Testing gravity on cosmological scales with weak lensing

JPL Postdoc Seminar Series August 13, 2009 Ali Vanderveld (JPL, Caltech)

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Abell 2218

(1) Lensing via refraction

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Photons gravitate towards the massive object

Massive body

The gravitational potential $\Phi(x)$

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The light bending angle is a function of this potential:

 $\Theta = G(\Phi)$

Massive body

The Eddington experiment

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Lensing changes image shapes

Intrinsic shapes

http://www.lsst.org

Why is this useful?

Dark matter mapping:

We think that about 85% of the matter in the Universe is dark

Massey et al. 2004

The bullet cluster – evidence for dark matter

Pink: hot gas

Blue: total matter

http://ccapp.osu.edu

Lensing constrains other cosmological parameters

Tereno et al. 2005

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The dark energy problem

Supernovae look dimmer than expected in a "normal" universe

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But we can fit the data if we add in some "dark energy"

Riess et al. 1998

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 $\Omega_{\rm M}$

Riess et al. 1998

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Modifying gravity

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Modify gravity theory

Imagine the Universe is a huge ball of matter

• Small scales – we found that GR work livw uu w • Large incompatible with GR One way: state ged How do we measure these modifications? Gravitational lensing

Modify gravity theory

Lensing in GR

The potential is a function of the matter distribution:

 $\Phi = F(\rho)$

The light bending angle is a function of this potential:

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Lensing in modified gravity

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 $\Phi = \overline{F}(\rho)$

The light bending angle is a function of this potential:

 $\Theta = \overline{G}(\Phi)$

Lensing in modified gravity

Big question: How do we best use gravitational lensing observations to constrain modifications of GR on very large scales?

Lessons from "small" scales

The parameterized post-Newtonian formalism –

the gravitational potentials of GR are modified, for instance:

> Matter distribution $p(x)$

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> **Matter** distribution $\rho(x)$

Photons "feel" the gravitational potential $\overline{\mathsf{y}\Phi(\mathsf{x})}$

Lessons from "small" scales

The parameterized post-Newtonian formalism –

the gravitational potentials of GR are modified, for instance:

Matter distribution $p(x)$ Photons "feel" the gravitational potential $\gamma \Phi(x)$

Can constrain the parameter γ

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Better: Constrain just the PPN parameters with observations

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Better: Constrain just the PPN parameters with observations

M

d

δθ

Light deflection due to the Sun:

 $\delta\theta =$ 1 2 $(1 + \gamma)$ 4*M d*

M

d

δθ

Light deflection due to the Sun:

$$
\delta\theta = \frac{1}{2}(1+\gamma)\frac{4M}{d}
$$

$\sqrt{y-1}=(-1.7\pm4.5)\times10^{-4}$ (VLBI)

Perihelion precession of Mercury:

Advance per orbit =

$$
\Delta \phi = \frac{2\pi M}{a(1-e^2)} (2+2\gamma-\beta) + \dots
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|2γ-β-1|<3×10-3 (Shapiro, 1990)

a → M

Cosmological constraints?

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Weak lensing data

Constrain PPN parameters

Weak lensing observables in modified gravity

Through the SURP program, we are working with Prof. Robert Caldwell at Dartmouth to compute:

- Light bending and image distortions
- Weak lensing power spectrum
- Higher-order statistics, such as the bispectrum

Example: Apparent non-**Gaussianity**

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The search for non-Gaussianity

Ways for these distributions to the non-Gaussian:

• Exotic new physics

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Modified gravity theories have different levels of nonlinearity via the PPN parameter β
Gaussian distributions have zero skewness

= a measure of asymmetry of the distribution

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Measure the weak lensing skewness

Gaussian distr skewness

cosmic acceleration Get one step closer to finding the cause of cosmic acceleration

Constrain β

the distribution

Measure the weak lensing skewness

Another big question: How do we obtain these weak lensing observations?

The High Altitude Lensing **Observatory**

PI: Jason Rhodes

Jeff Booth (JPL), Kurt Liewer (JPL), Michael Seiffert (JPL),Wesley Traub (JPL), Richard Key (JPL), Adam Amara (ETH Zurich), Richard Ellis (Caltech), Richard Massey (University of Edinburgh), Satoshi Miyazaki (NOAJ Japan), Harry Teplitz (Spitzer Science Center, Caltech), Calvin Barth Netterfield (University of Toronto), Alexandre Refregier (CEA Saclay, Paris), Roger Smith (Caltech)

Space quality weak lensing past and future

Past: COSMOS survey

Space quality weak lensing past and future

Statistical Potential of Space Quality Surveys

Past: COSMOS survey

Future: the Joint Dark Energy Mission (JDEM) and ESA's Euclid

Space quality weak lensing past and future

Statistical Potential of Space Quality Surveys 100000

Area x Redshift^{1.2}

Past: COSMOS survey

Future: the Joint Dark Energy Mission (JDEM) and ESA's Euclid

Now: the balloon-based High Altitude Lensing Observatory (HALO) fills the gap

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Need to get above the atmosphere

But JDEM and Euclid are at least a decade away

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Use a balloon!

HALO

- 15-20 day flight Australia to Australia (can stop in South America if needed)
- 1.2m lightweight primary mirror
- 48 2k×4k Hamamatsu CCDs
- Single wide optical filter
- Solar panel to recharge batteries

• Need to pick up the disk drives (2 Tb) afterwards to do the science

Survey strategy

- 500-720 nm filter
- 1500 second integration time
- 13 square degrees per night
- 15-20 galaxies per square arcminute
- Redshifts from the ground

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- Amount and distribution
- Weak and strong lensing

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An exciting new opportunity!

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Conclusions

- Modifications of GR on cosmological scales could explain cosmic acceleration
- Weak lensing is an excellent probe of modified gravity
- The High Altitude Lensing Observatory (HALO) will produce the high-quality weak lensing data needed to carry this out in the near term