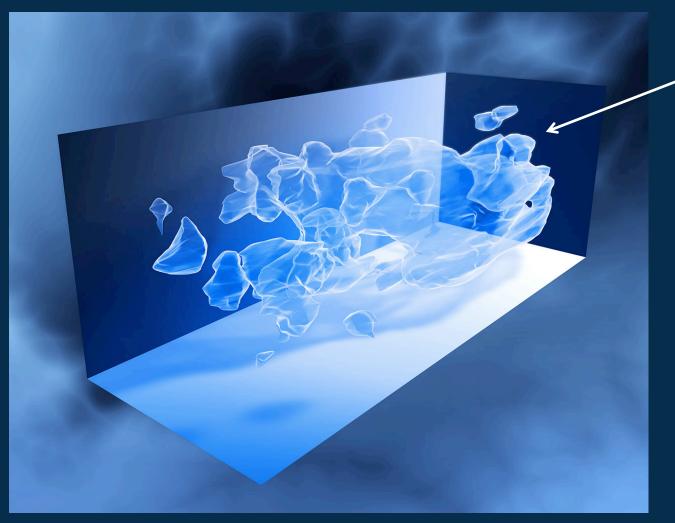
Intrinsic shapes



http://www.lsst.org

Why is this useful?

Dark matter mapping:



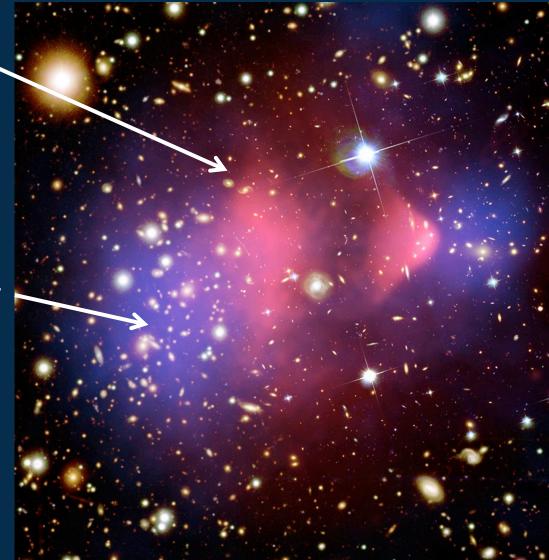
We think that about 85% of the matter in the Universe is dark

Massey et al. 2004

The bullet cluster – evidence for dark matter

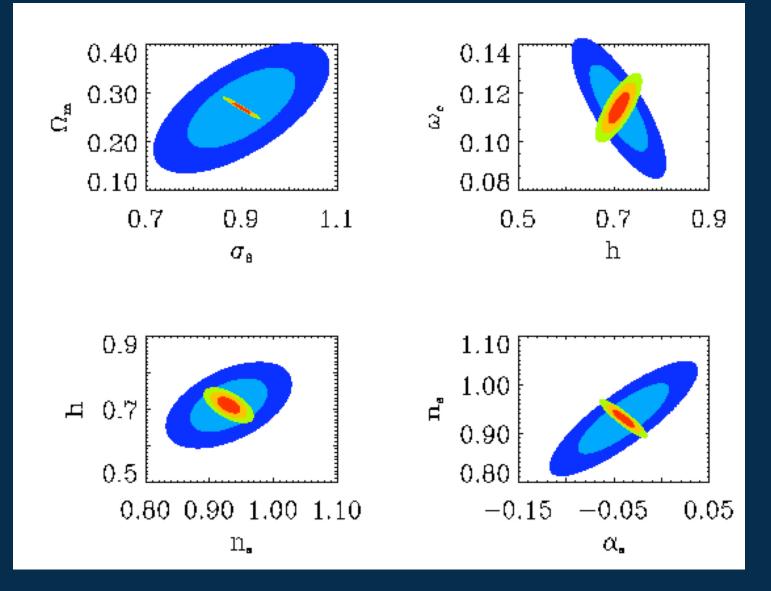
Pink: hot gas

Blue: total matter



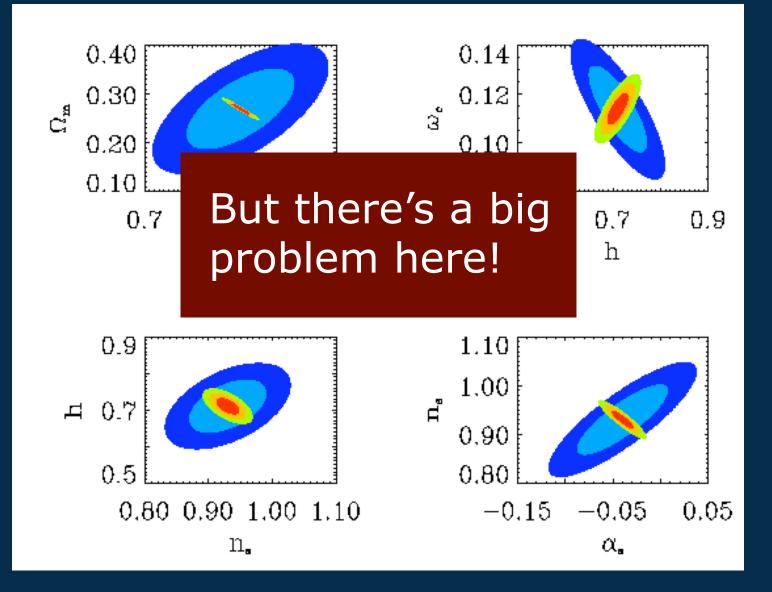
http://ccapp.osu.edu

Lensing constrains other cosmological parameters



Tereno et al. 2005

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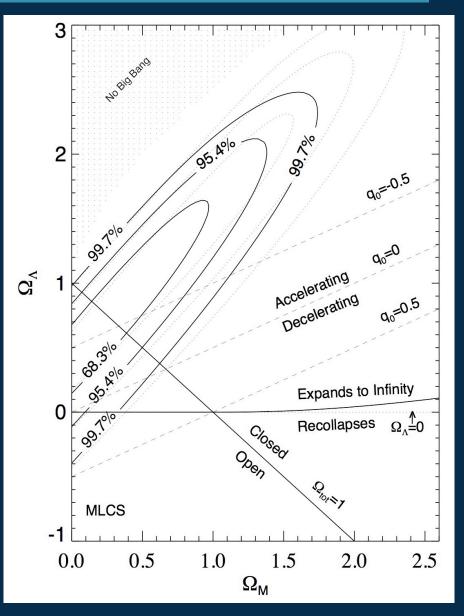
The dark energy problem

Supernovae look dimmer than expected in a "normal" universe

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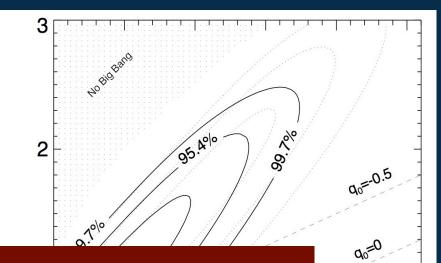
But we can fit the data if we add in some "dark energy"



Riess et al. 1998

The dark energy problem

Supernovae look dimmer than expected in a "normal" universe



q0=0.5

2.5

2.0

But ng **Problem:** the expansion data of the Universe appears to som nds to Infinity lapses Ω,[≜]0 be accelerating

0.0

0.5

Riess et al. 1998

1.0

1.5

 $\Omega_{\rm M}$

 Measurement error – maybe we did the measurements or data analysis wrong? (Unlikely)

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Modifying gravity

Imagine the Universe is a huge ball of matter

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• Large scales – acceleration is incompatible with GR

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Imagine the Universe is a huge ball of matter

- Sma work
 How do we measure these modifications?
- Large incompatible with GR



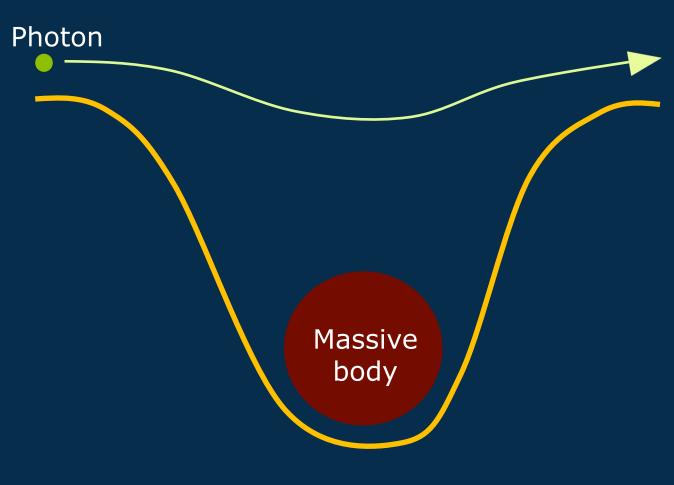
Modify gravity theory

Imagine the Universe is a huge ball of matter

Sma work
How do we measure these modifications? One way: One way: Gravitational lensing
Large
Incompatible with GR

Modify gravity theory

Lensing in GR



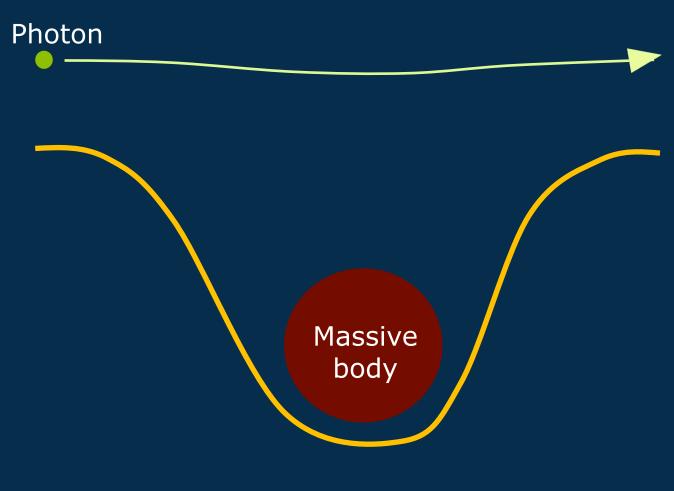
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Lensing in modified gravity



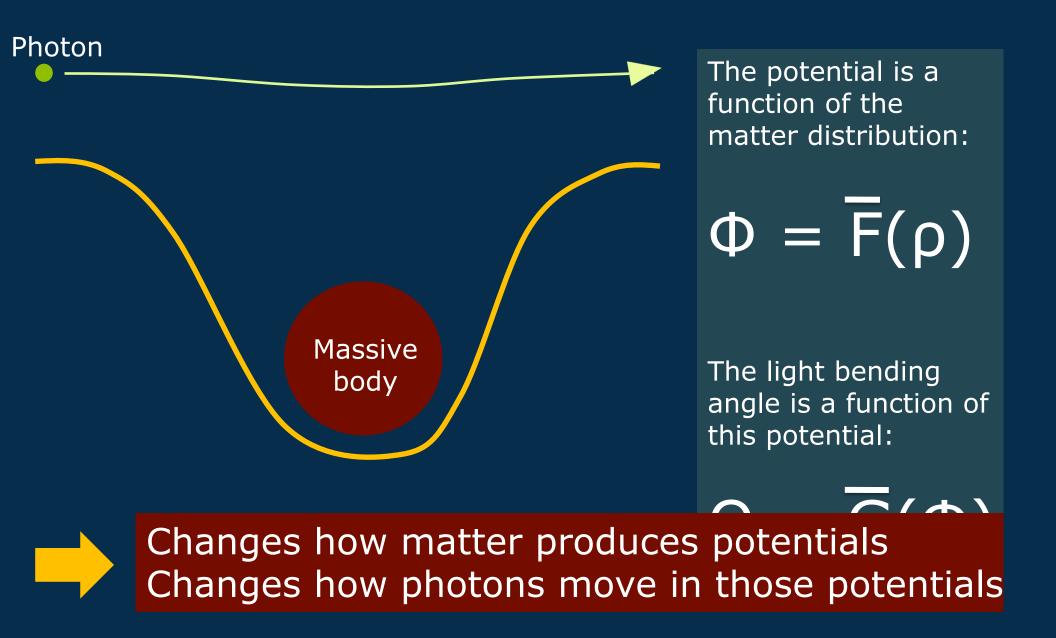
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Lensing in modified gravity



Big question: How do we best use gravitational lensing observations to constrain modifications of GR on very large scales?

Lessons from "small" scales

The parameterized post-Newtonian formalism –

the gravitational potentials of GR are modified, for instance:

Matter distribution $\rho(x)$

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Photons "feel" the gravitational potential γΦ(x)

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Photons "feel" the gravitational potential γΦ(x)



Parameter	What it measures relative to GR	Value in GR	Value in semi- conservative theories	Value in fully conservative theories
γ	How much space-curva- ture produced by unit rest mass?	1	γ	γ
β	How much "nonlinearity" in the superposition law for gravity?	1	β	eta
ξ	Preferred-location effects?	0	ξ	ξ
α_1	Preferred-frame effects?	0	α_1	0
α_2		0	α_2	0
α_3		0	0	0
α_3	Violation of conservation	0	0	0
ζ_1	of total momentum?	0	0	0
ζ_2		0	0	0
ζ_3		0	0	0
ζ_4		0	0	0

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$lpha_1$ $lpha_2$ $lpha_3$	Preferred-frame effects?	dev	iations from	oarameterizin GR in "weak	
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ζ_2					
ζ_3					
× .					

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$\begin{array}{c} \alpha_1 \\ \alpha_2 \\ \alpha_3 \end{array}$	Preferred-frame effects?	dev	Framework for parameterizing deviations from GR in "weak field" gravity:		
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ζ_2		(1)		/stem	
Č.					
ζ_3					

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ζ_3 ζ_4		(2)	Large scale	cosmology

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δθ

d

Μ

Light deflection due to the Sun:

 $\delta\theta = \frac{1}{2}(1+\gamma)\frac{4M}{d}$

δθ

d

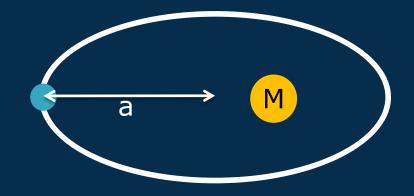
Μ

Light deflection due to the Sun:

$$\delta\theta = \frac{1}{2}(1+\gamma)\frac{4M}{d}$$

$\gamma - 1 = (-1.7 \pm 4.5) \times 10^{-4} (VLBI)$

Perihelion precession of Mercury:



Advance per orbit =

$$\Delta\phi = \frac{2\pi M}{a(1-e^2)}(2+2\gamma-\beta)+\dots$$

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a M

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|2γ-β-1|<3×10⁻³ (Shapiro, 1990)

Cosmological constraints?

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One way – weak gravitational lensing

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Weak lensing data

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Constrain PPN

parameters

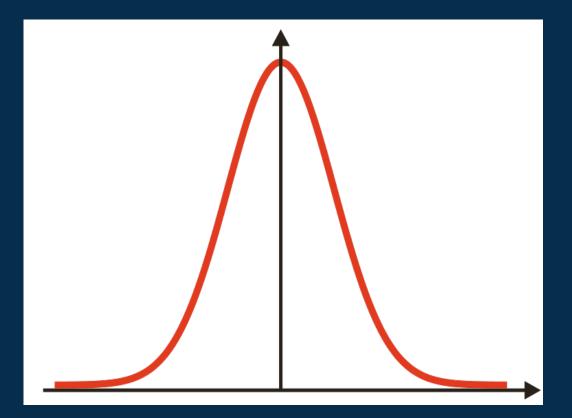
Weak lensing observables in modified gravity

Through the SURP program, we are working with Prof. Robert Caldwell at Dartmouth to compute:

- Light bending and image distortions
- Weak lensing power spectrum
- Higher-order statistics, such as the bispectrum

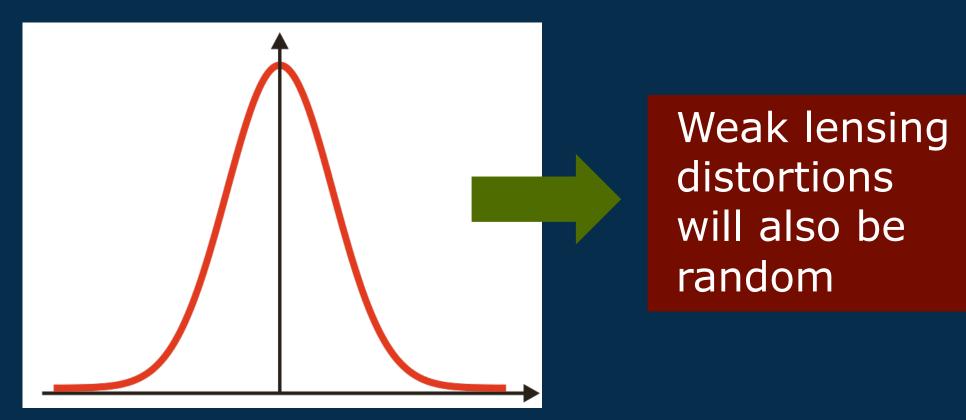
Example: Apparent non-Gaussianity

If density fluctuations are very small, they can be approximated as a random field with a Gaussian distribution:



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The search for non-Gaussianity

Ways for these distributions to the non-Gaussian:

Exotic new physics

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 Nonlinearity – density fluctuations are actually nonlinear and therefore correlated

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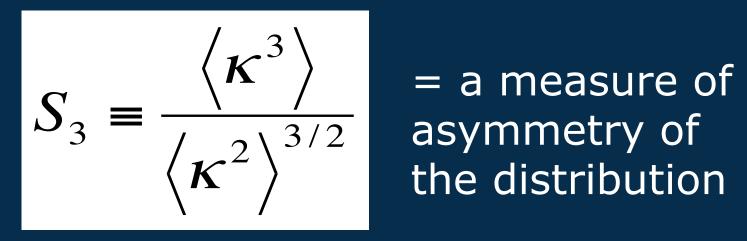
 Nonlinearity – density fluctuations are actually nonlinear and therefore correlated



Modified gravity theories have different levels of nonlinearity via the PPN parameter β

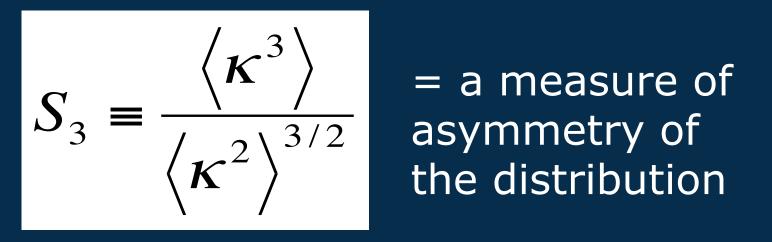
Non-Gaussianity from modified gravity

Gaussian distributions have zero skewness



Non-Gaussianity from modified gravity

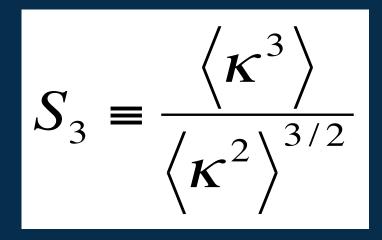
Gaussian distributions have zero skewness



 $= f(\beta)$

Non-Gaussianity from modified gravity

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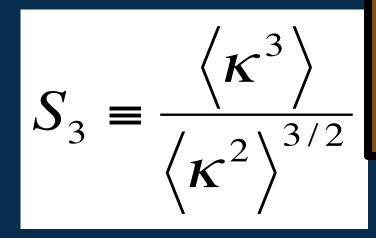
= a measure of $= f(\beta)$ asymmetry of the distribution

Measure the weak lensing skewness



Non-Gaussianity from modified gravity

Gaussian distr skewness



Get one step closer to finding the cause of cosmic acceleration

Constrain B

the distributi

Measure the weak lensing skewness



Another big question: How do we obtain these weak lensing observations?

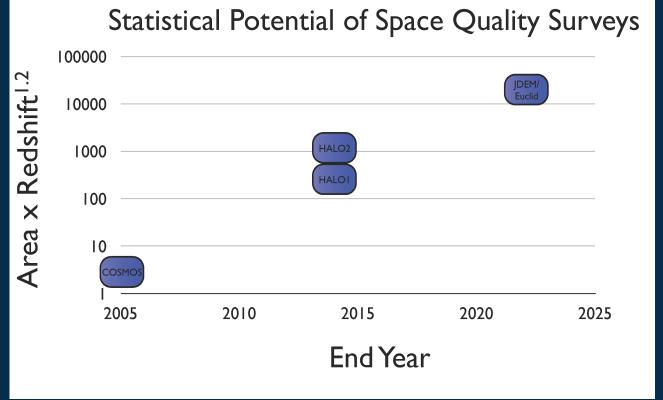
The High Altitude Lensing Observatory



PI: Jason Rhodes

Jeff Booth (JPL), Kurt Liewer (JPL), Michael Seiffert (JPL), Wesley Traub (JPL), Richard Key (JPL), Adam Amara (ETH Zurich), Richard Ellis (Caltech), Richard Massey (University of Edinburgh), Satoshi Miyazaki (NOAJ Japan), Harry Teplitz (Spitzer Science Center, Caltech), Calvin Barth Netterfield (University of Toronto), Alexandre Refregier (CEA Saclay, Paris), Roger Smith (Caltech)

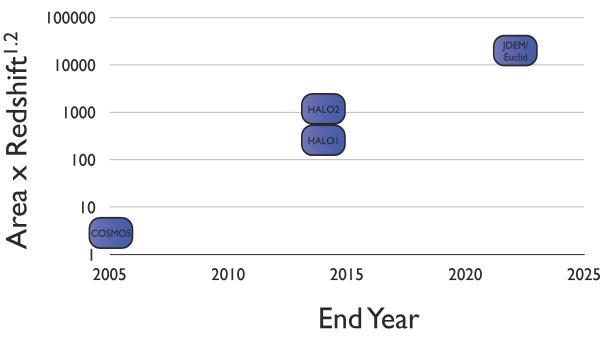
Space quality weak lensing past and future



Past: COSMOS survey

Space quality weak lensing past and future

Statistical Potential of Space Quality Surveys

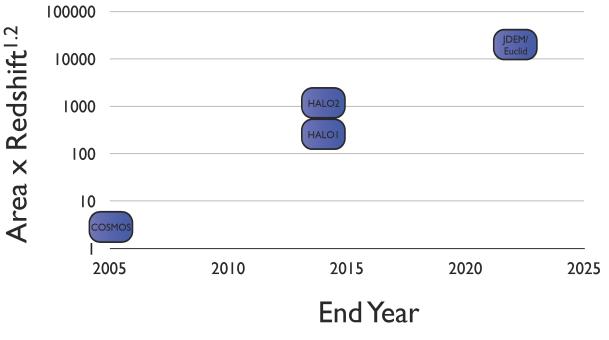


Past: COSMOS survey

Future: the Joint Dark Energy Mission (JDEM) and ESA's Euclid

Space quality weak lensing past and future

Statistical Potential of Space Quality Surveys



<u>**Past</u>: COSMOS** survey</u>

Future: the Joint Dark Energy Mission (JDEM) and ESA's Euclid

<u>Now</u>: the balloon-based High Altitude Lensing Observatory (HALO) fills the gap

The Dark Energy Task Force (Albrecht et al. 2006) Weak lensing is the best method for constraining dark energy and modified gravity, provided numerous systematic effects can be overcome

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Ground-based surveys are fundamentally limited by systematics

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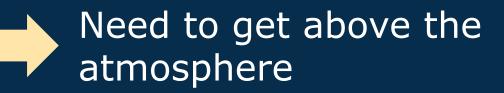
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But JDEM and Euclid are at least a decade away

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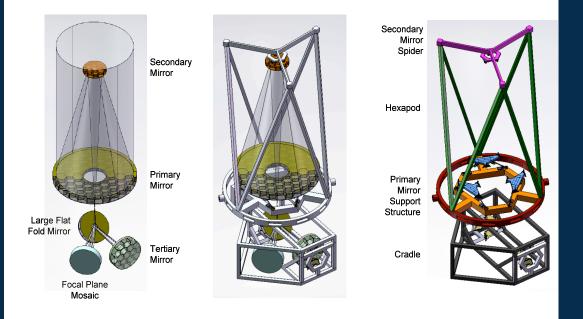


Use a balloon!



HALO

- 15-20 day flight Australia to Australia (can stop in South America if needed)
- 1.2m lightweight primary mirror
- 48 2k×4k Hamamatsu CCDs
- Single wide optical filter
- Solar panel to recharge batteries



 Need to pick up the disk drives (2 Tb) afterwards to do the science

Survey strategy

- 500-720 nm filter
- 1500 second integration time
- 13 square degrees per night
- 15-20 galaxies per square arcminute
- Redshifts from the ground



Understand dark matter:

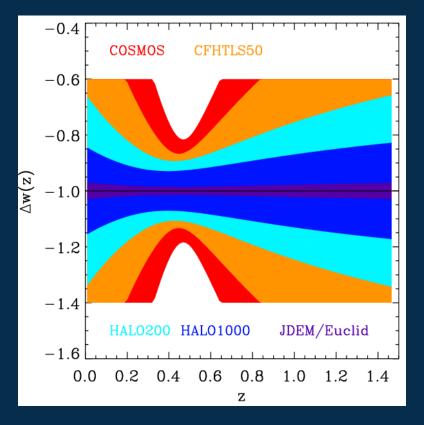
- Amount and distribution
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Understand dark matter:

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Explore dark energy and modified gravity:

- Examine expansion history
- Growth of structure



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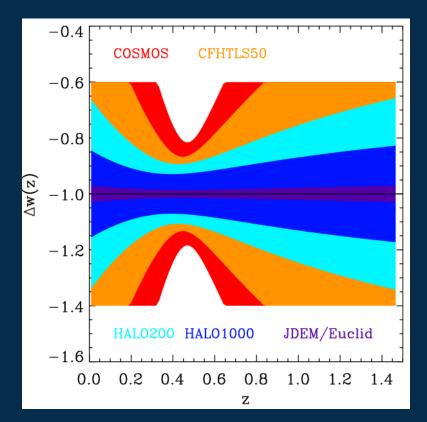
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Ancillary science:

- Galaxy morphology and evolution
- Stellar counts
- Surface brightness fluctuations



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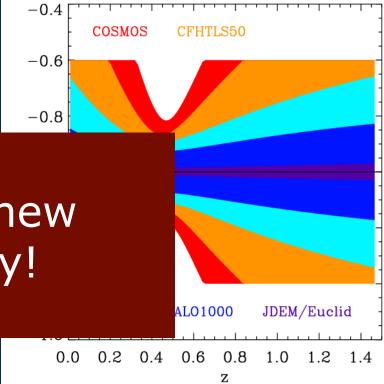
Explore darl modified gra

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An exciting new opportunity!

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Conclusions

- Modifications of GR on cosmological scales could explain cosmic acceleration
- Weak lensing is an excellent probe of modified gravity
- The High Altitude Lensing Observatory (HALO) will produce the high-quality weak lensing data needed to carry this out in the near term