

# $e+A$ physics at a future Electron-Ion Collider

## Outline

- What is the nature of glue?
- Non-linear QCD and Saturation physics
  - Nuclear “Oomph” Factor
- The 4 key gluon measurements in  $e+A$ 
  - Momentum distributions
  - Space-time distributions
  - The role of the Pomeron
  - Energy loss in cold nuclear matter

Matthew A. C. Lamont  
Brookhaven National Lab

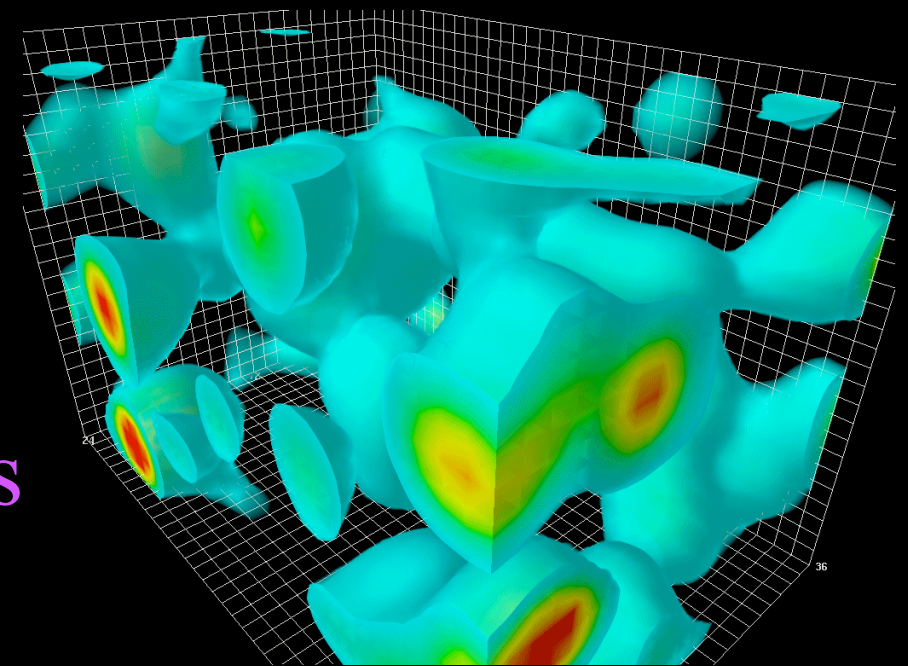
# What do we know about gluons?

## Glue and the QCD Lagrangian:

$$L_{QCD} = \bar{q}(i\gamma^\mu \partial_\mu - m)q - g(\bar{q}\gamma^\mu T_a q)A_\mu^a - \frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu}$$

- Despite the success of QCD, our knowledge of glue is very limited. We know:

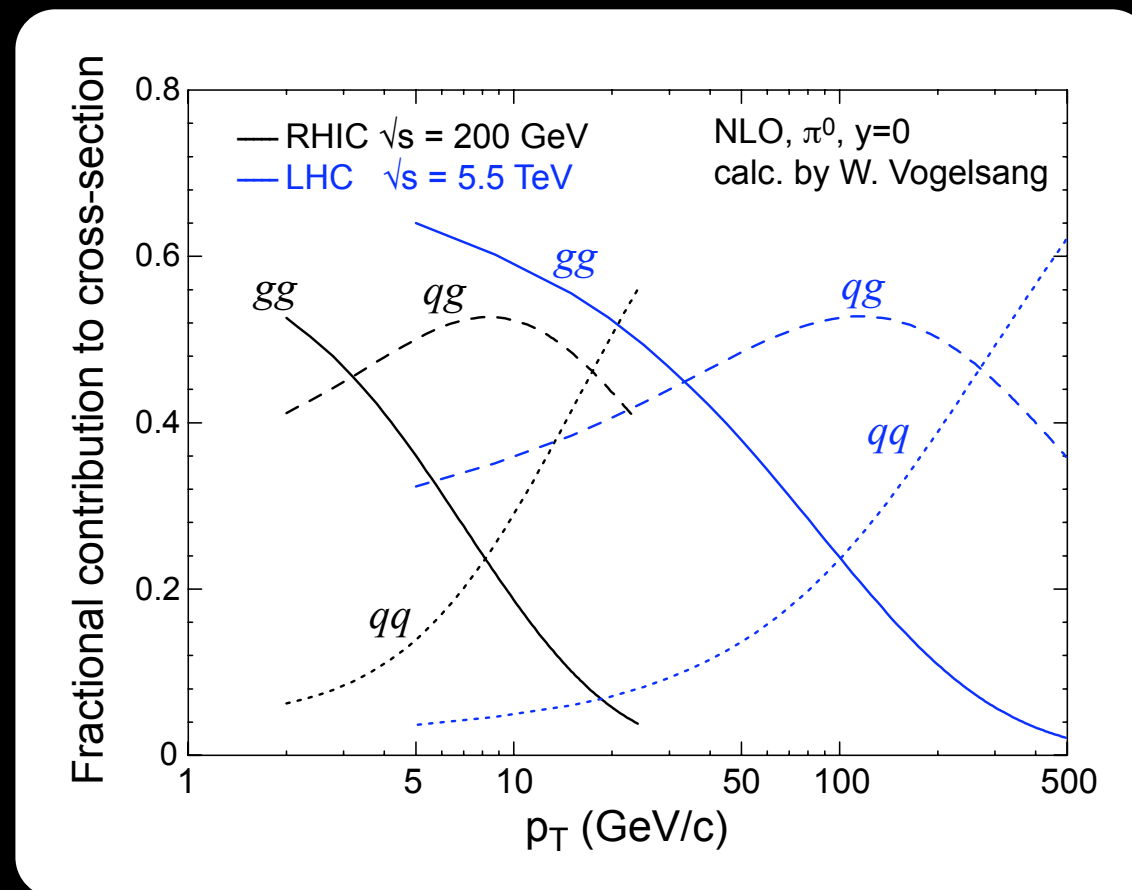
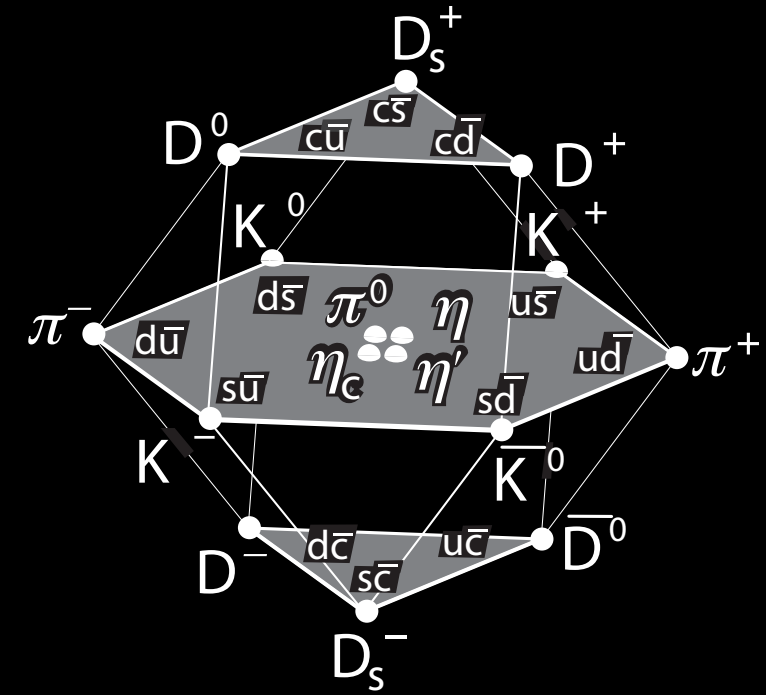
- Mediators of the strong interaction
- Determine essential features of QCD
- Asymptotic freedom from gluon loops
- Dominate structure of QCD vacuum ( $\chi$ SB)
  - Quenched  $L_{QCD}$  gets hadron masses correct to  $\sim 10\%$



Action ( $\sim$ energy) density fluctuations of gluon-fields in QCD vacuum ( $2.4 \times 2.4 \times 3.6$  fm) (Derek Leinweber)

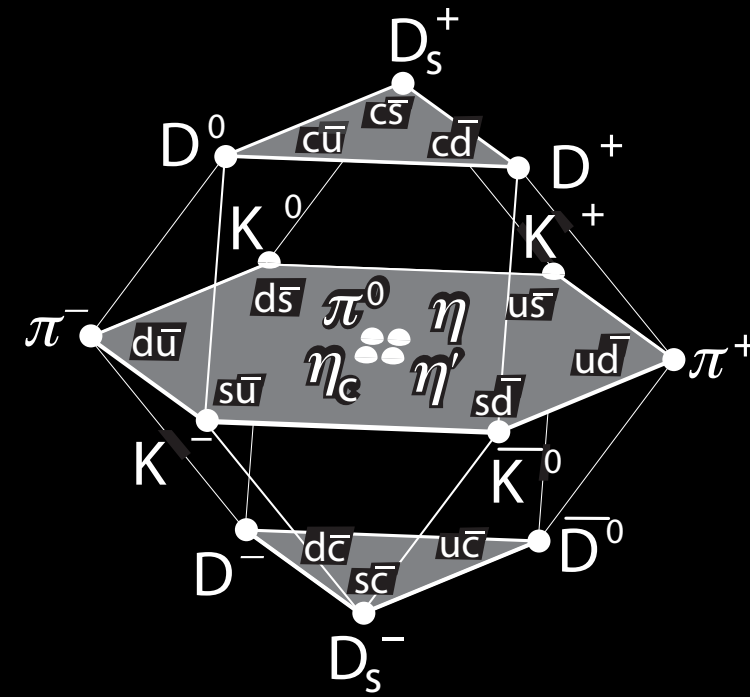
# Glue and the Lagrangian

- **Hard to “see” glue in the low-energy world**
  - ➔ Gluon degrees of freedom “missing” in hadronic spectrum
  - ➔ Constituent Quark Picture?
- **From DIS:**
  - ➔ Drive the structure of baryonic matter already at medium-x
- **Crucial players at RHIC and the LHC**
  - ➔ Drive the entropy

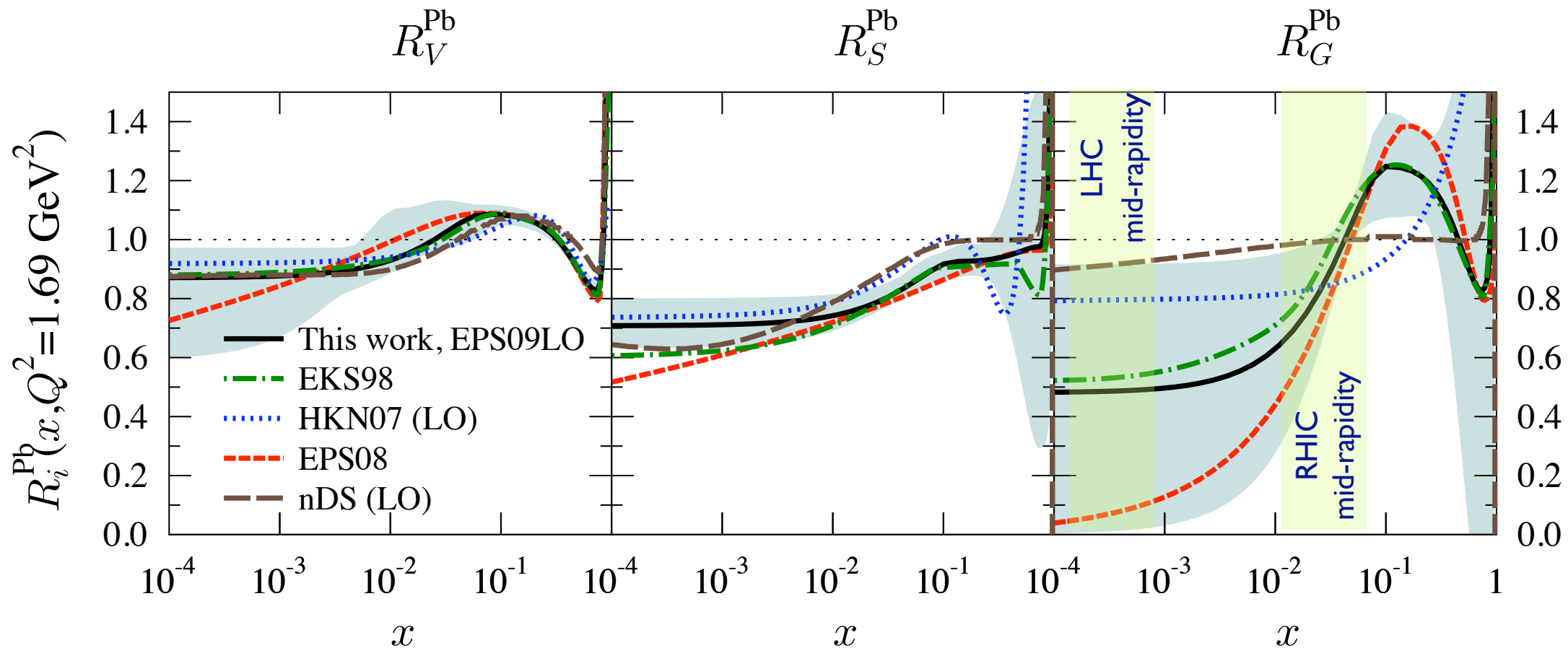


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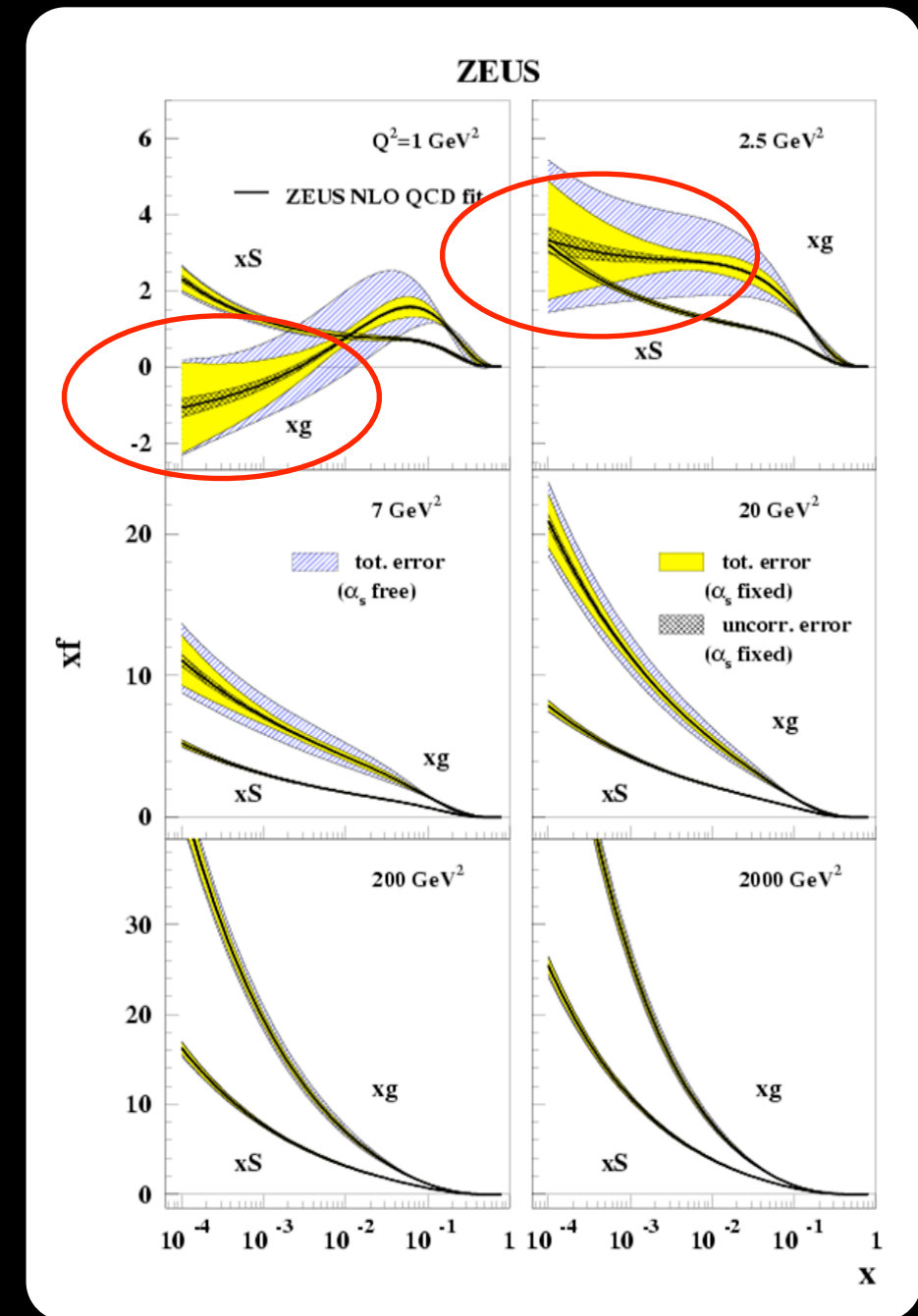
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Eskola, Paukkonen, Salgado: arXiv0902.4154

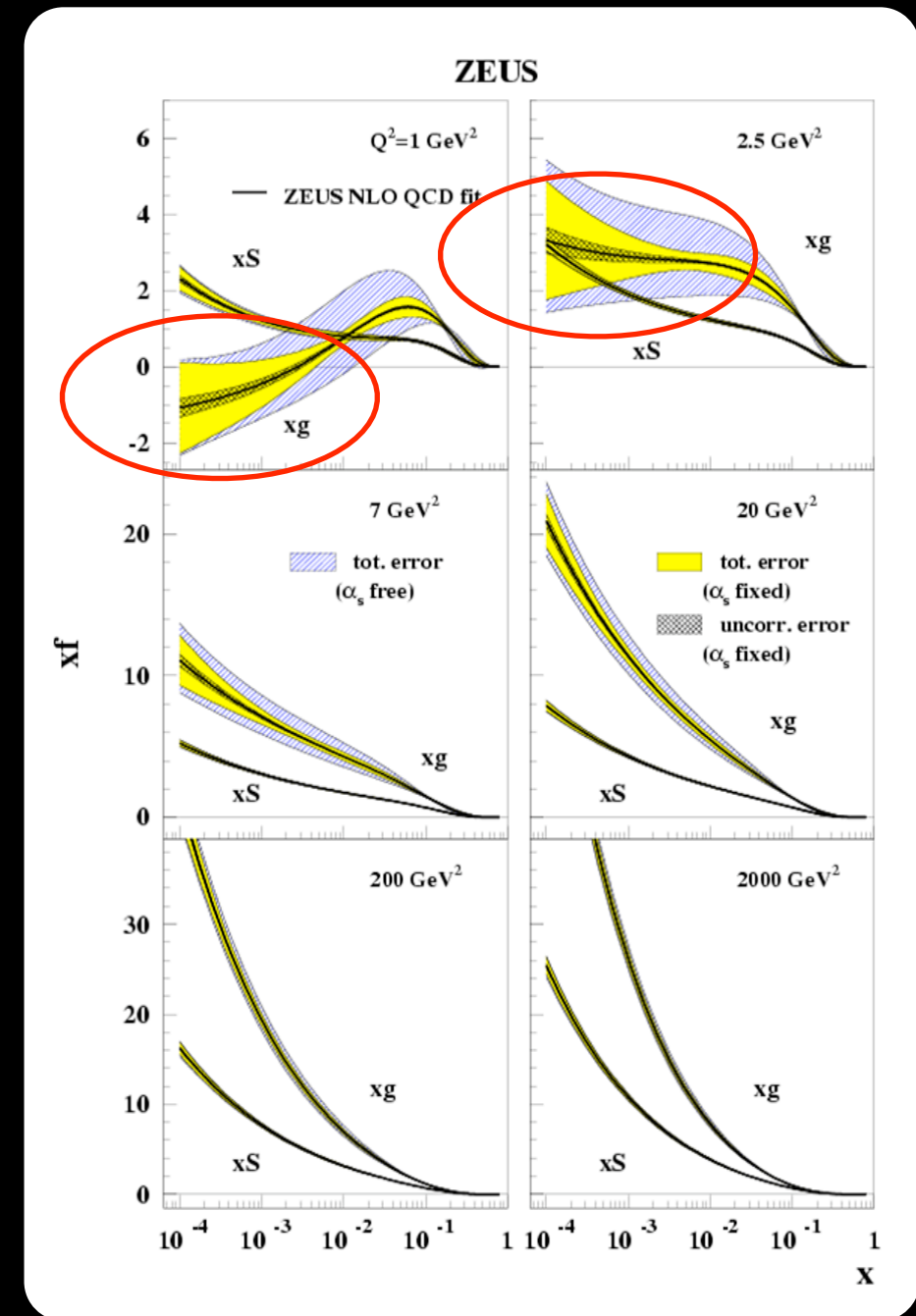


# The problem with our current understanding



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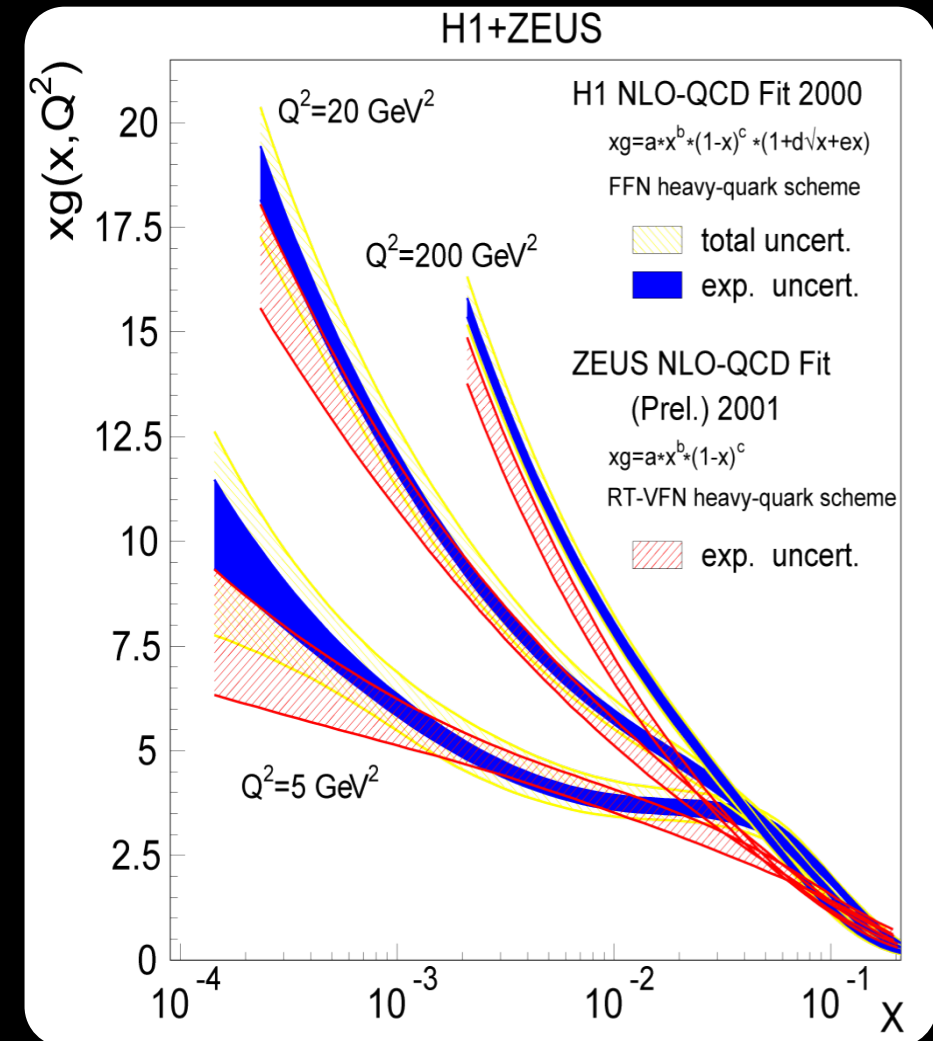
- Using the **Linear DGLAP** evolution model:
- Weird behaviour of  $xG$  at low- $x$  and low- $Q^2$  in HERA data
  - $xG$  goes negative
  - $xG < xS$  (even though sea quarks come from gluon splitting)





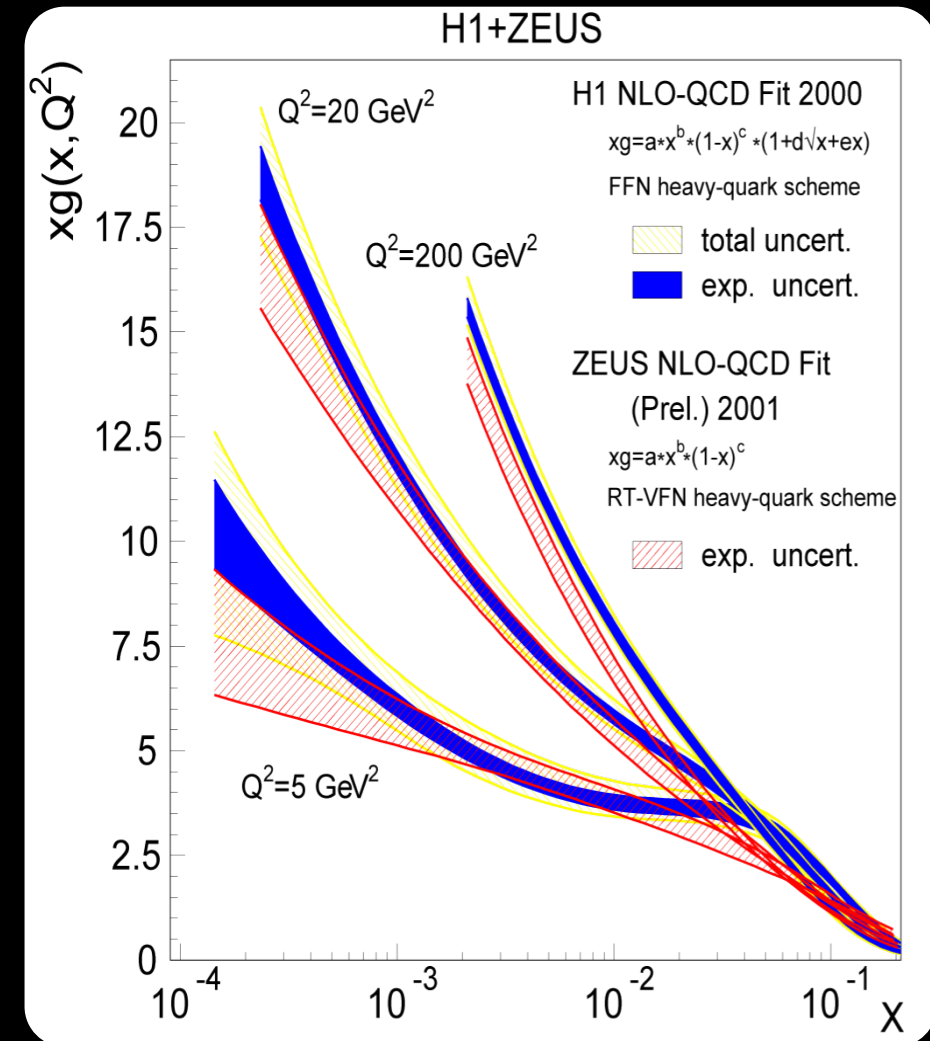
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- More Severe:
  - Linear evolution has a built-in high energy “catastrophe”
  - $xG$  has a rapid rise with decreasing  $x$  (and increasing  $Q^2$ )  $\Rightarrow$  violation of Froissart unitarity bound
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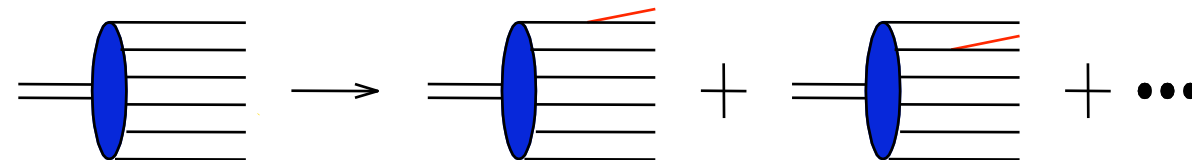
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    - Must have saturation !!
- **What is the underlying dynamics?**





# Non-linear QCD and Saturation

proton

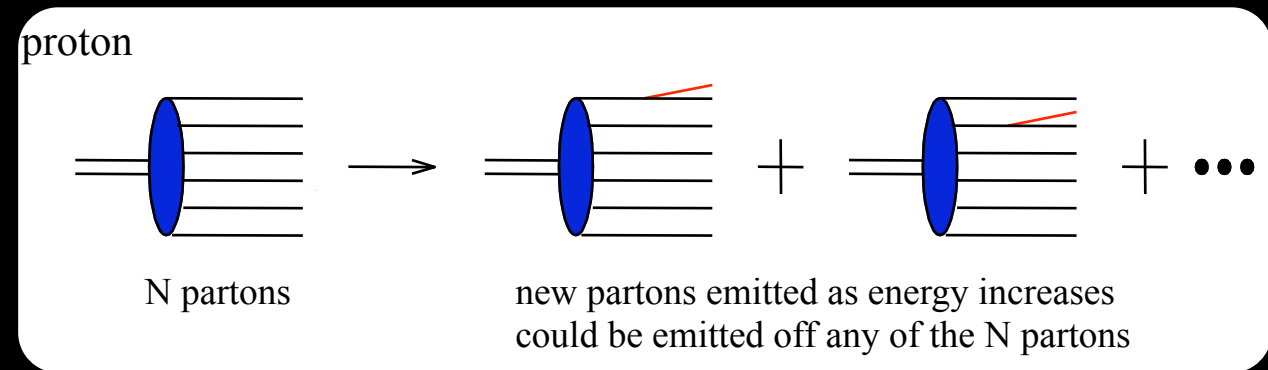


N partons

new partons emitted as energy increases  
could be emitted off any of the N partons

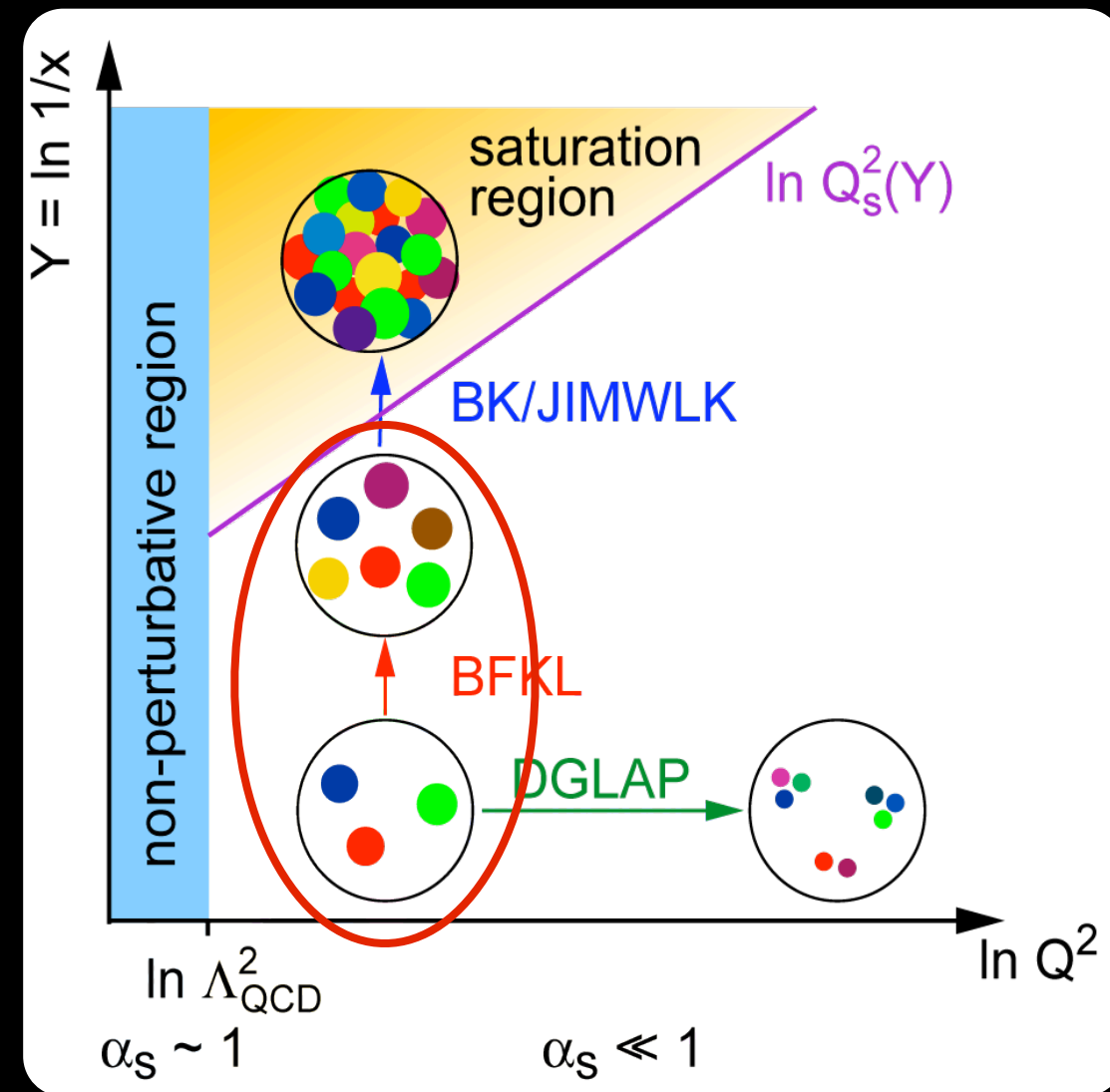
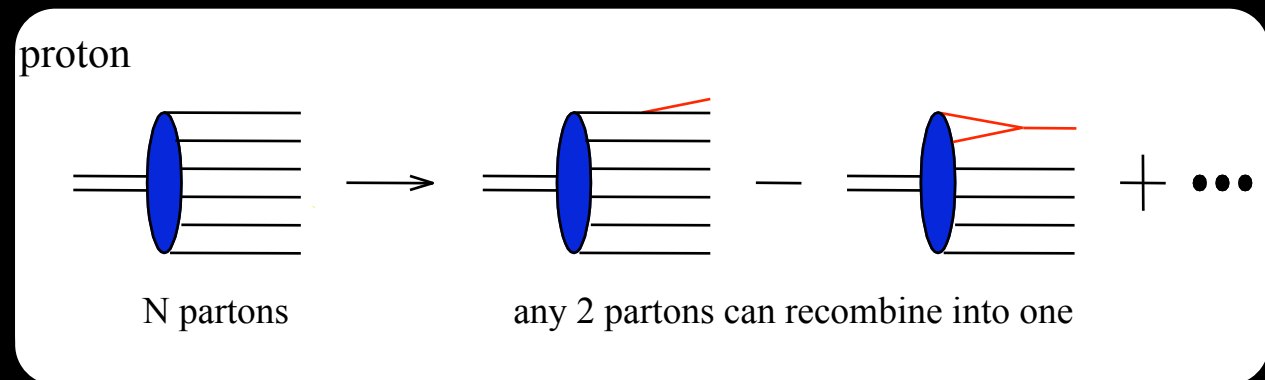
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  - linear
  - explosion in colour field at small- $x$



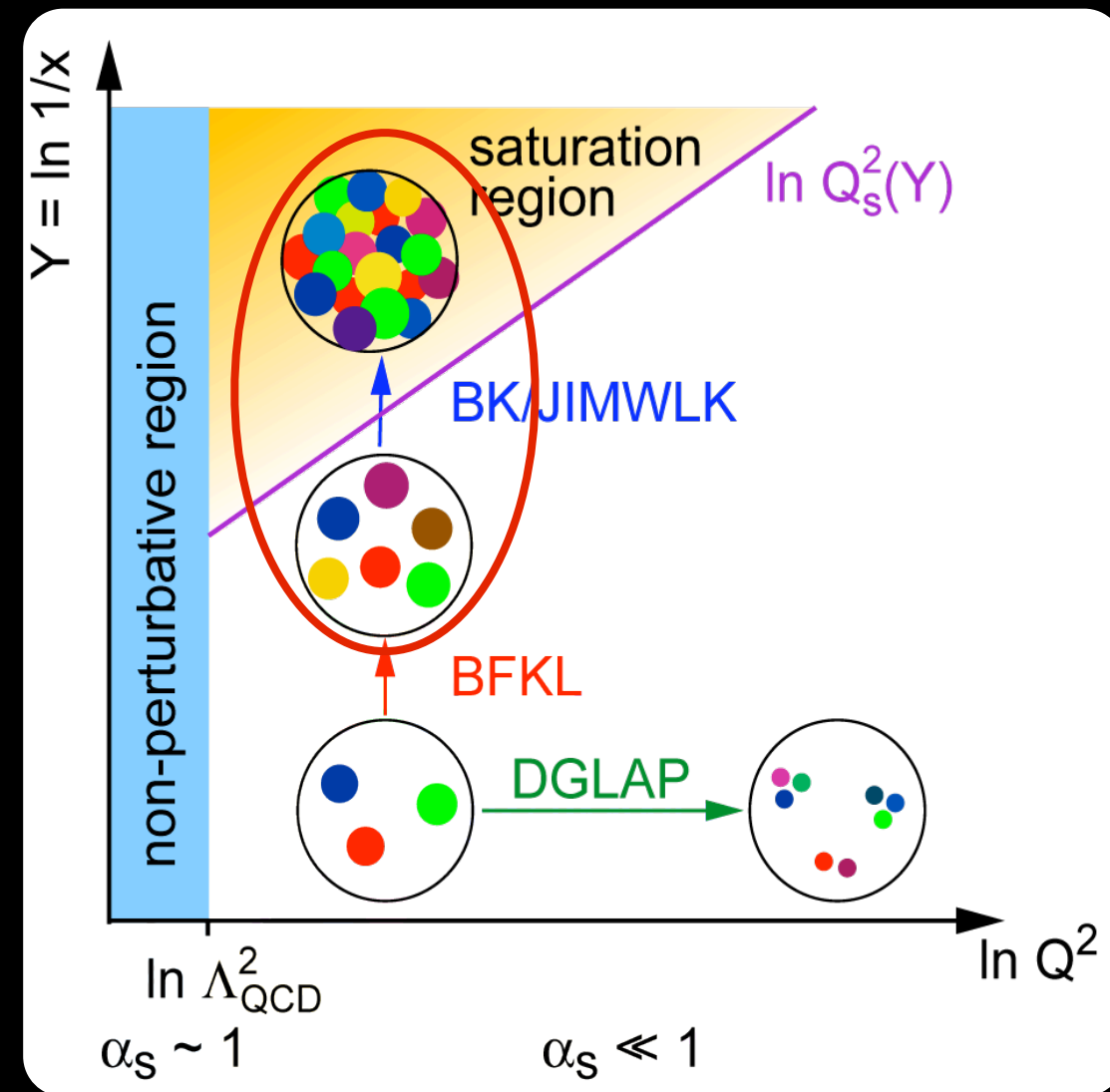
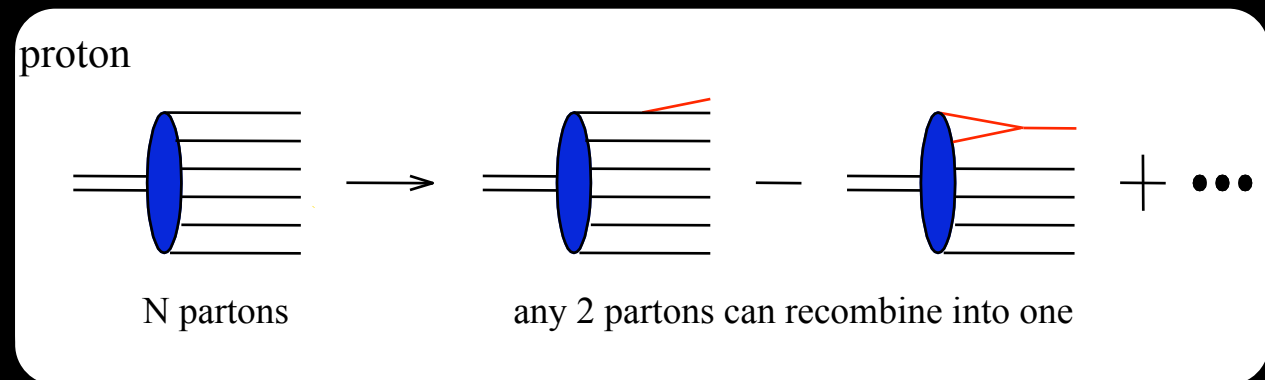
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- characterised by the saturation scale,  $Q_s(x,A)$
- arises naturally in the Colour-Glass Condensate EFT

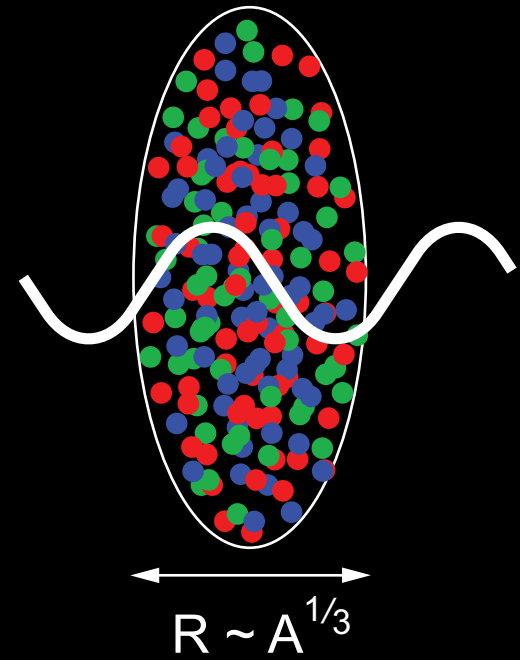


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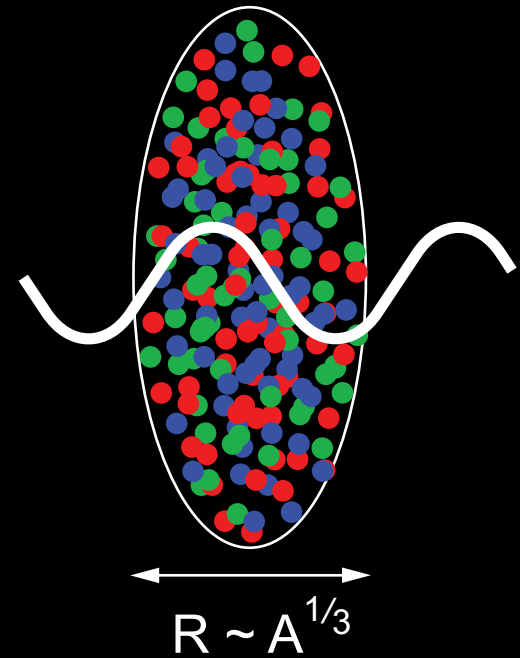


# The Nuclear “Oomph Factor”



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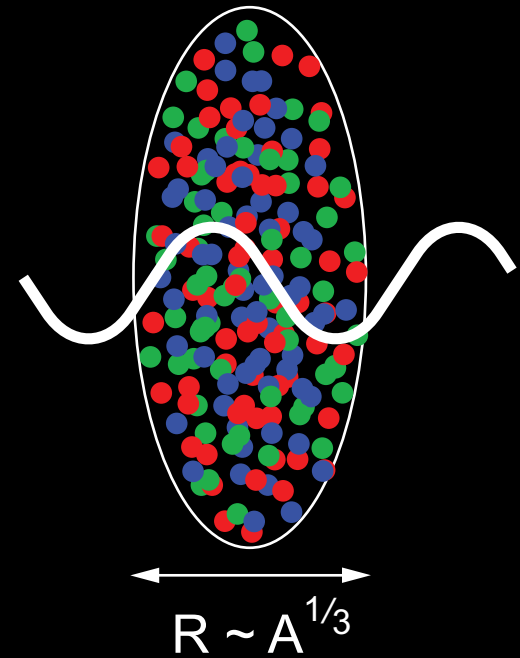
- Enhancing Saturation effects:
  - Probes interact over distances  $L \sim (2mnx)^{-1}$
  - For probes where  $L > 2R_A (\sim A^{1/3})$ , cannot distinguish between nucleons in front or back of the nucleus.
- Probe interacts coherently with all nucleons.





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  - Probe interacts coherently with all nucleons.
  - Probes with transverse resolution  $1/Q^2 (\ll \Lambda^2_{\text{QCD}}) \sim 1 \text{ fm}^2$  will see large colour charge fluctuations.
  - This kick experienced in a random walk is the resolution scale.



# The Nuclear “Oomph Factor”

Simple geometric considerations lead to:

$$Q_S^2 \propto \frac{\alpha_s x G(x, Q_S^2)}{\pi R_A^2} \quad \text{HERA: } xG \propto \frac{1}{x^{1/3}} \quad \text{A dependence: } xG_A \propto A$$

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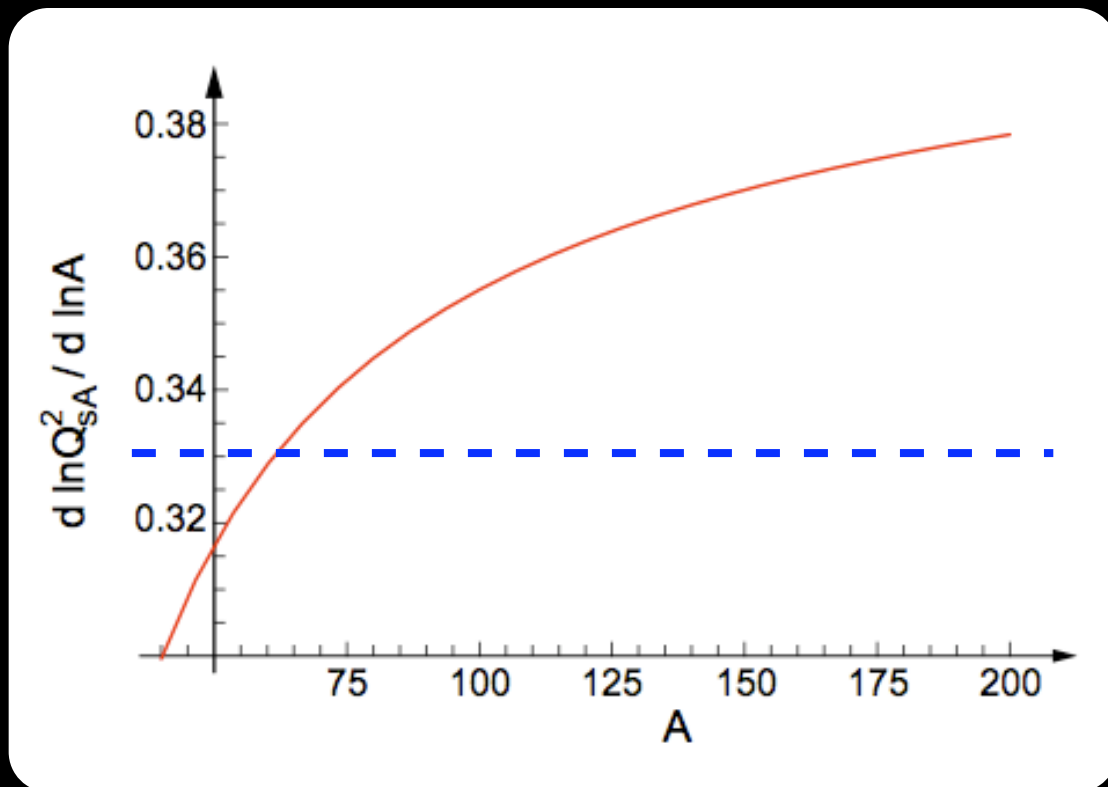
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*Enhancement of  $Q_S$  with  $A$ :  $\Rightarrow$  non-linear QCD regime reached at significantly lower energy in  $e+A$  than in  $e+p$*

# The Nuclear “Oomph Factor”

More sophisticated analyses  
⇒ confirm (exceed) pocket  
formula for high A

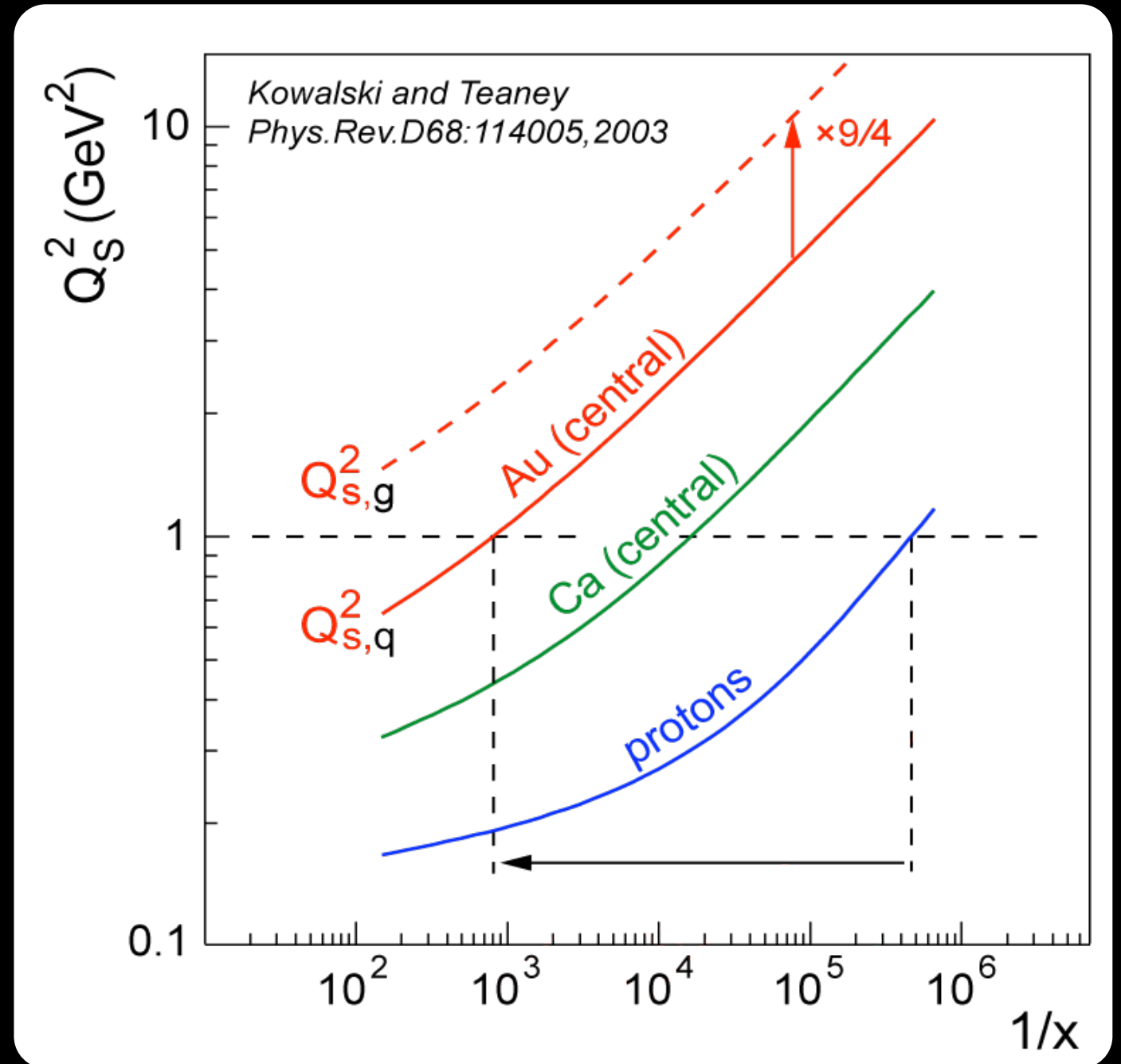
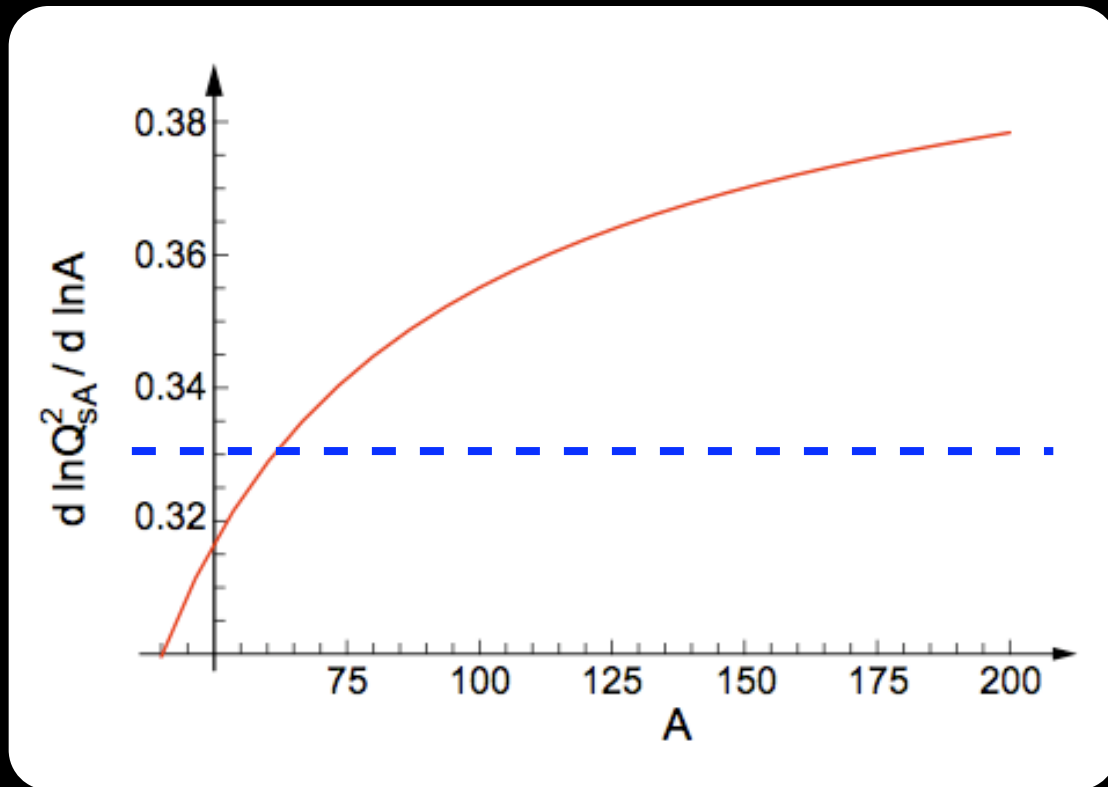
e.g. Kowalski, Lappi and Venugopalan,  
PRL 100, 022303 (2008); Armesto et  
al., PRL 94:022002; Kowalski, Teaney,  
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One would require an energy in e+p  
 $\sim 10\text{-}100 \times e+A$  to get to same  $Q_s^2$

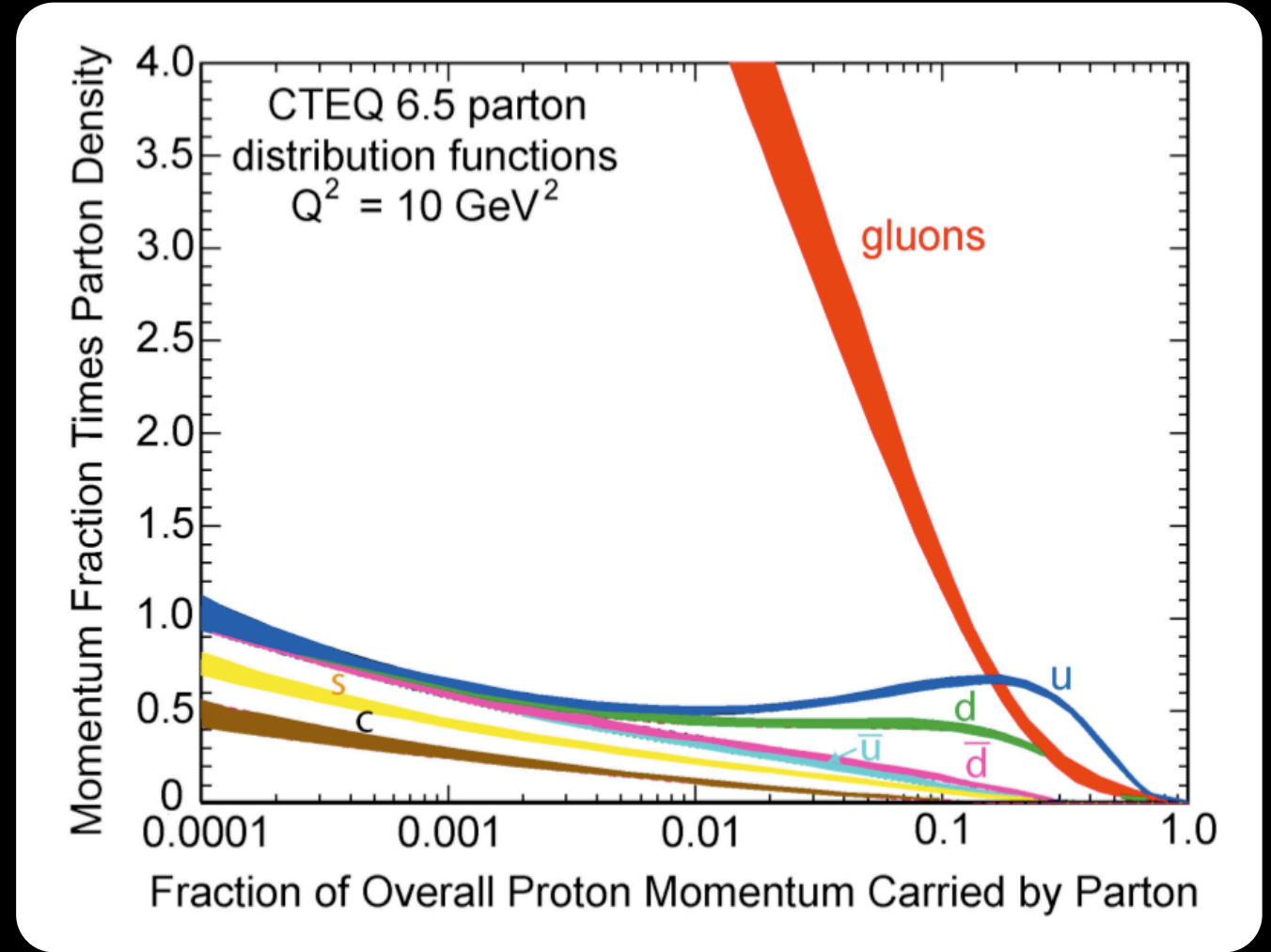
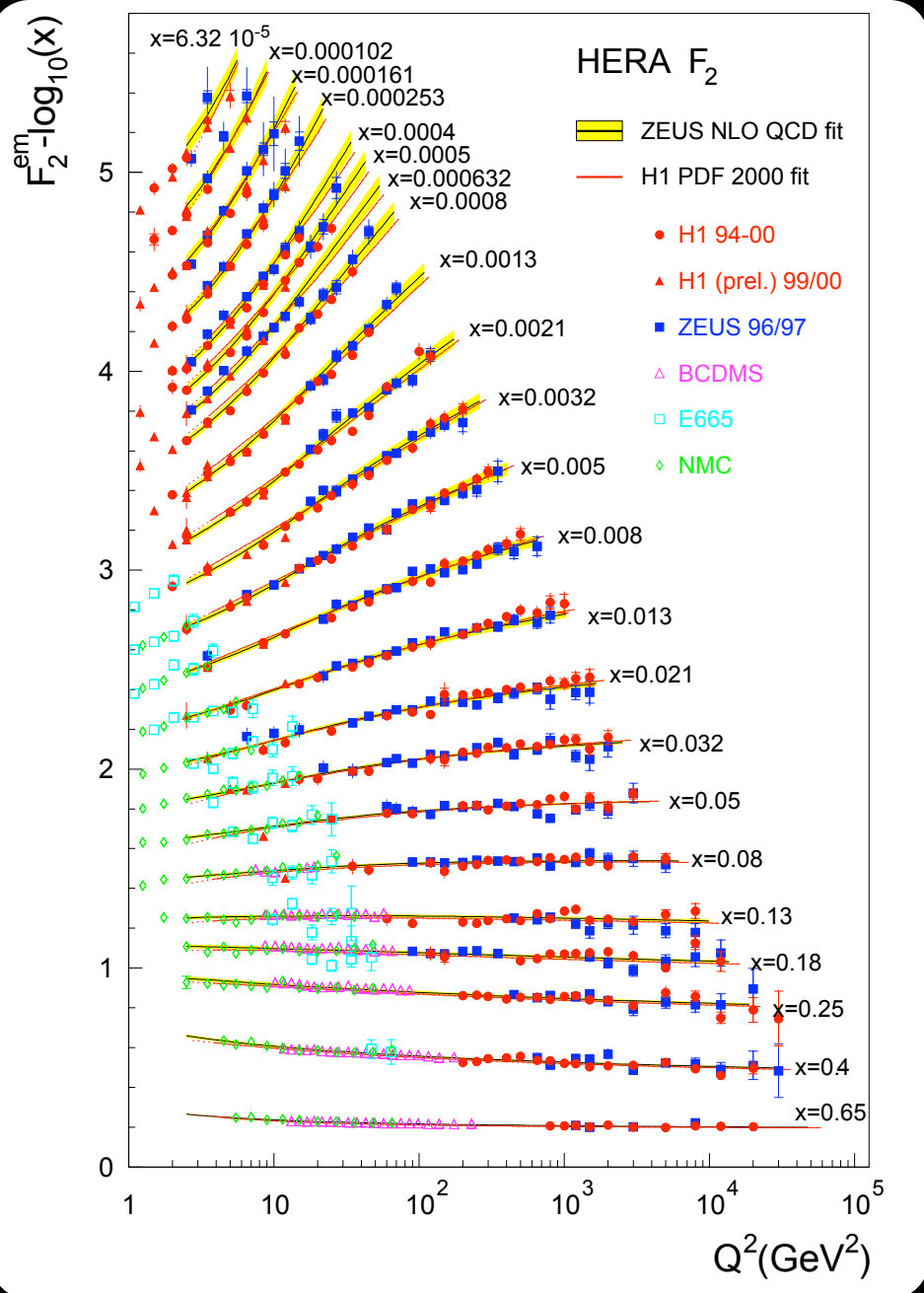
# 4 Key Measurements in e+A Physics

- Momentum distribution of gluons in nuclei?
  - Extract via scaling violation in  $F_2$ :  $\partial F_2 / \partial \ln Q^2$
  - Direct Measurement:  $F_L \sim xG(x, Q^2)$  - requires  $\sqrt{s}$  scan
  - Inelastic vector meson production (e.g.  $J/\Psi$ ,  $\rho$ )
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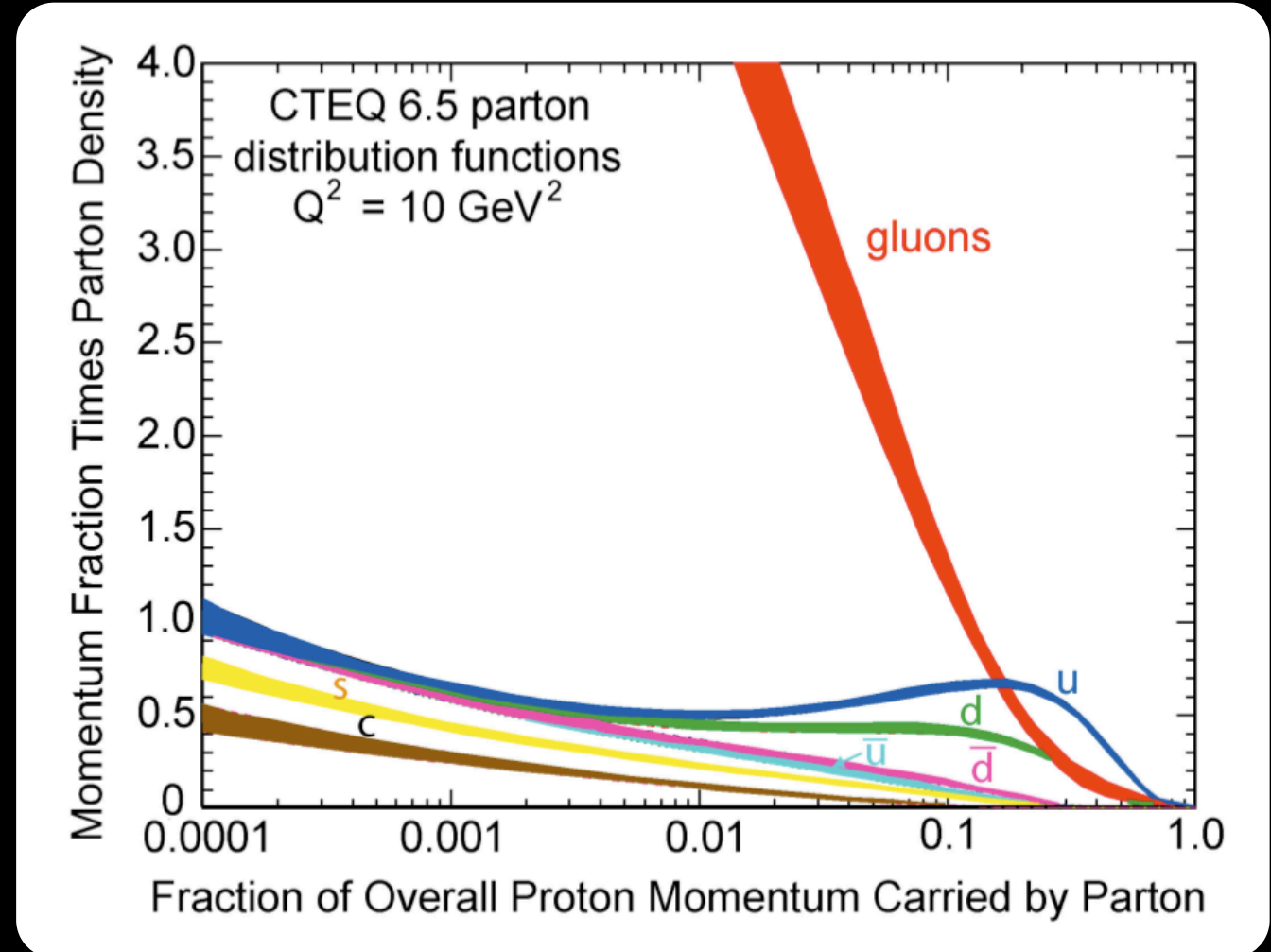
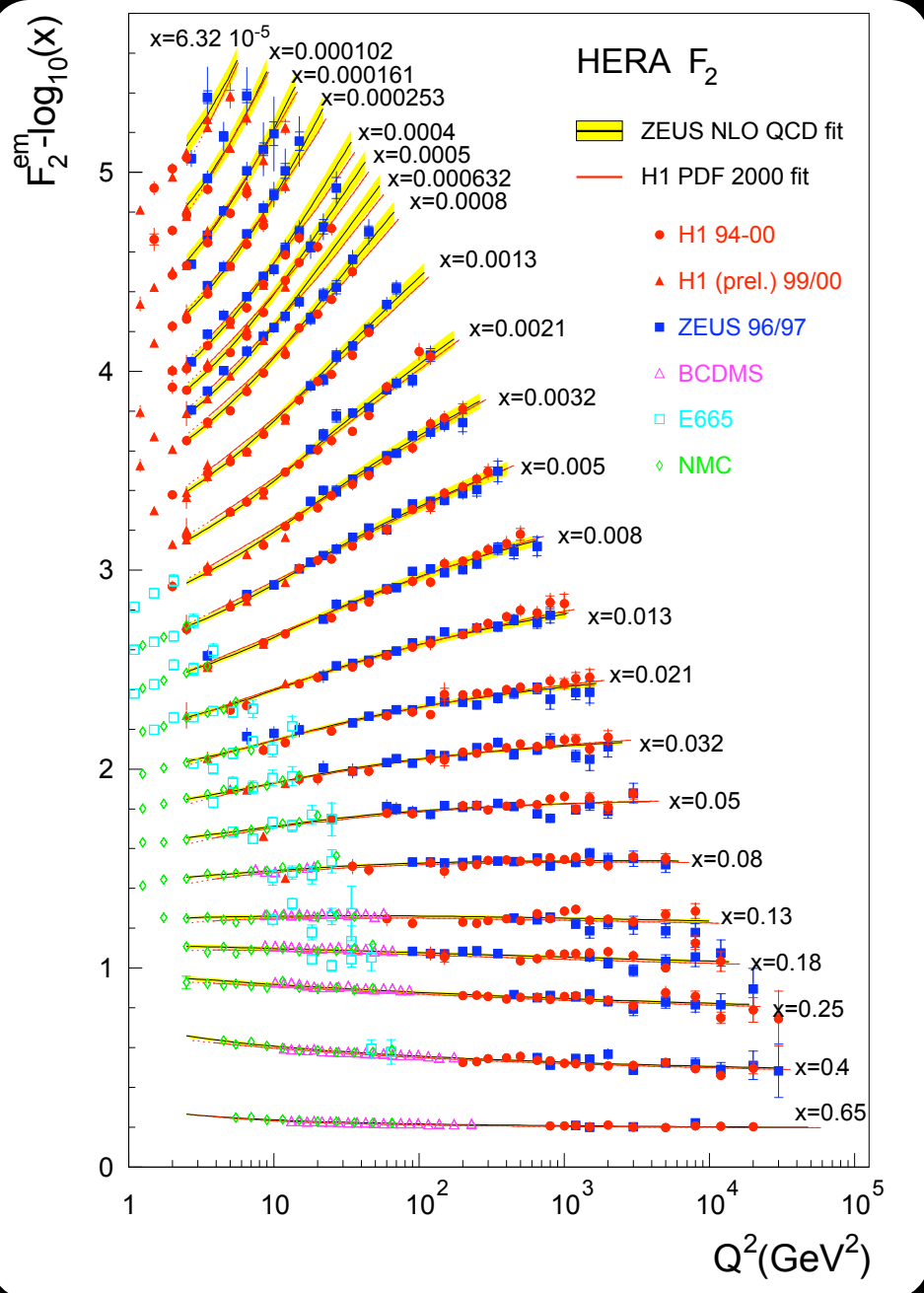
# Glun momentum distributions: i) $F_2$ scaling violation

$$\frac{d^2 \sigma^{ep \rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[ \left( 1 - y + \frac{y^2}{2} \right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$



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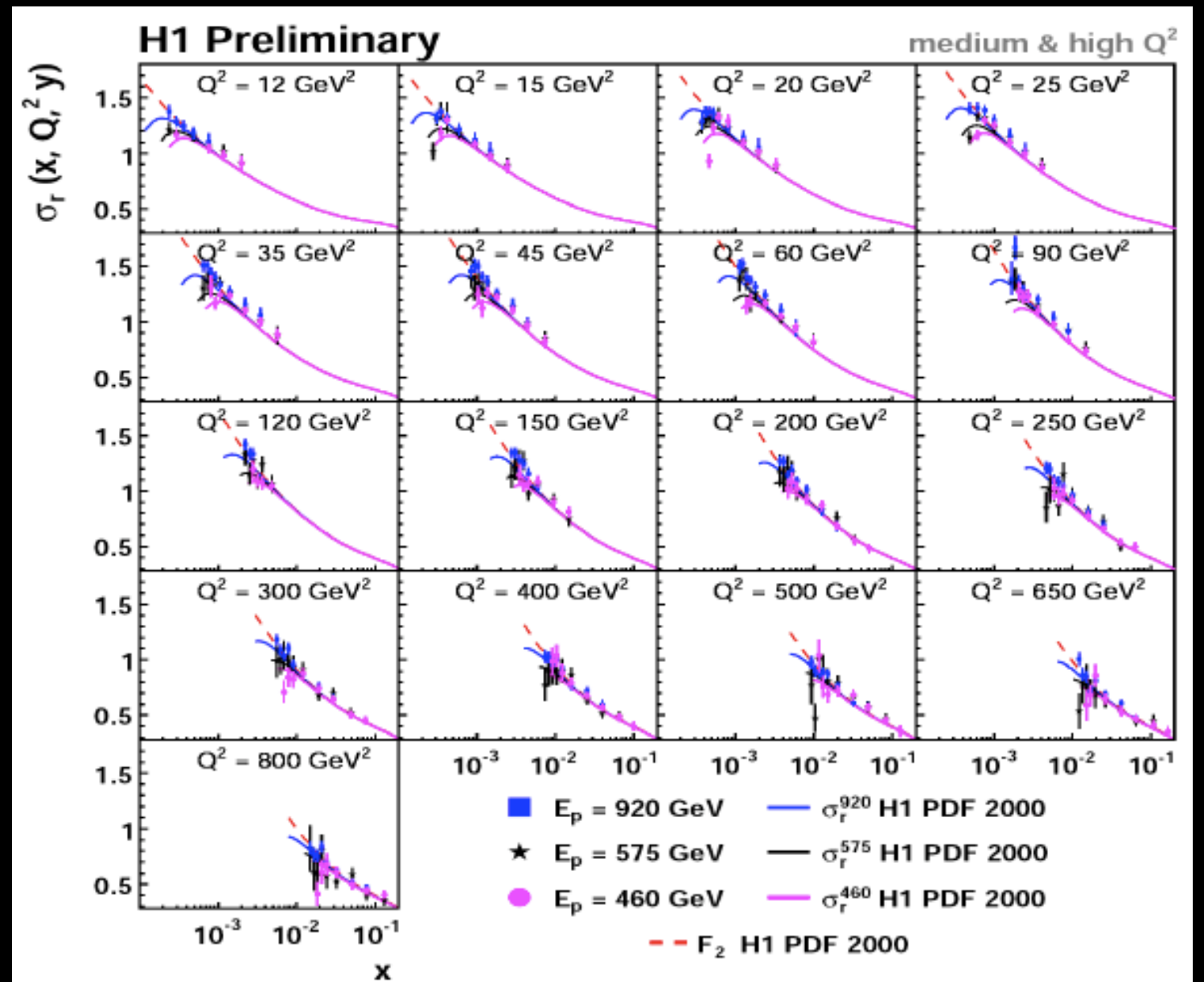
Scaling violation of  $dF_2/d\ln Q^2$  and linear DGLAP evolution  $\Rightarrow xG(x, Q^2)$

# Gluon momentum distributions: ii) $F_L$ measured directly

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$$F_L \sim \alpha_s xG(x, Q^2)$$

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Preliminary  $F_L$  measurements

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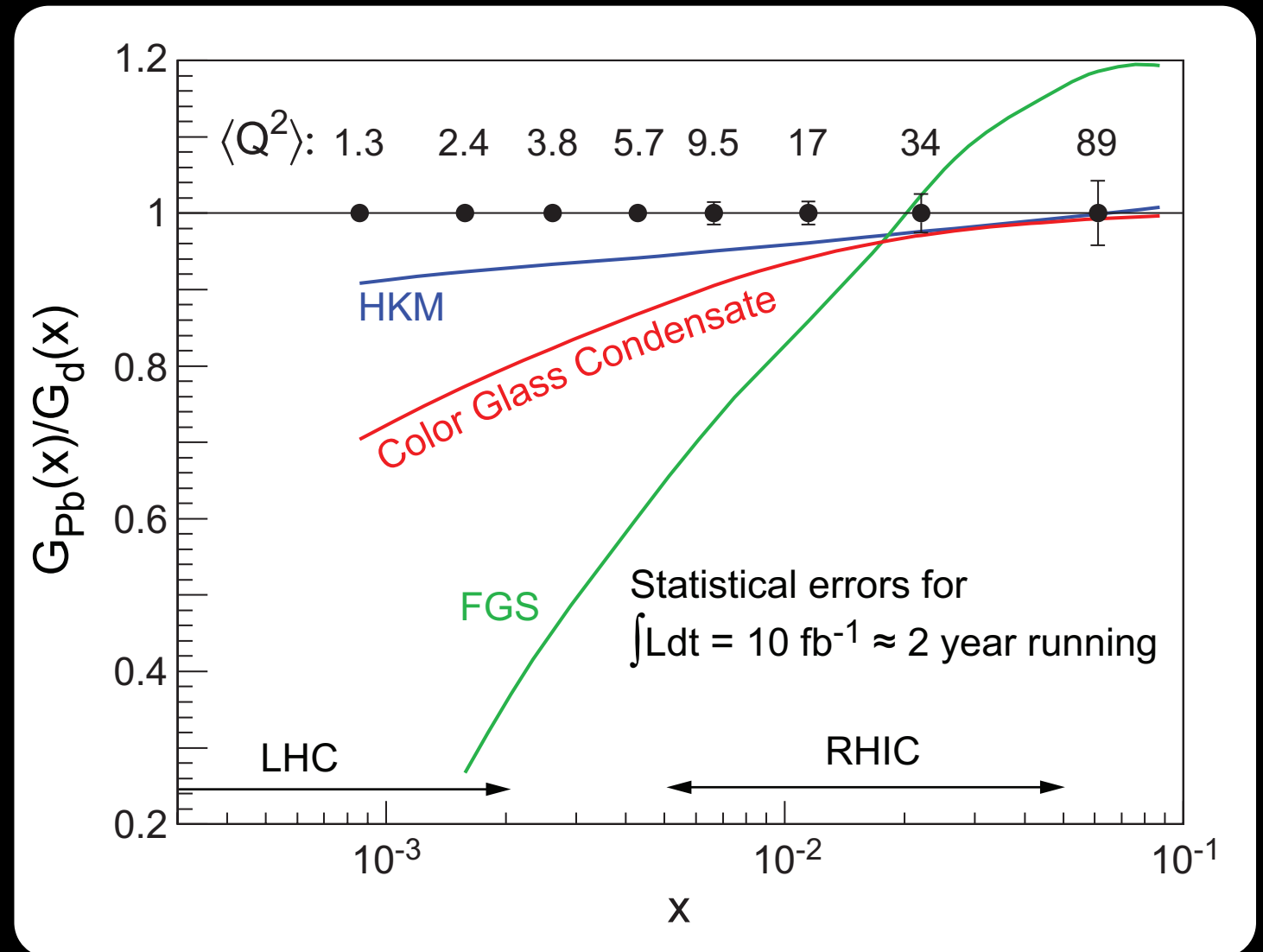
Here:

$$\begin{aligned} \int L dt &= 4/A \text{ fb}^{-1} \text{ (10+100) GeV} \\ &= 4/A \text{ fb}^{-1} \text{ (10+50) GeV} \\ &= 2/A \text{ fb}^{-1} \text{ (5+50) GeV} \end{aligned}$$

statistical error only

Syst. studies of  $F_L(A, x, Q^2)$ :

- $xG(x, Q^2)$  with great precision
- Distinguish between models



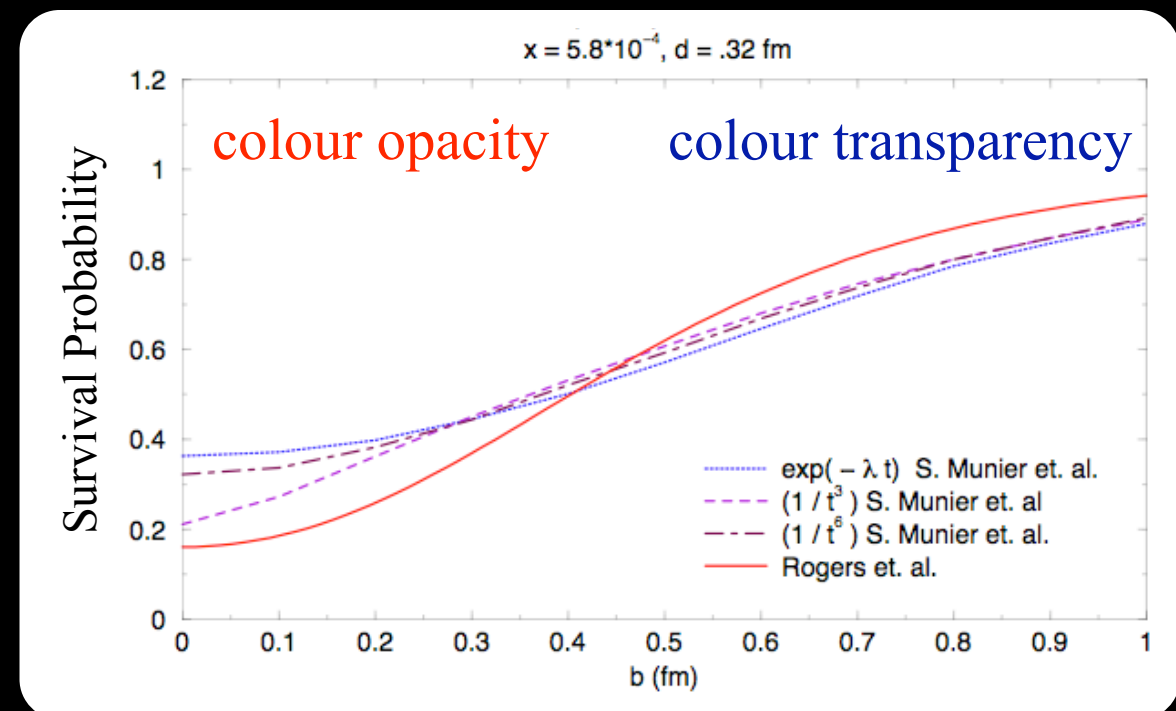
HKM and FGS are "standard" shadowing parameterizations that are evolved with DGLAP

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  - Inelastic vector meson production (e.g.  $J/\Psi$ ,  $\rho$ )
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- Space-time distribution of gluons in nuclei?
  - Exclusive final states (e.g.  $\rho$ ,  $J, \Psi$ )
  - Deep Virtual Compton Scattering (DVCS) -  $\sigma \sim A^{4/3}$
  - $F_2$ ,  $F_L$  for various impact parameters

# Gluon space-time distributions

- In the “colour dipole” picture:
  - virtual photon fluctuates into a qq-bar dipole and scatters coherently on the nucleus
  - calculate the survival probability of qq-bar pair to propagate through the target without interacting
  - Calculate by measuring the vector meson cross-section
- pQCD  $\Rightarrow$  survival  $\sim 1$
- dipole models  $\Rightarrow \sim x5$  smaller
- HERA data limited on this
- $b$  profile of nuclei more uniform



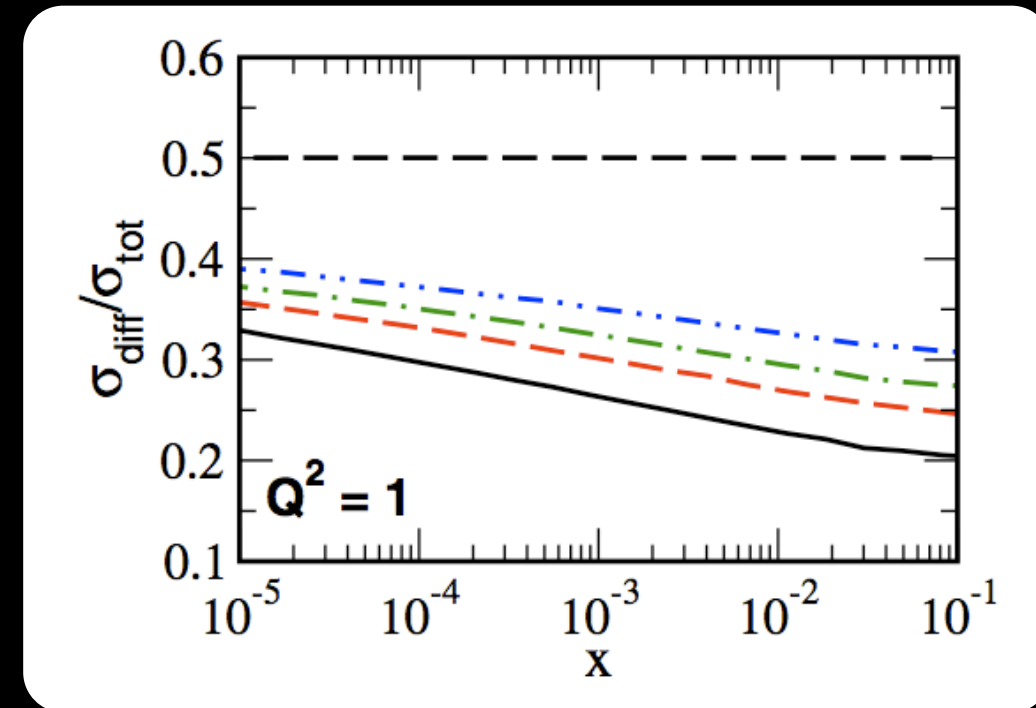
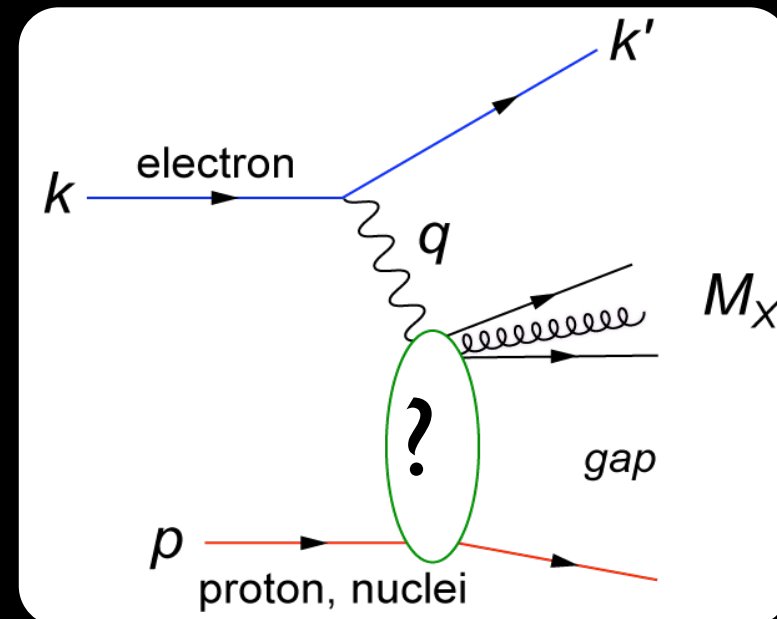
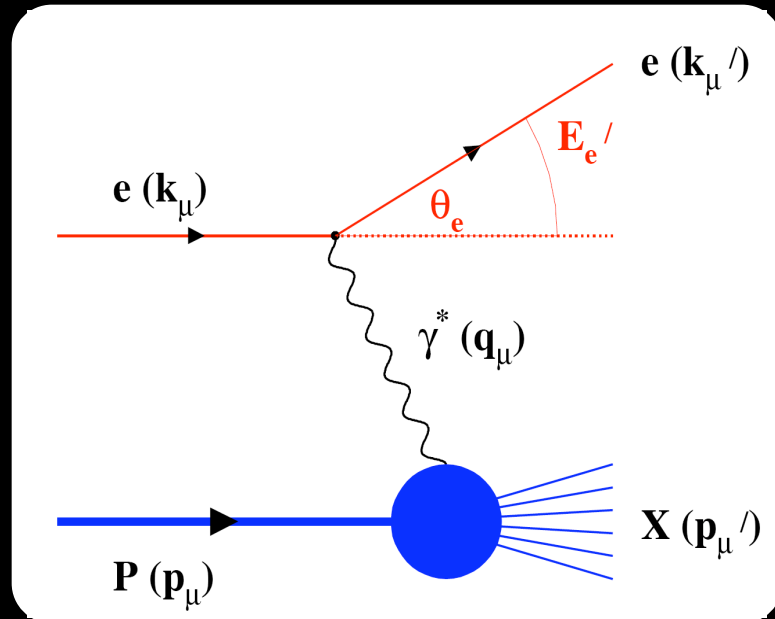


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- **Role of colour-neutral (Pomeron) excitations?**
  - Diffractive cross-section:  $\sigma_{\text{diff}}/\sigma_{\text{tot}}$  ( $\sim 10\%$ : HERA e+p; 30%? EIC e+A?)
  - Diffractive structure functions and vector meson productions
  - Abundance and distribution of rapidity gaps

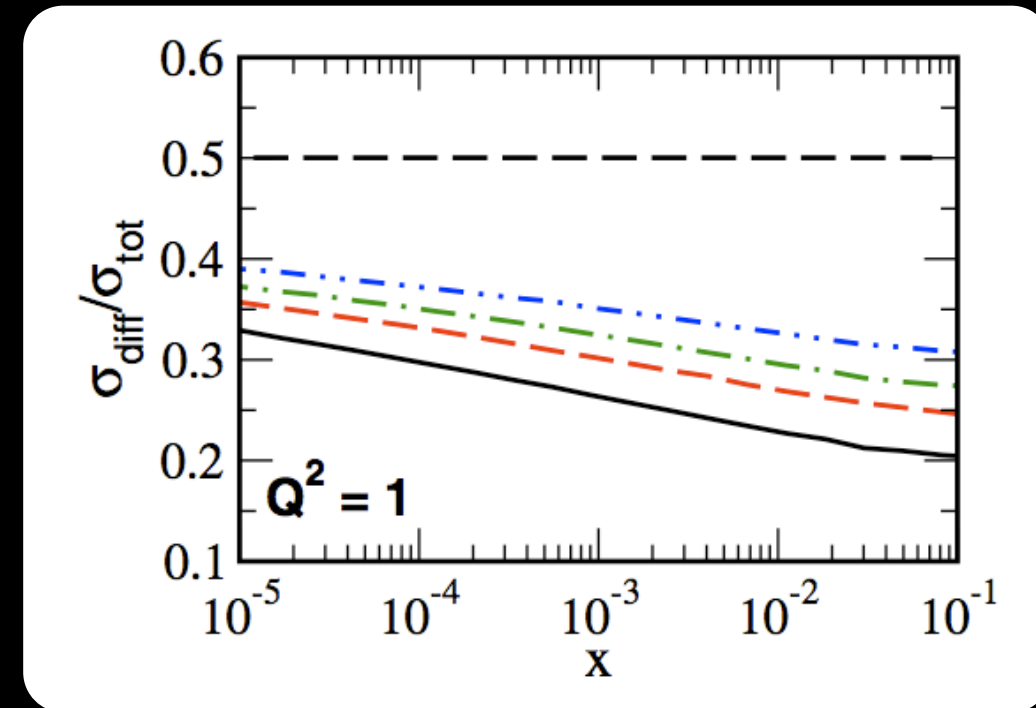
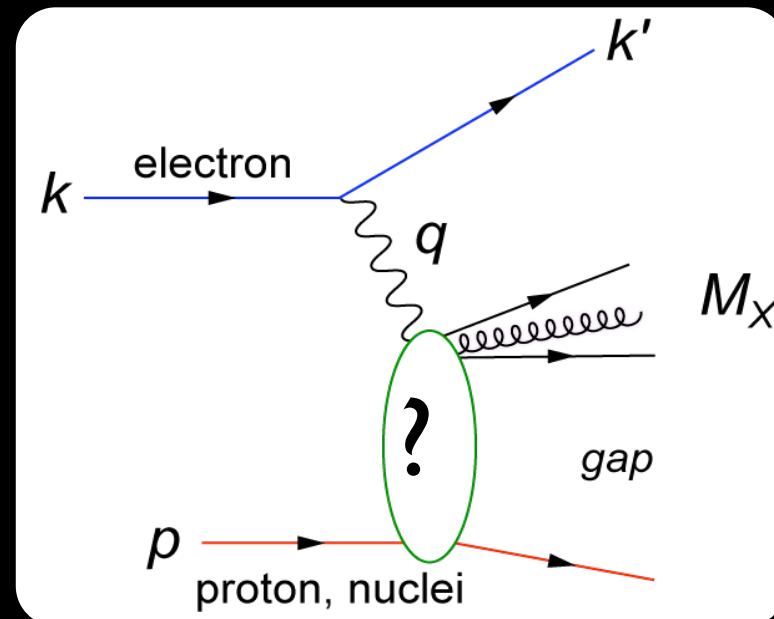
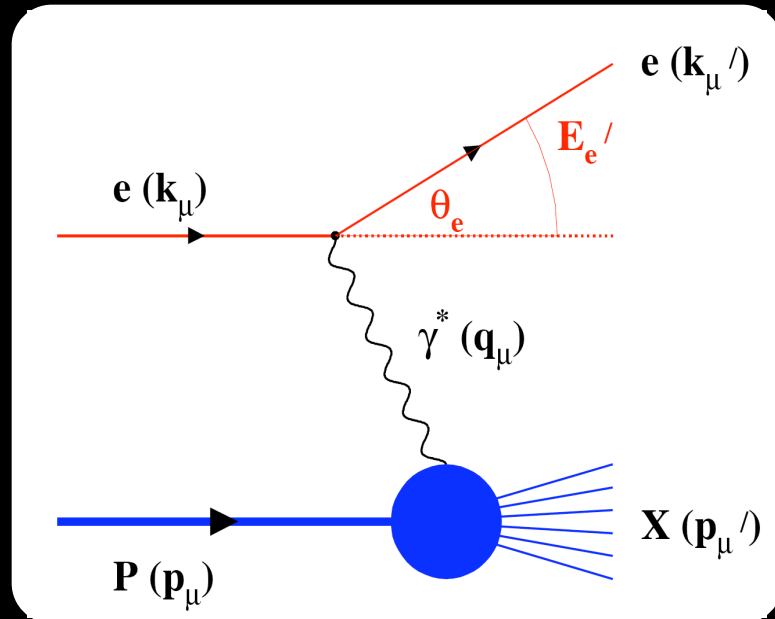
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## Diffractive physics in e+A:



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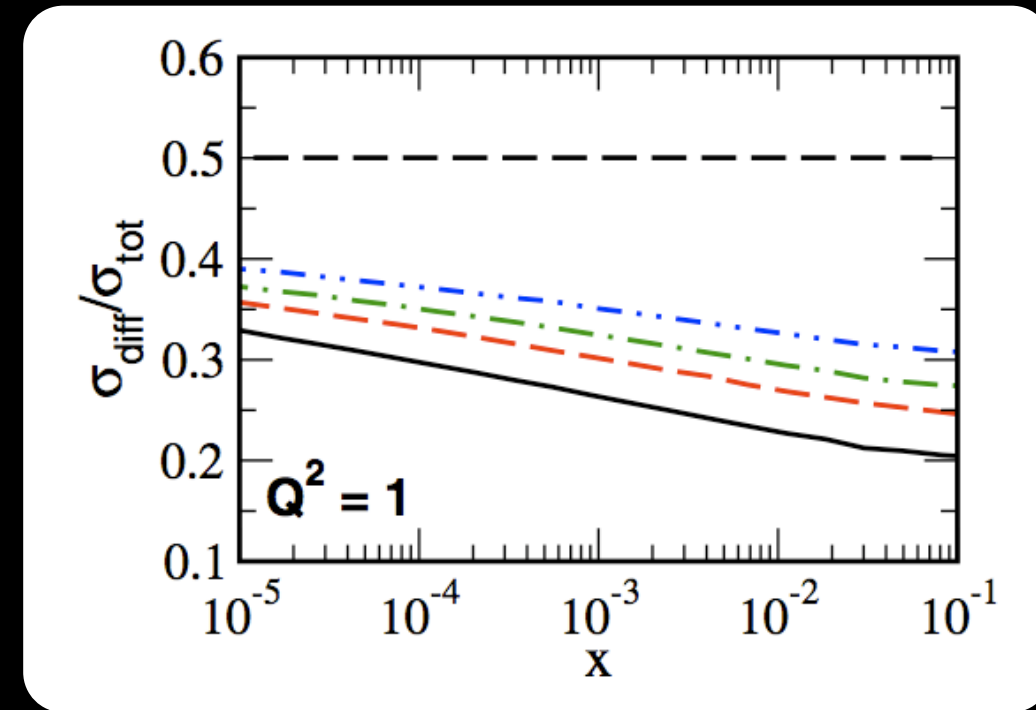
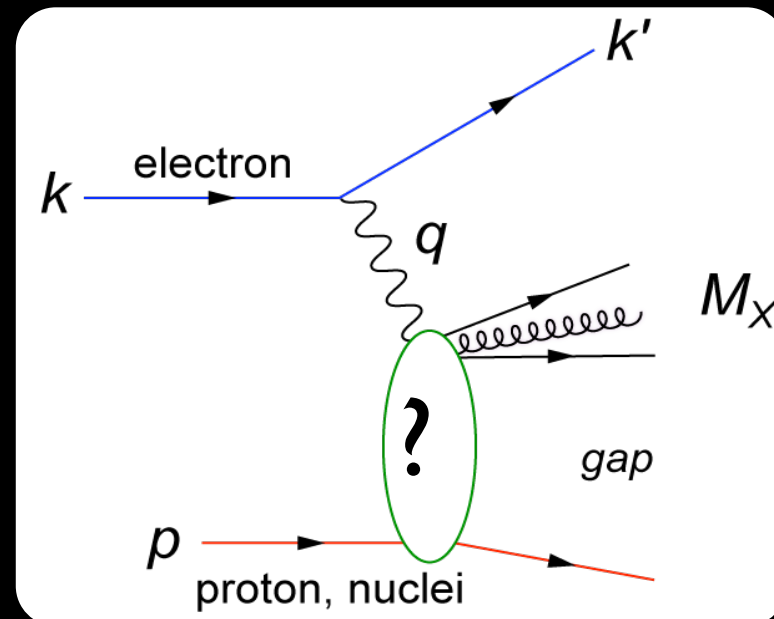
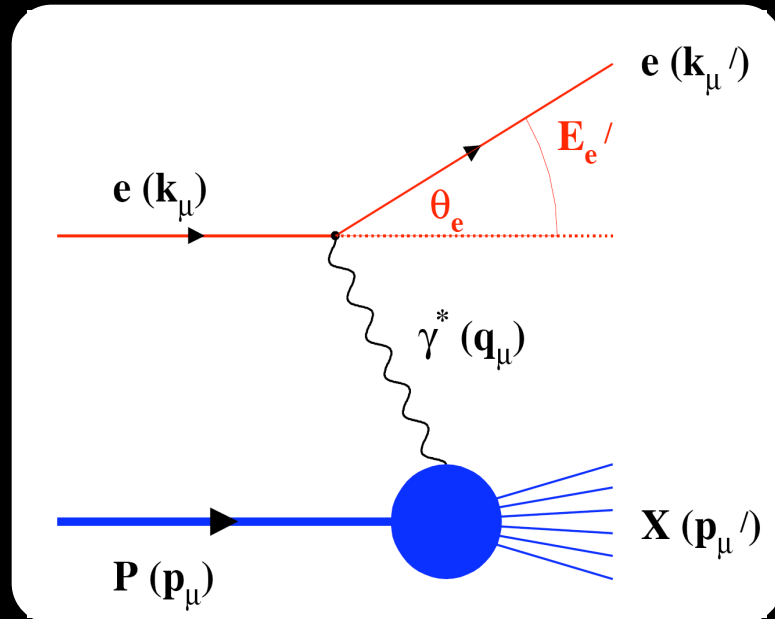
## Diffractive physics in $e+A$ :



- HERA/ep: 15% of all events are hard diffractive
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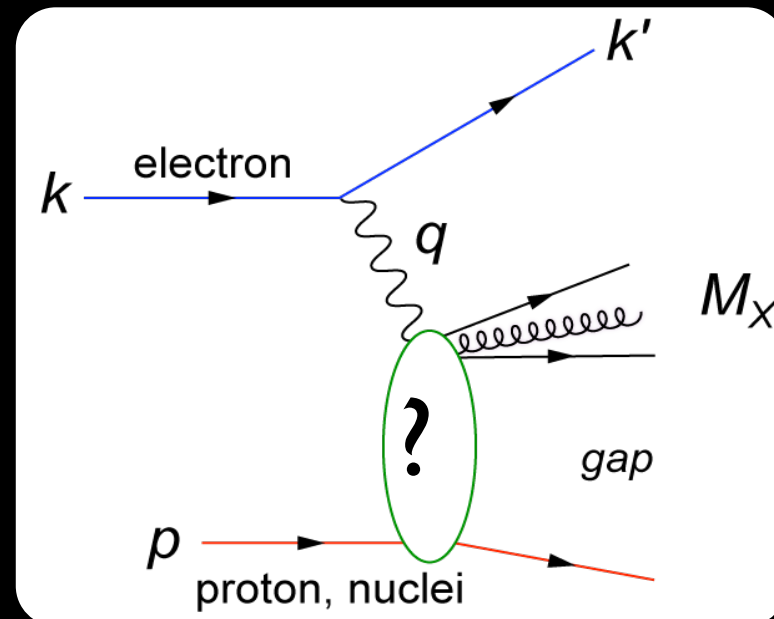
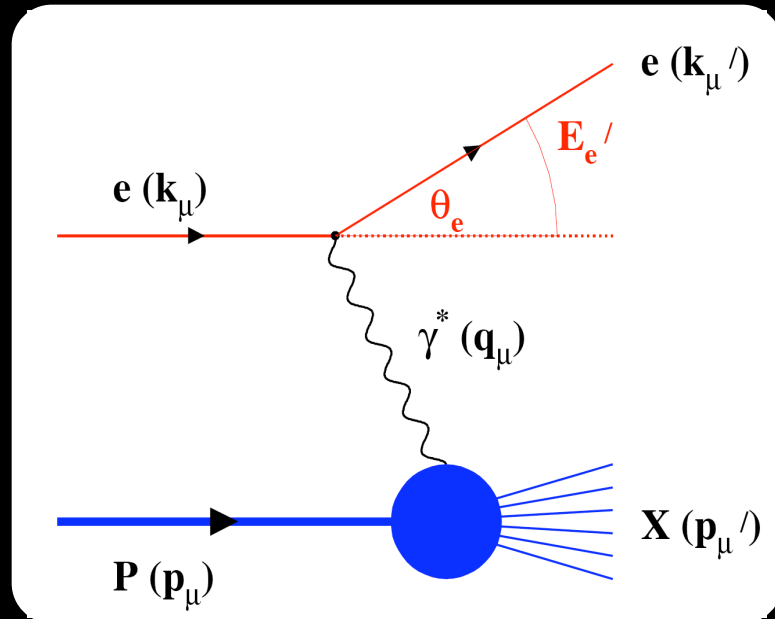


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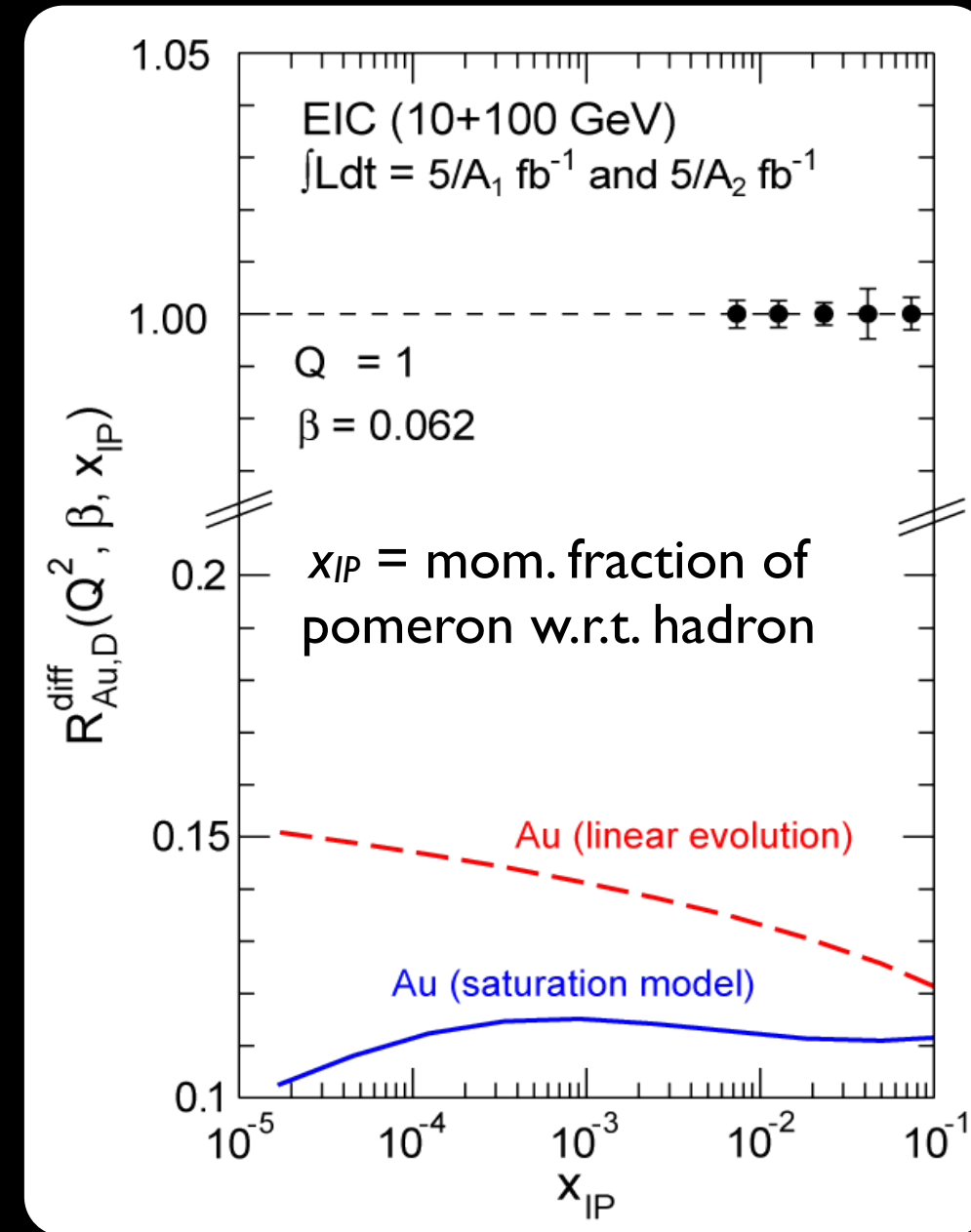
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Curves: Kugeratski, Goncalves, Navarra, EPJ C46, 413

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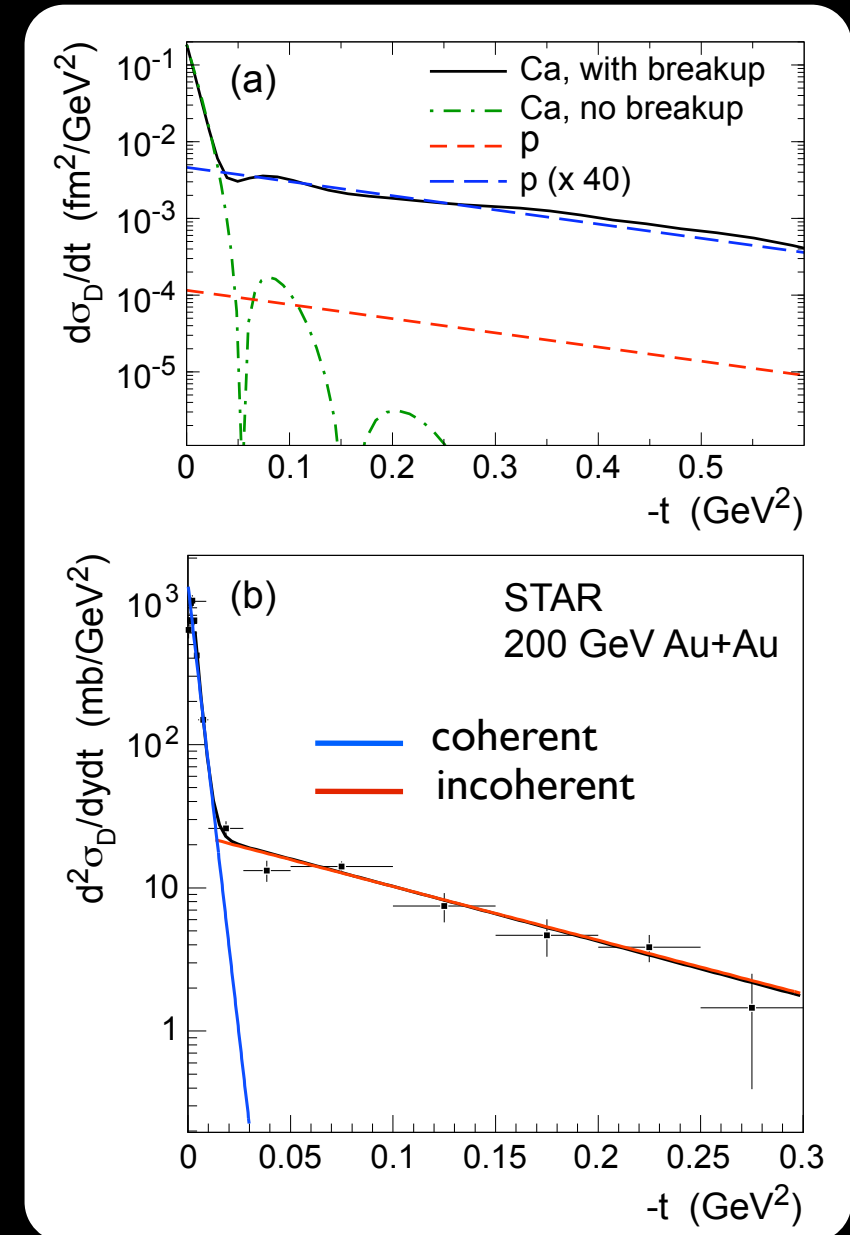
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- Distinguish between linear evolution and saturation models



# Role of colour-neutral (Pomeron) excitations

$$\frac{d\sigma}{dt} \Big|_{t=0} (\gamma^* A \rightarrow V A) \propto \alpha_s^2 [G_A(x, Q^2)]^2$$

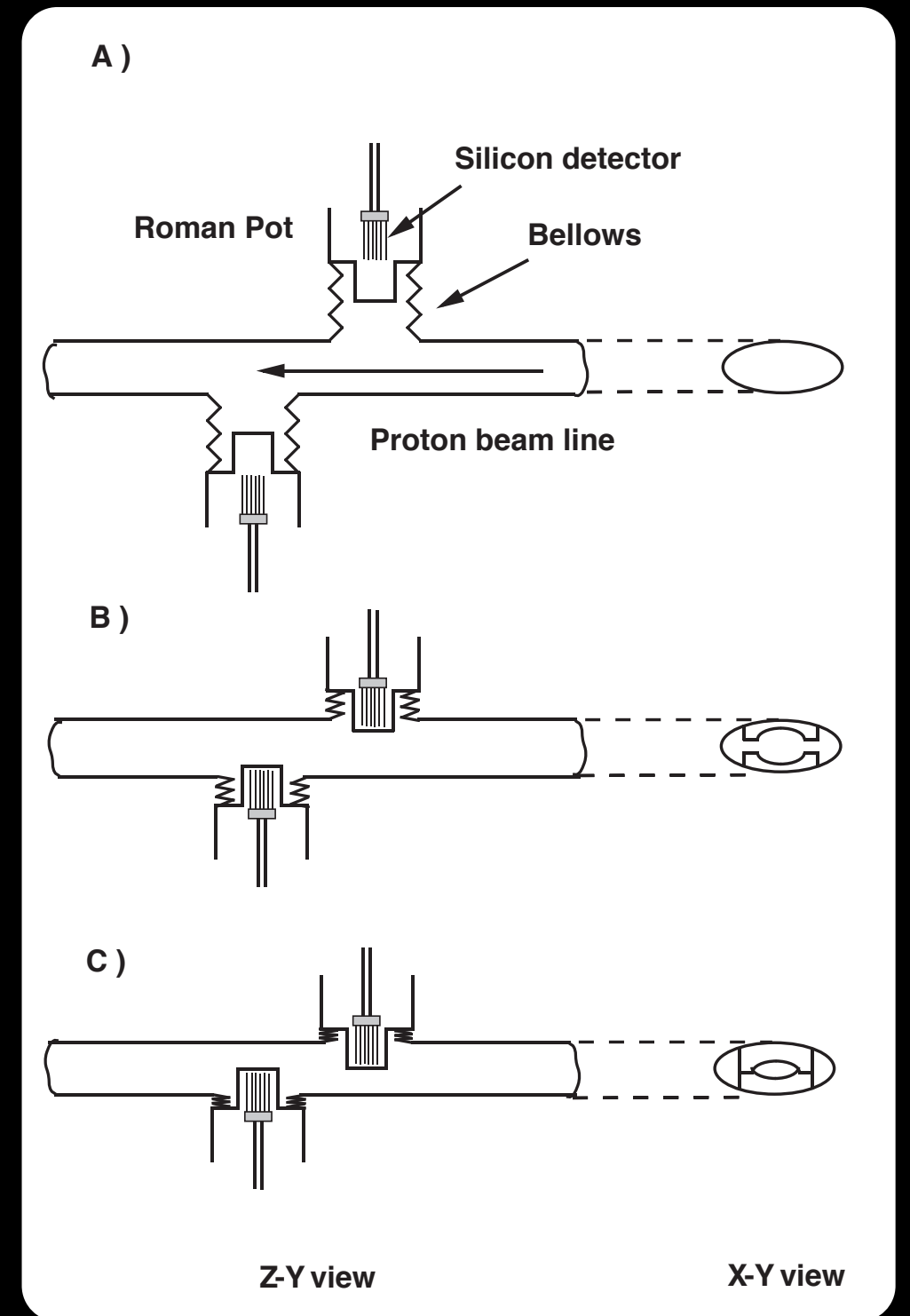
- Knowledge of  $t$  is important
  - small  $t \Rightarrow$  coherent diffraction
  - large  $t \Rightarrow$  incoherent diffraction
- Results from STAR UPC Au+Au collisions
  - coherent diffraction  $\Rightarrow t < 0.03 \text{ GeV}^2$
  - incoherent diffraction  $\Rightarrow t > 0.03 \text{ GeV}^2$



STAR Ultra-Peripheral Au+Au Collisions:  
Phys. Rev. C 77 (2008) 34910

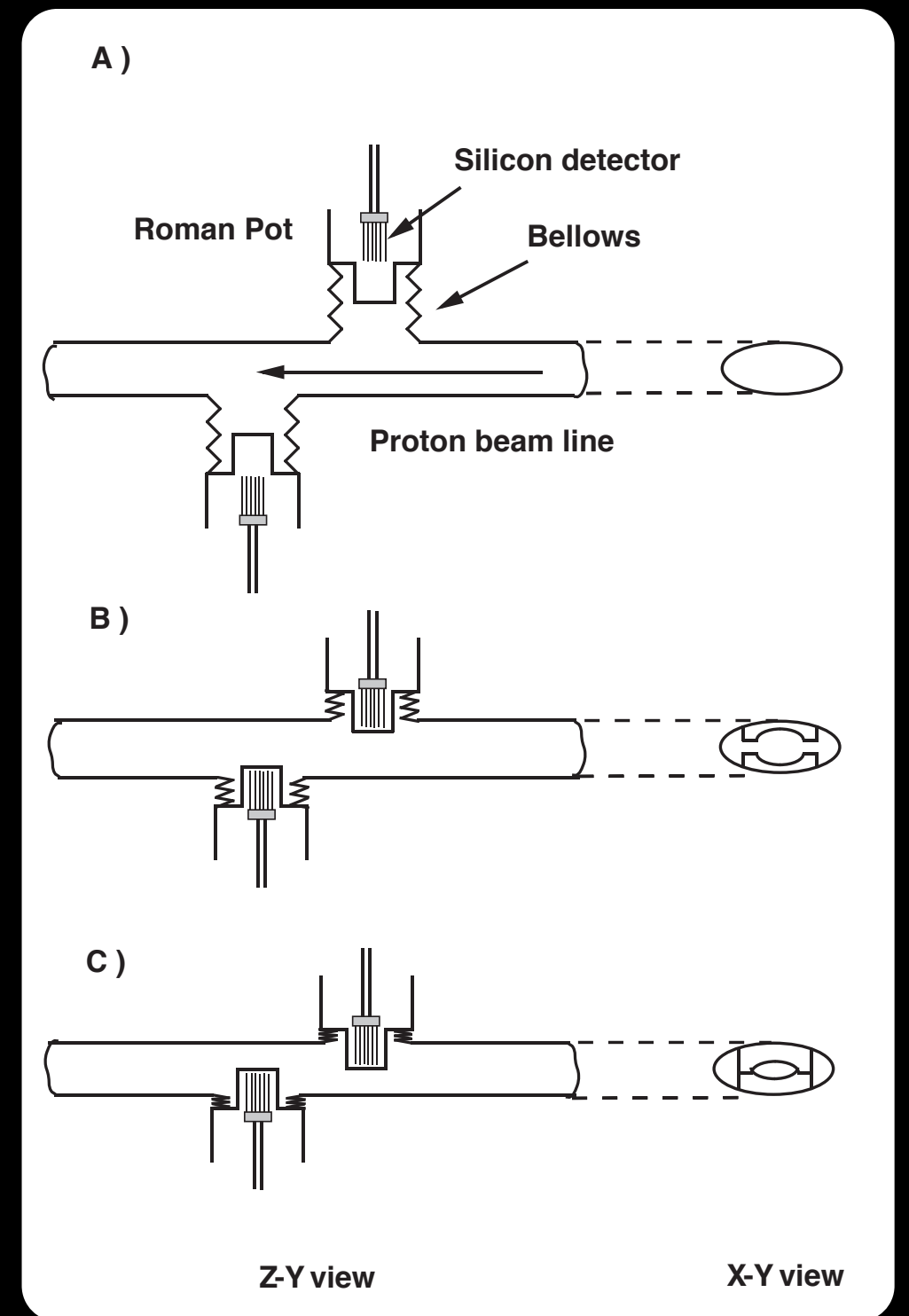


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- Can measure the nucleus if it is separated from the beam in Si (Roman Pot) “beamline” detectors
  - $p_T^{\min} \sim pA\theta_{\min}$ 
    - For beam energies = 100 GeV/n and  $\theta_{\min} = 0.08$  mrad:
- These are large momentum kicks, much greater than the binding energy ( $\sim 8$  MeV)
- Therefore, for large  $A$ , coherently diffractive nucleus cannot be separated from beamline without breaking up

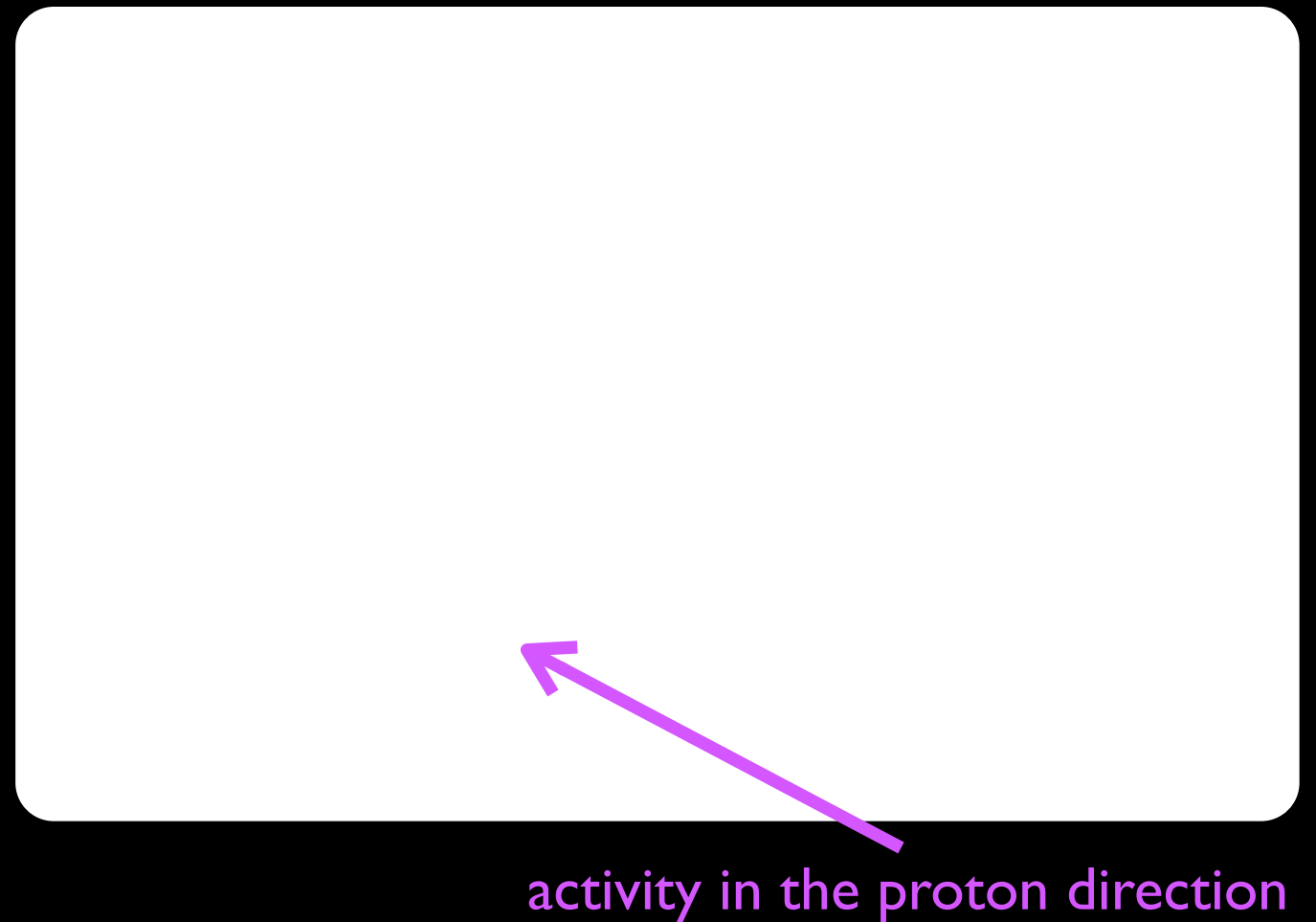
species (A)	$p_T^{\min}$ (GeV/c)
d (2)	0.02
Si (28)	0.22
Cu (64)	0.51
In (115)	0.92
Au (197)	1.58
U (238)	1.90

# How else to measure diffraction in $e+A$ ?

Method used at HERA:

Large Rapidity Gap Method:

In diffractive events, a large gap in rapidity occurs between outgoing  $p$  and final state particles

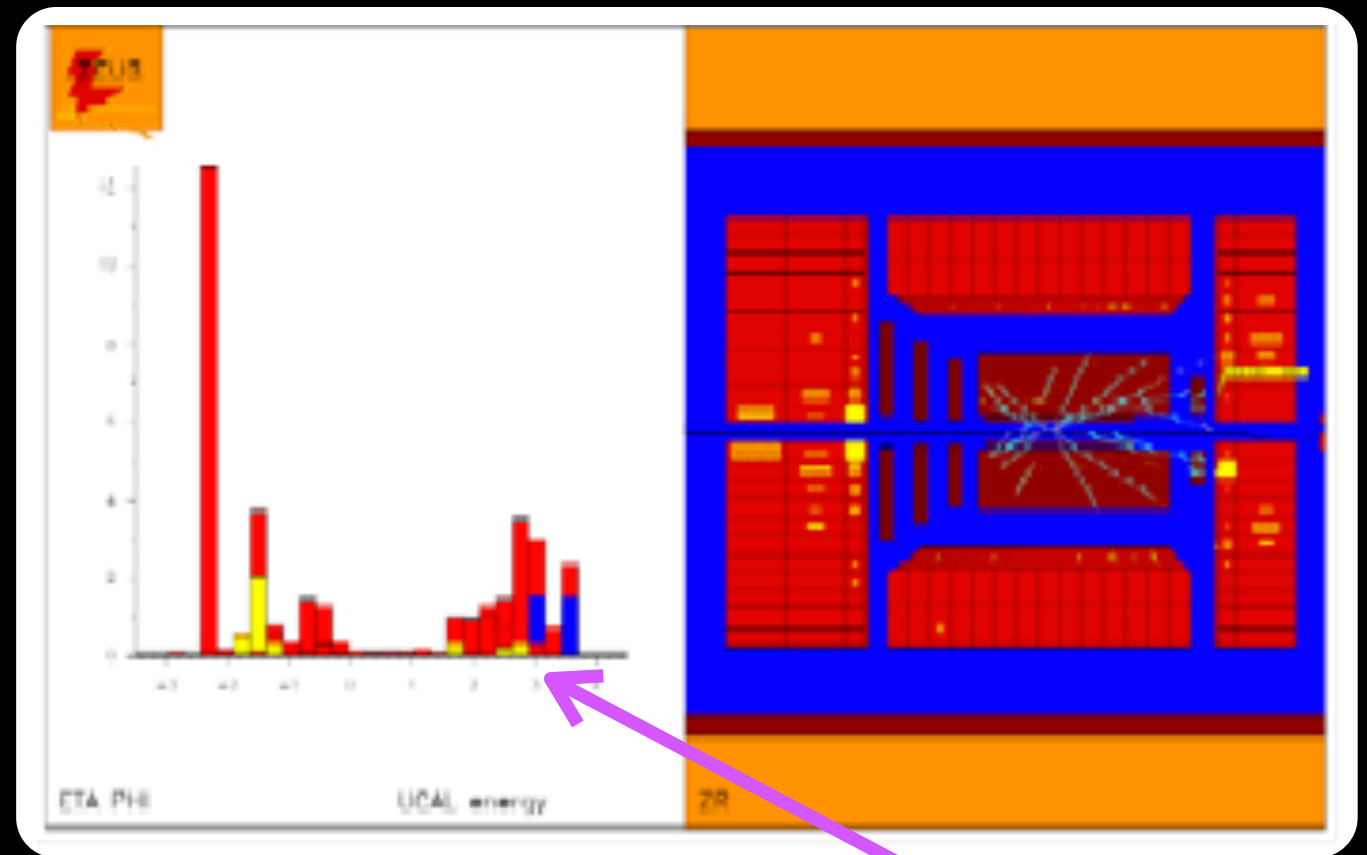


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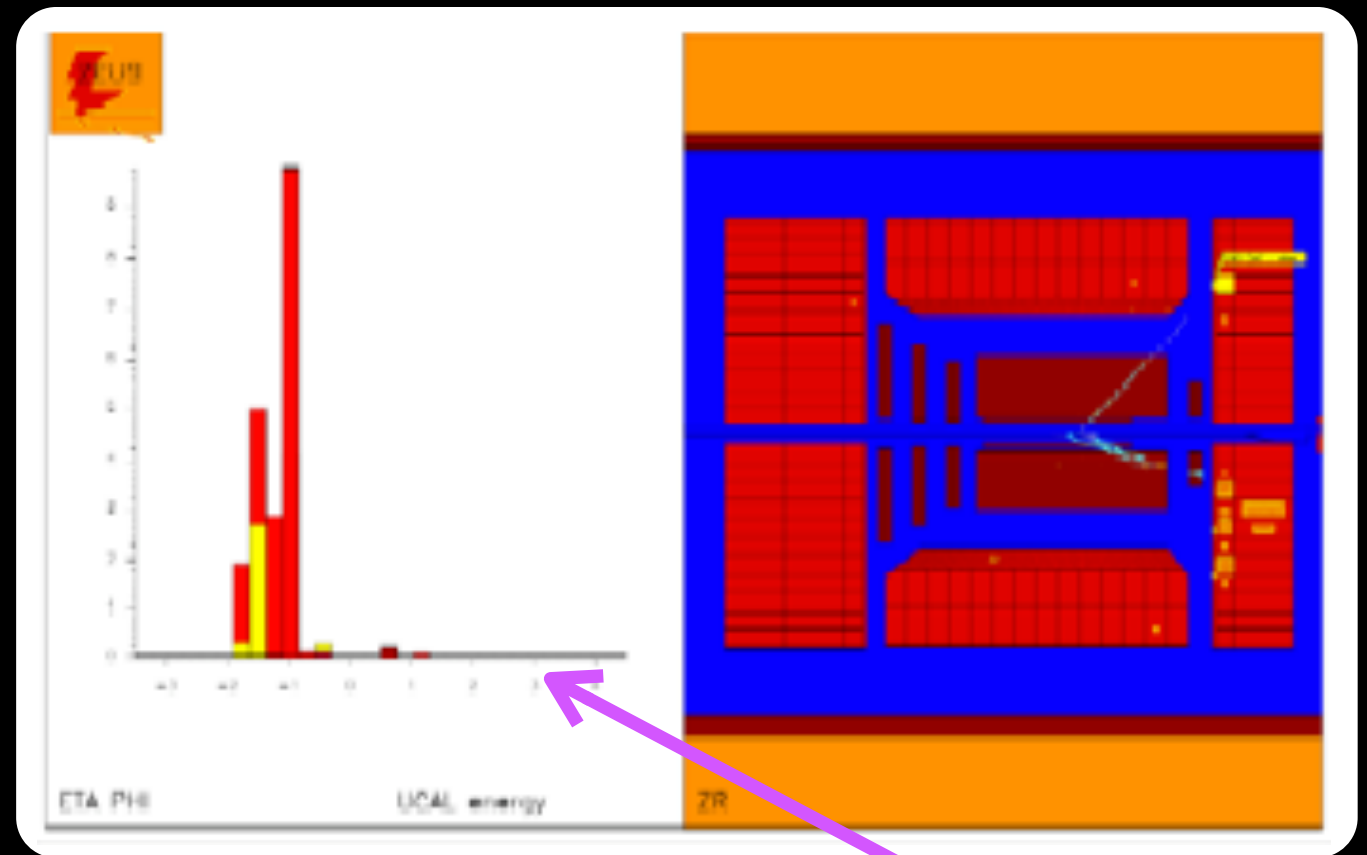
activity in the proton direction

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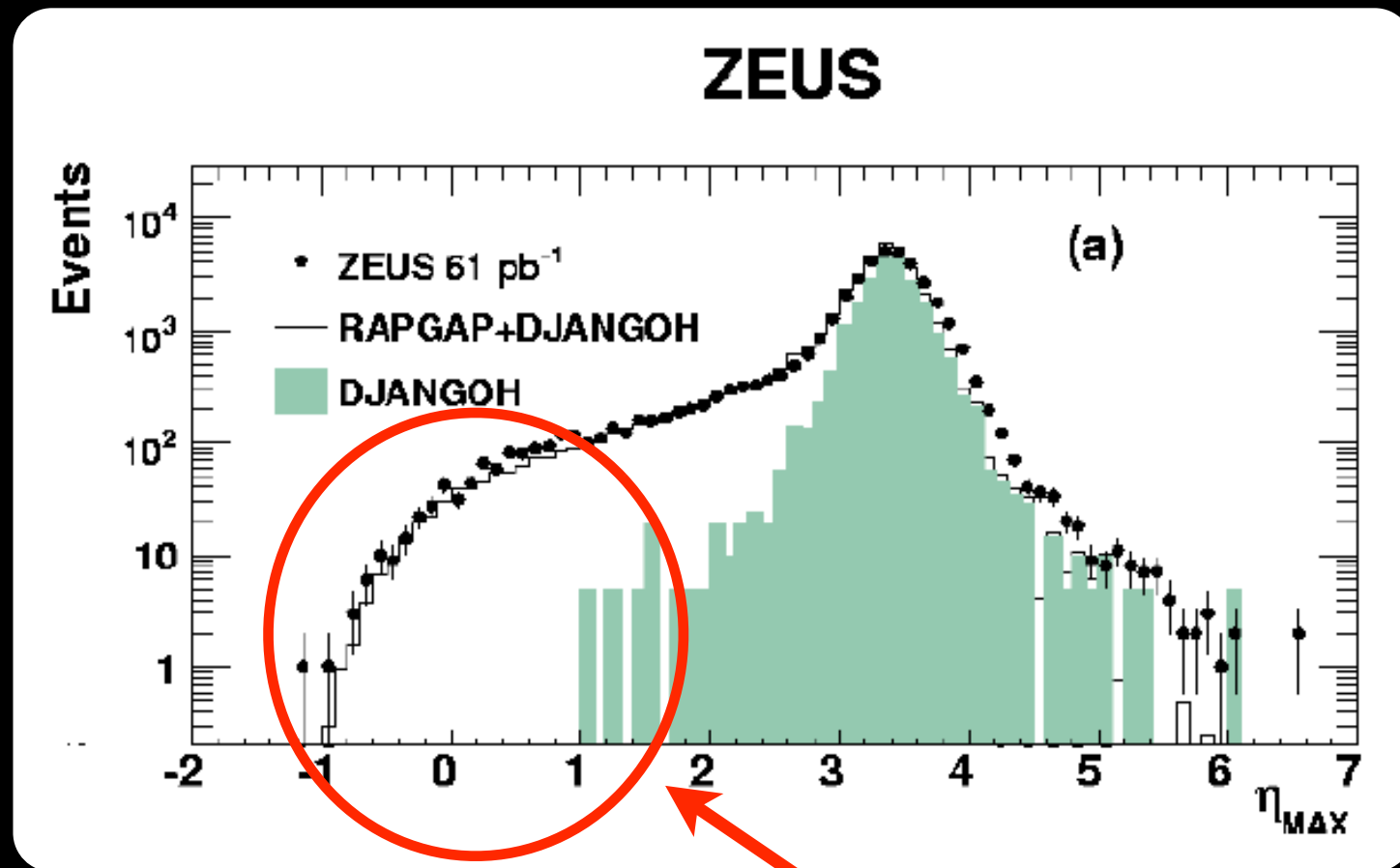
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- At HERA:  $\Delta\eta \sim 7 \Rightarrow$  hadronization reduces this to  $\sim 2.5$
- Pros
  - Lots of statistics
- Cons
  - Sensitive to hadronization models
  - No information on t





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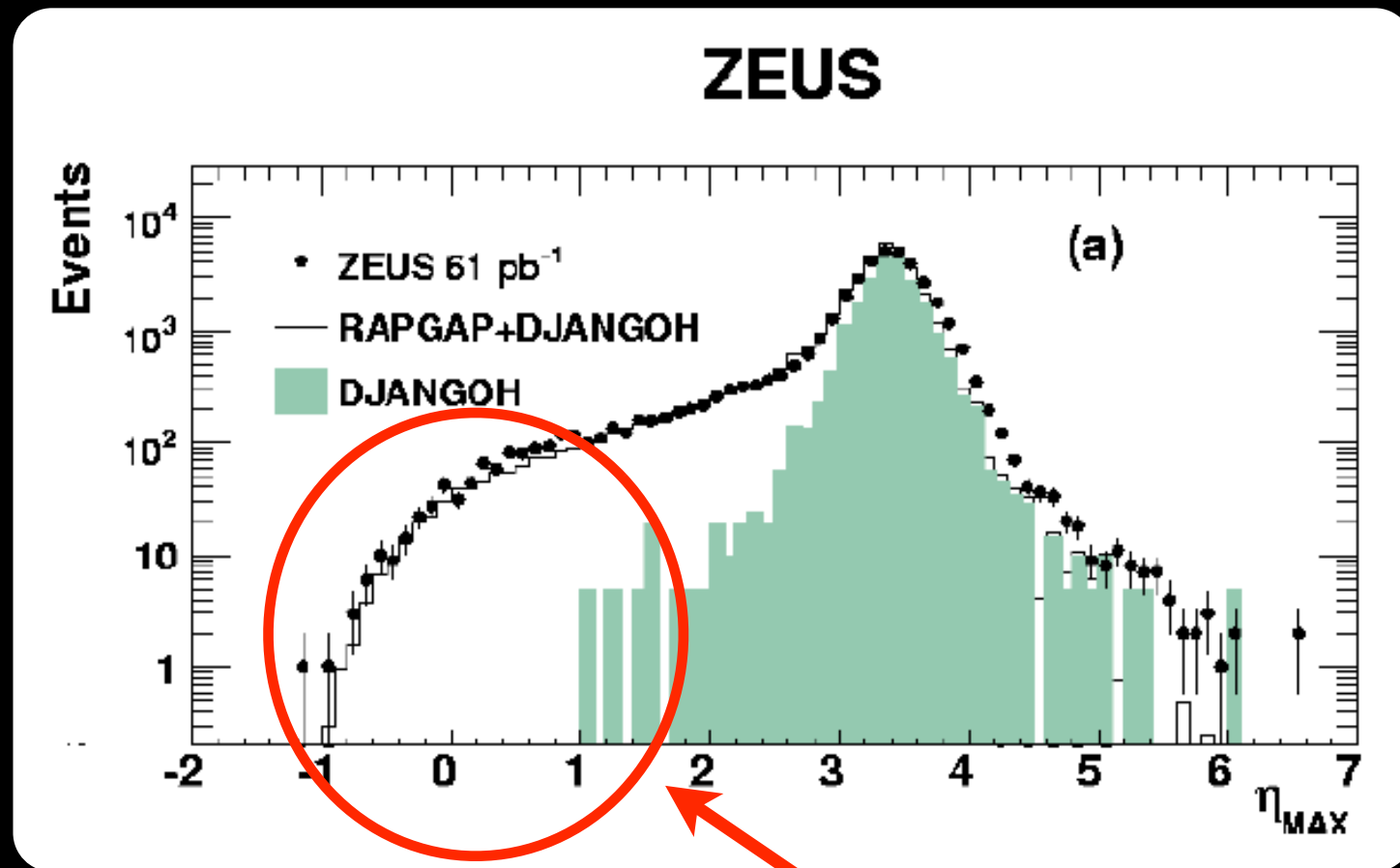
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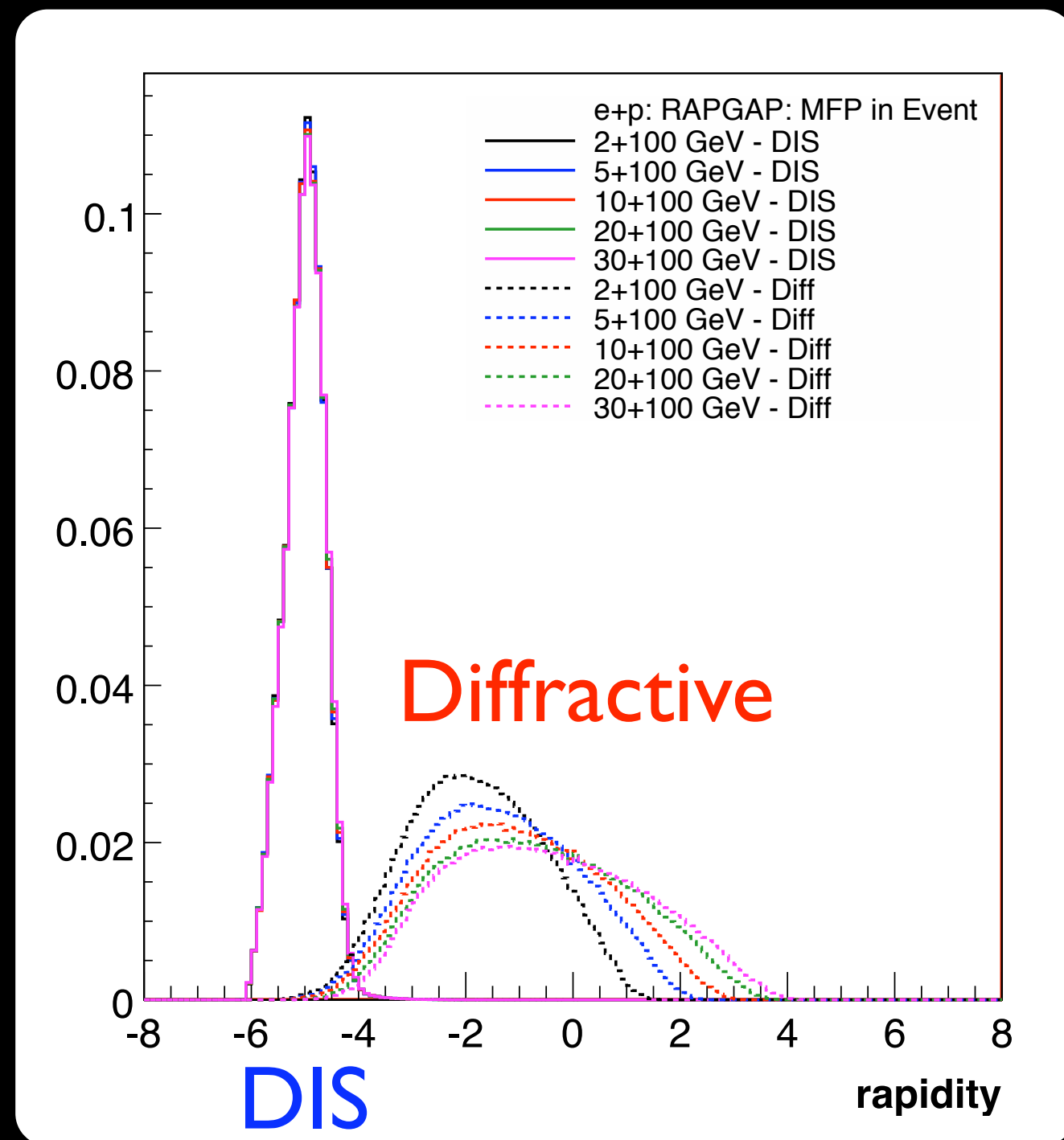
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Can this method be used at an EIC?

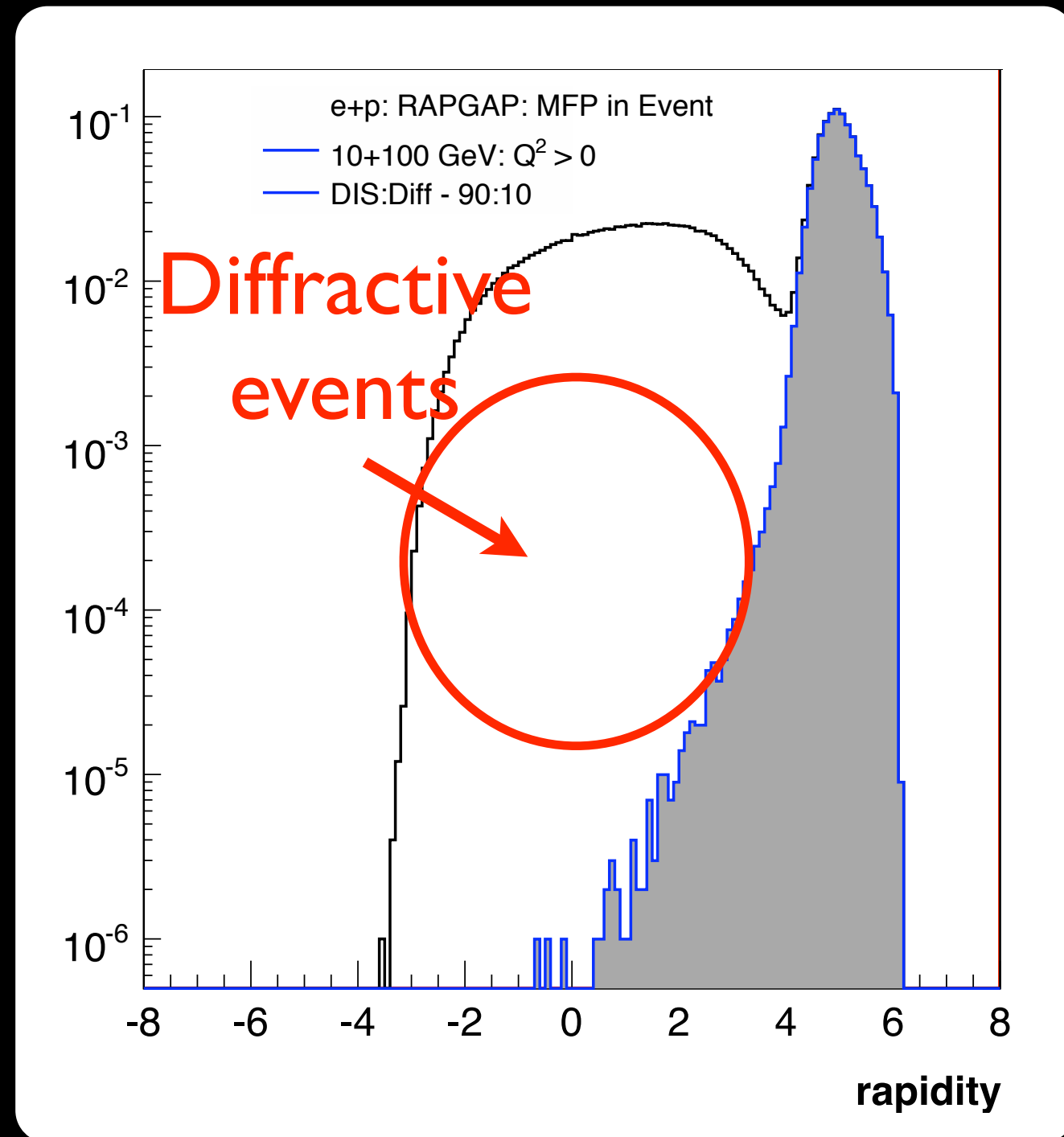
# Large rapidity gaps at an EIC

- Method:
  - Use RAPGAP in diffractive and DIS modes to simulate  $e+p$  collisions at EIC energies
  - Clear difference between DIS and Diffractive modes in “most forward particle in event” distributions
  - Little change in distributions with increasing energy



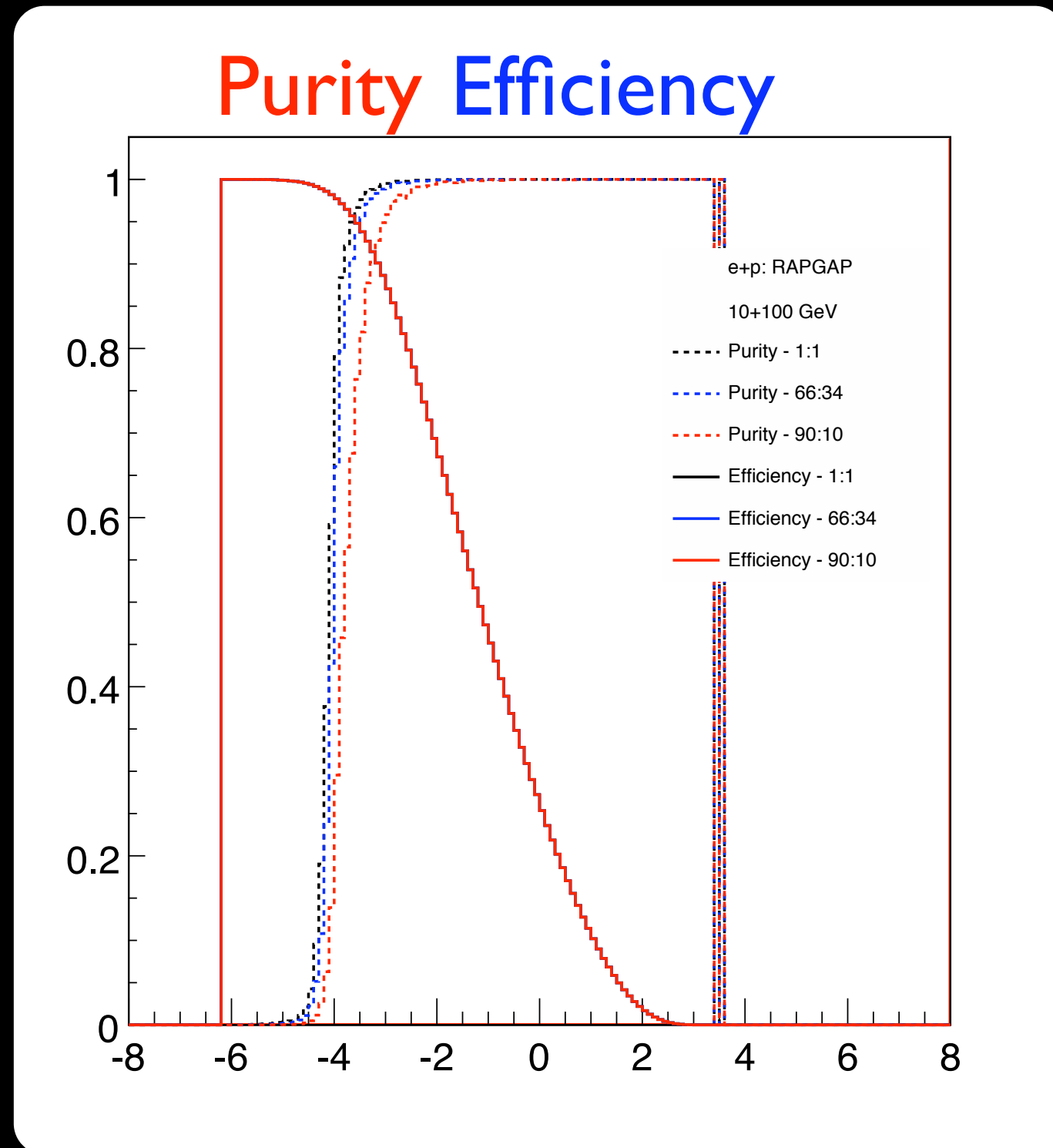
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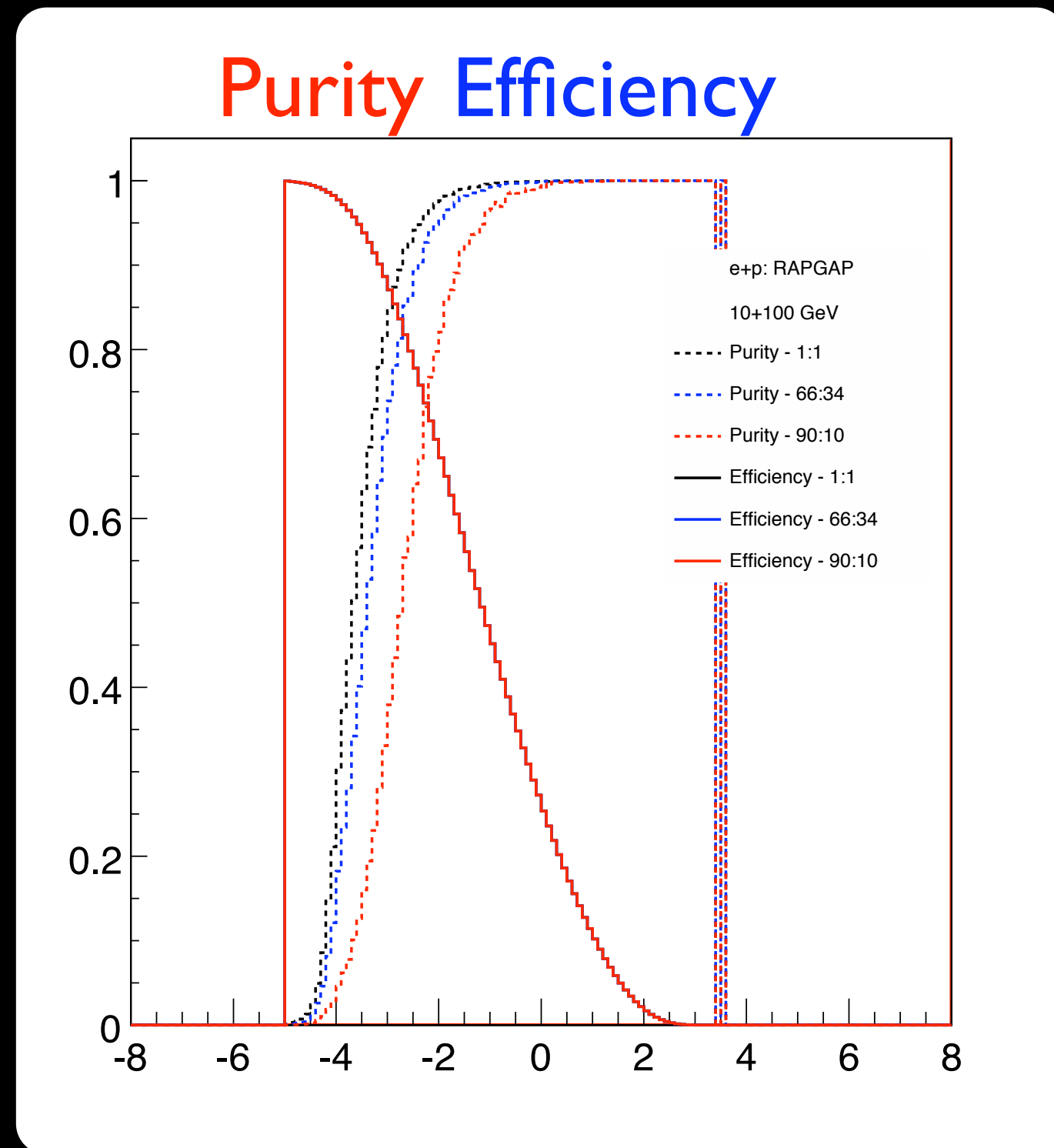
# Large rapidity gaps at an EIC

- Efficiency vs Purity:
  - Efficiency = fraction of diffractive events out of all diffractive events in sample
  - Purity = fraction of diffractive events out of all events in sample
  - Possible to place a cut to have both high efficiency and high purity
    - However, reduce the acceptance by 1 or 2 units of rapidity and these values drop significantly
  - Need hermetic detector coverage!!



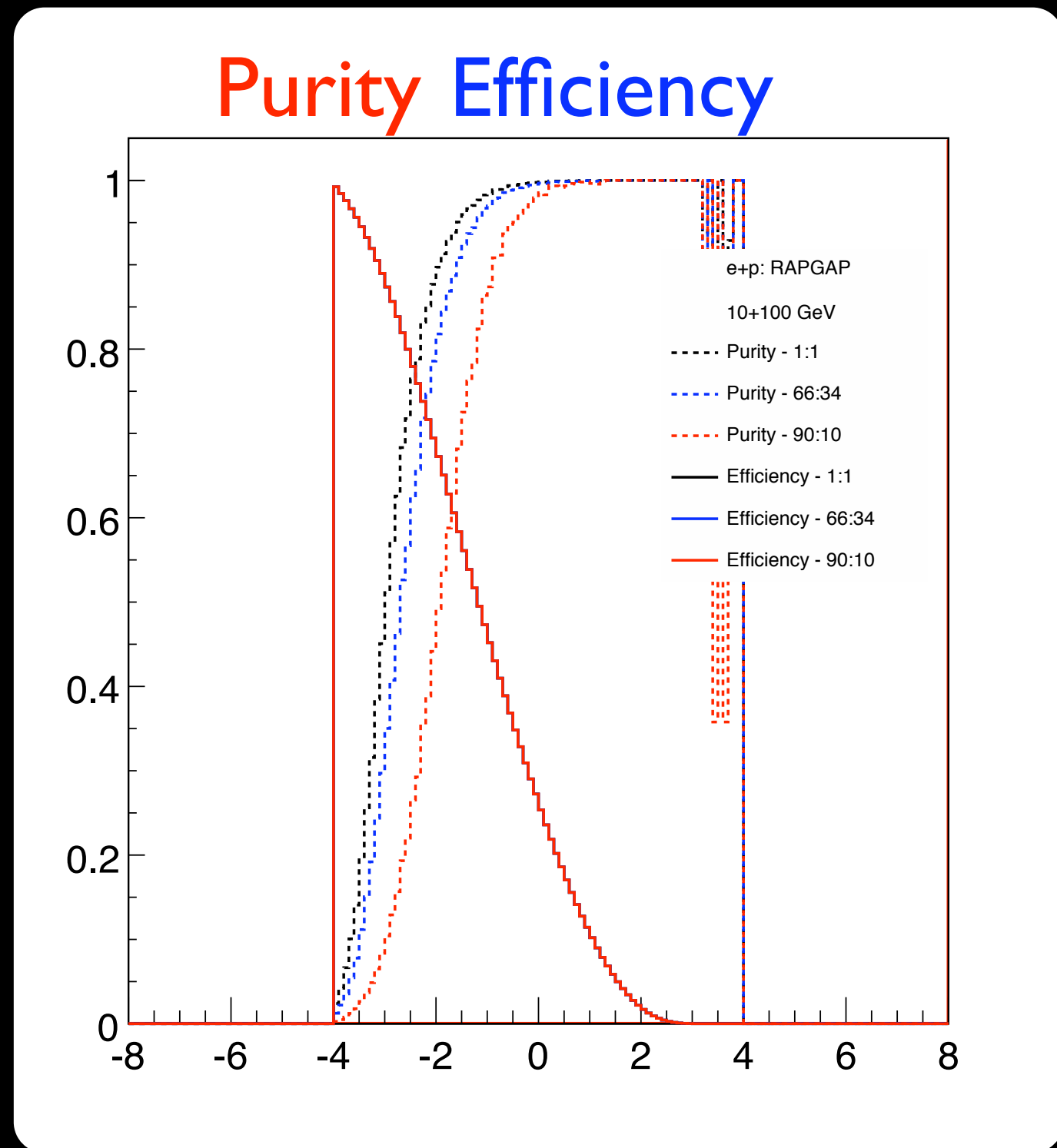
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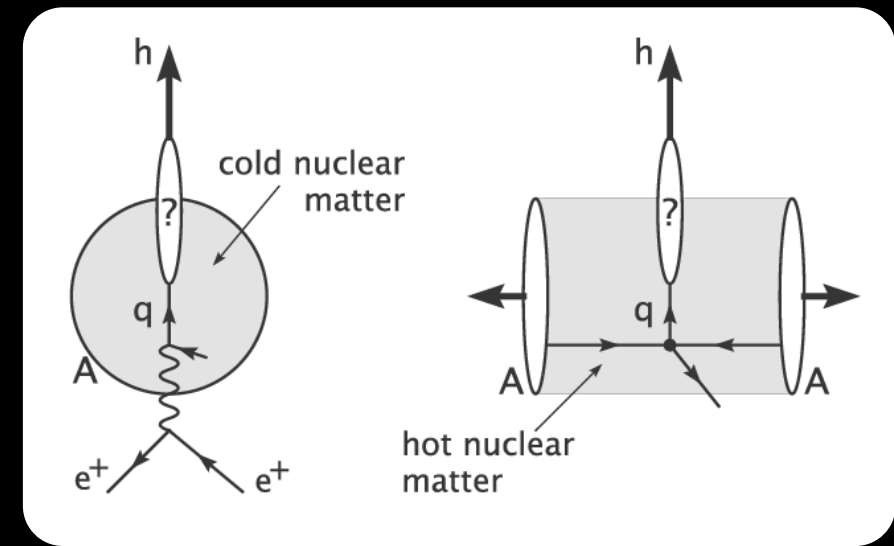


# 4 Key Measurements in e+A Physics

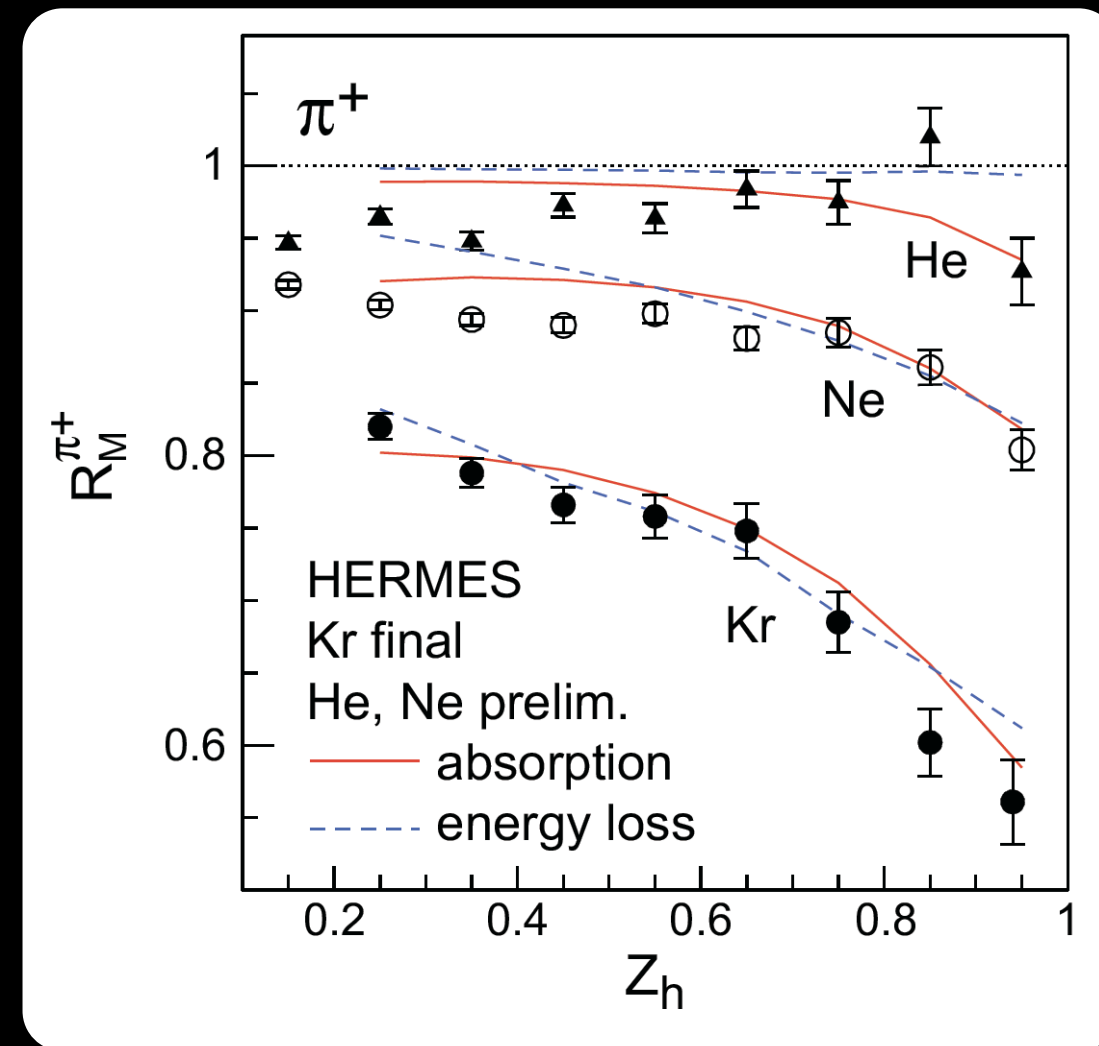
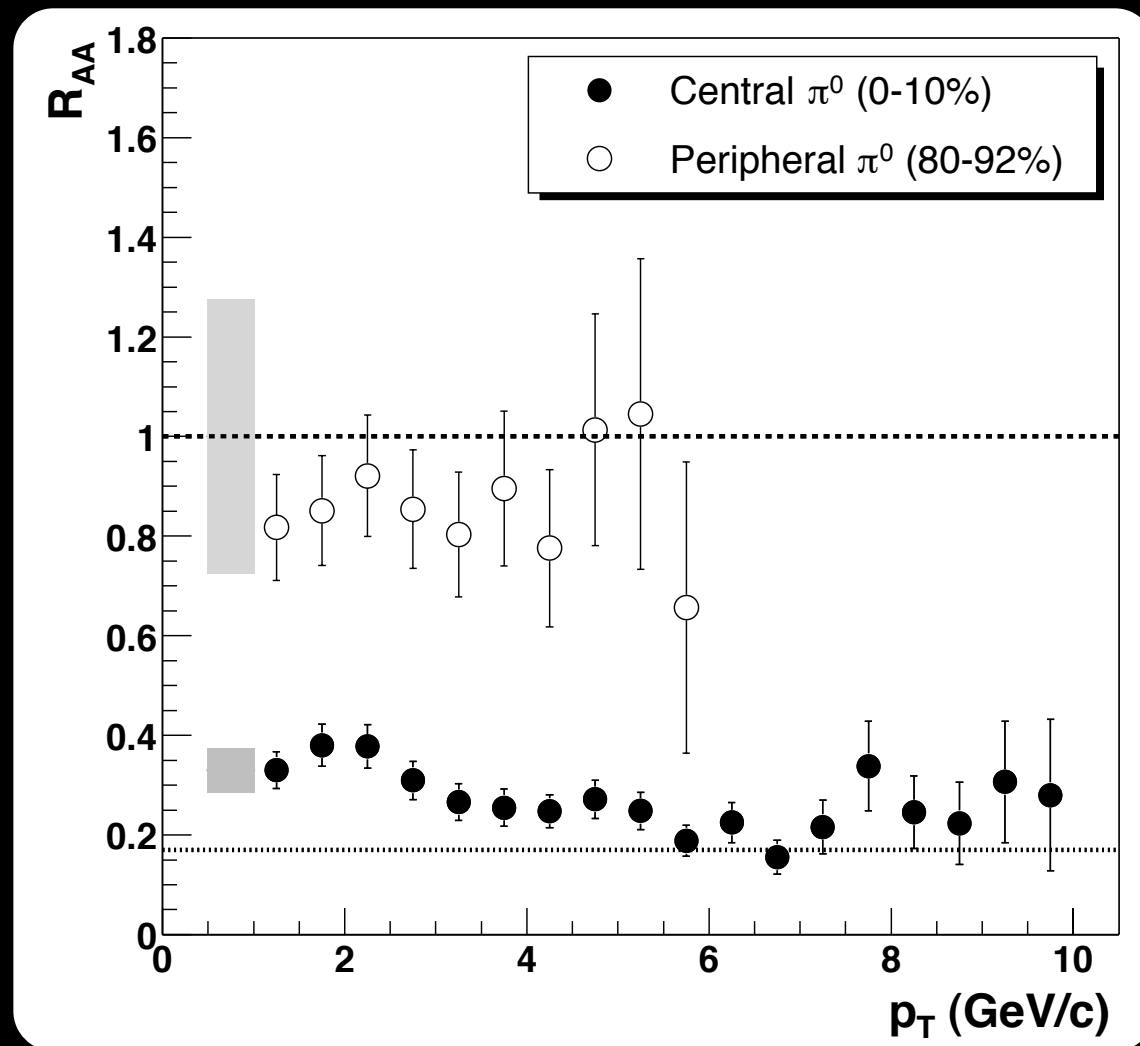
- Momentum distribution of gluons in nuclei?
  - Extract via scaling violation in  $F_2$ :  $\partial F_2 / \partial \ln Q^2$
  - Direct Measurement:  $F_L \sim xG(x, Q^2)$  - requires  $\sqrt{s}$  scan
  - Inelastic vector meson production (e.g.  $J/\Psi$ ,  $\rho$ )
  - Diffractive vector meson production ( $\sim [xG(x, Q^2)]^2$ )
- Space-time distribution of gluons in nuclei?
  - Exclusive final states (e.g.  $\rho$ ,  $J, \Psi$ )
  - Deep Virtual Compton Scattering (DVCS) -  $\sigma \sim A^{4/3}$
  - $F_2$ ,  $F_L$  for various impact parameters
- Role of colour-neutral (Pomeron) excitations?
  - Diffractive cross-section:  $\sigma_{\text{diff}}/\sigma_{\text{tot}}$  ( $\sim 10\%$ : HERA e+p; 30%? EIC e+A?)
  - Diffractive structure functions and vector meson productions
  - Abundance and distribution of rapidity gaps
- Interaction of fast probes with gluonic medium?
  - Hadronization, Fragmentation
  - Energy loss (charm!!)

# Interaction of fast probes with gluonic medium

- nDIS:
- Clean measurement in 'cold' nuclear matter
- Suppression of high- $p_T$  hadrons analogous to, but weaker than at RHIC



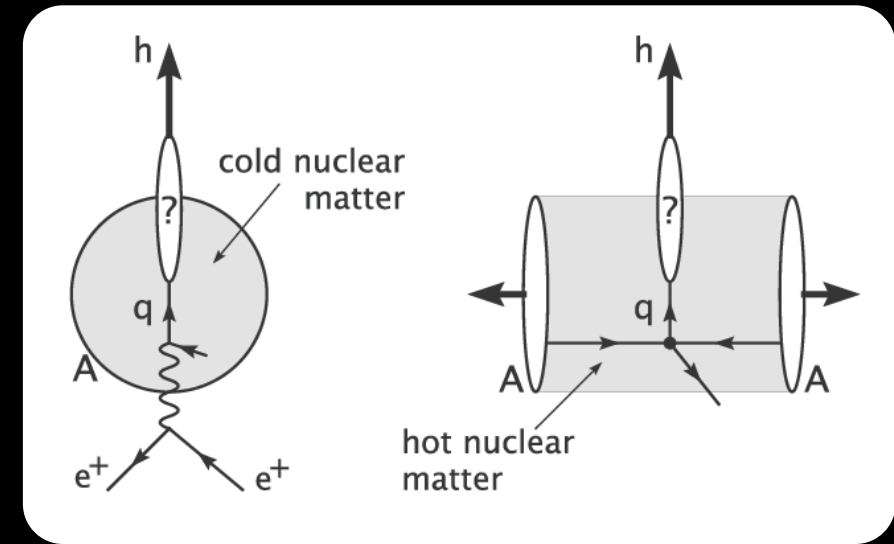
PHENIX expt: Phys.Rev.Lett.91:072301 (2003)





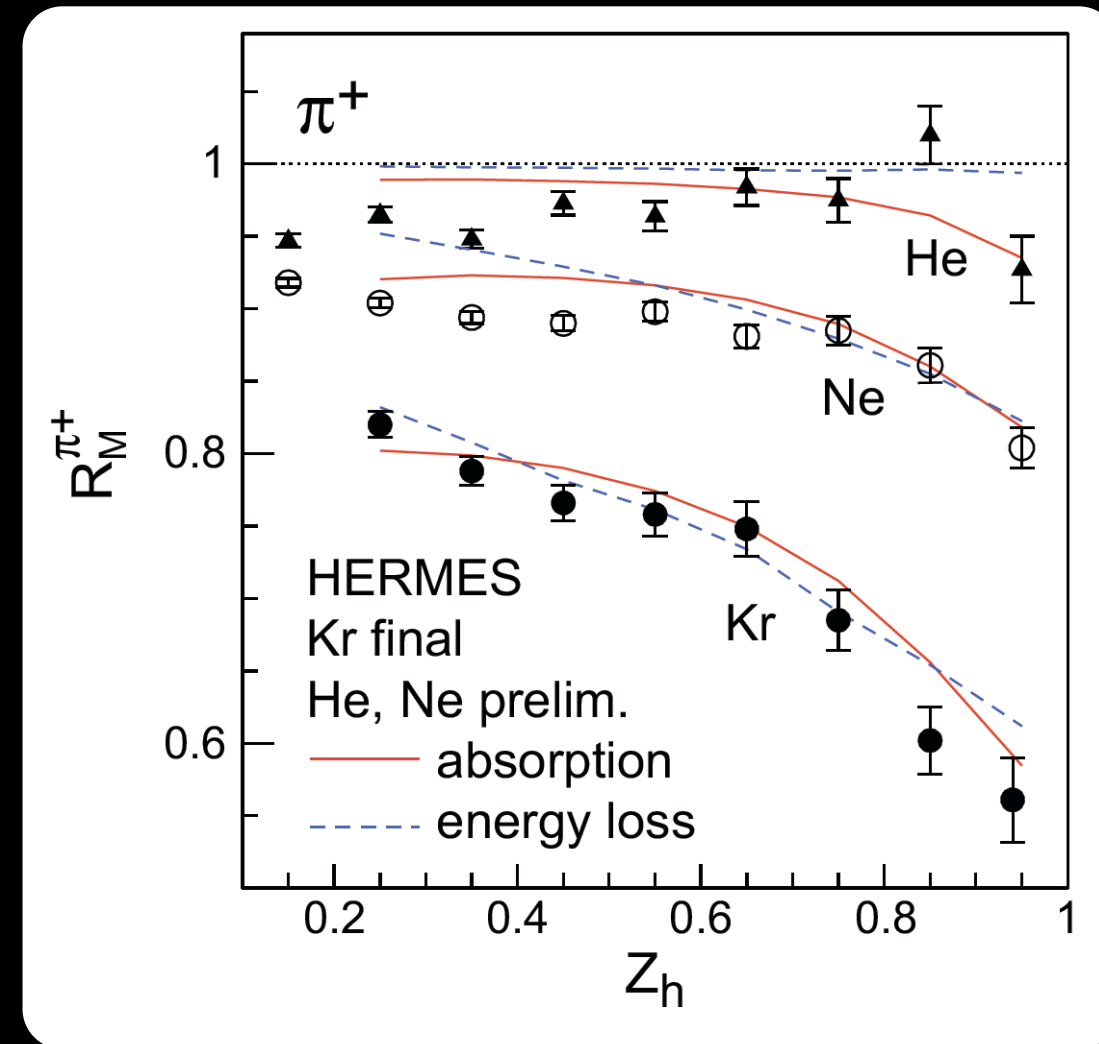
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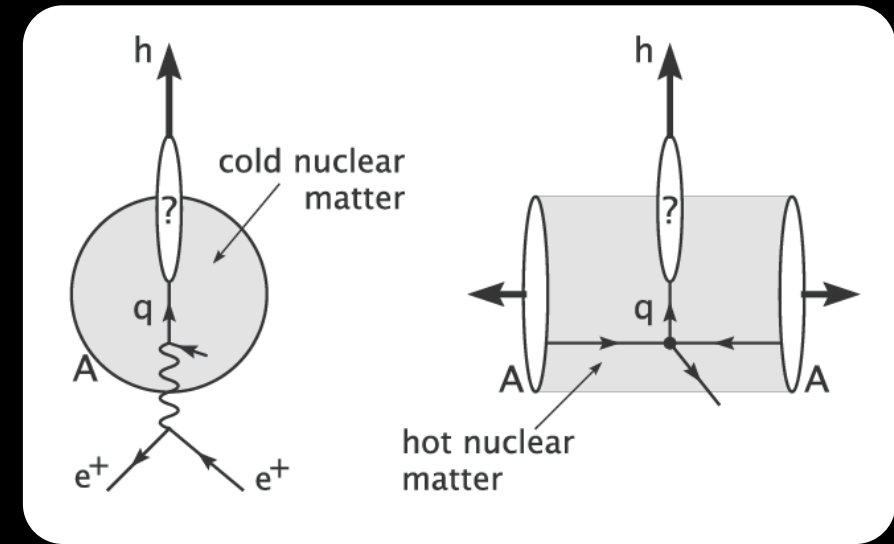
- Fundamental question:
  - When do partons get colour neutralized?

Parton energy loss vs. (pre)hadron absorption



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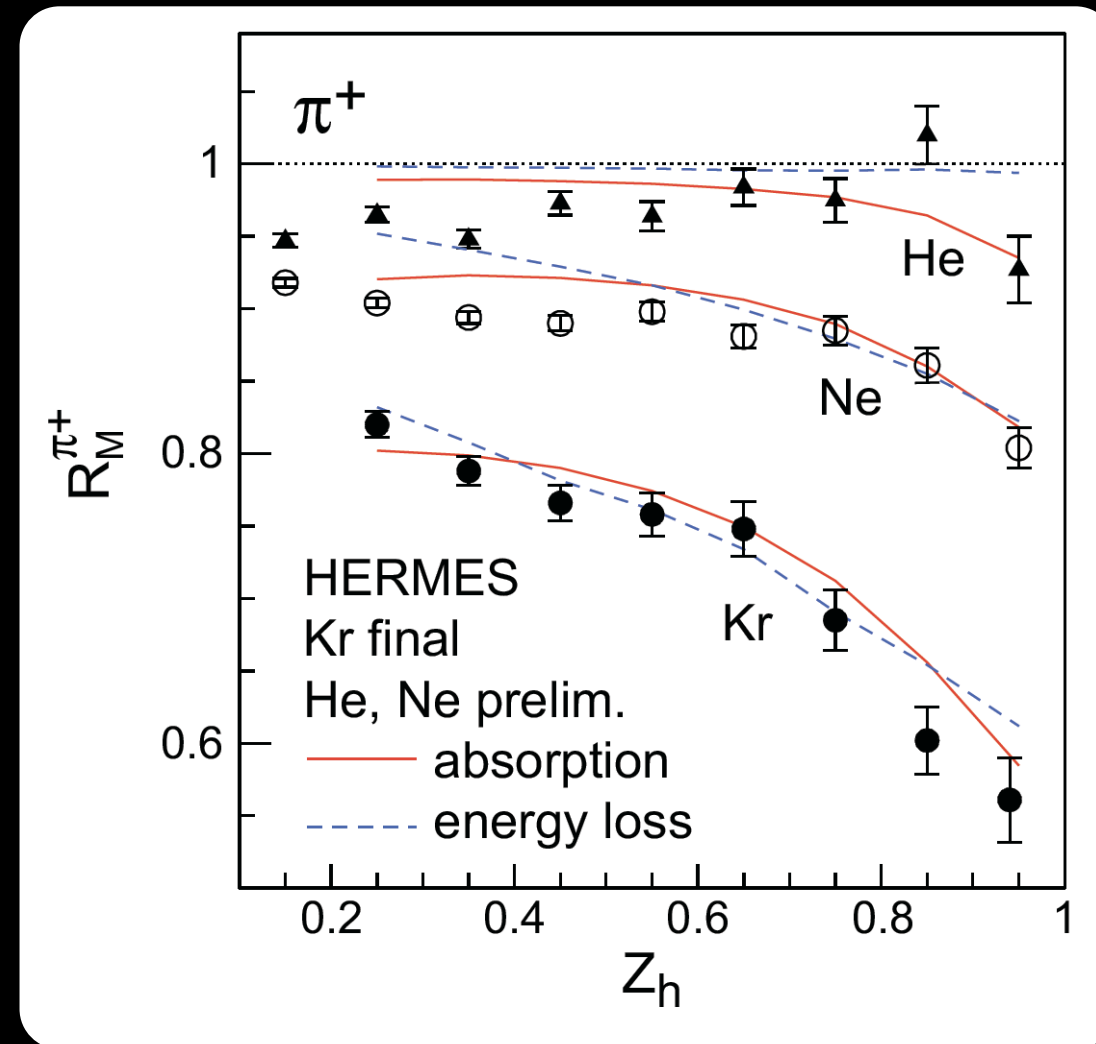
Parton energy loss vs. (pre)hadron absorption

Energy transfer in lab rest frame:

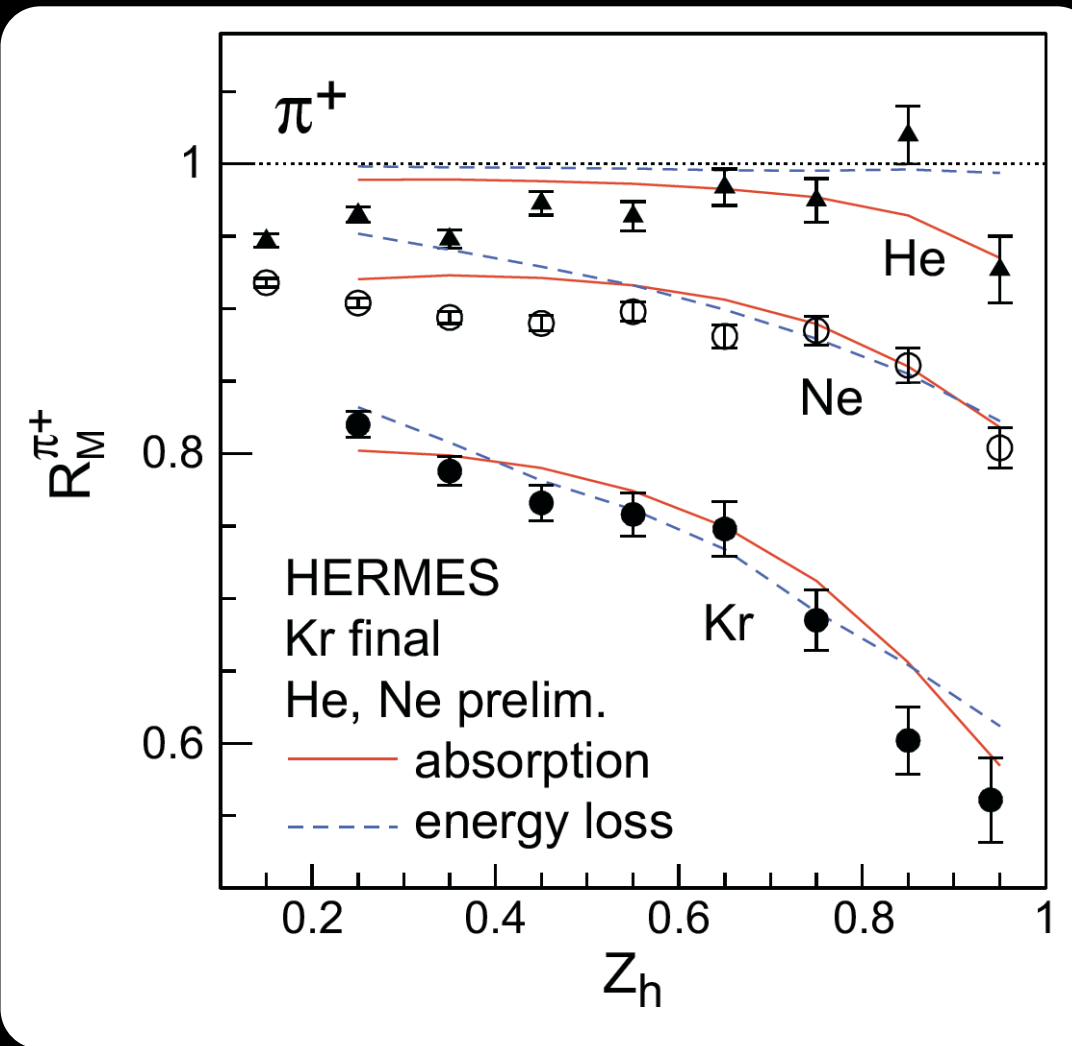
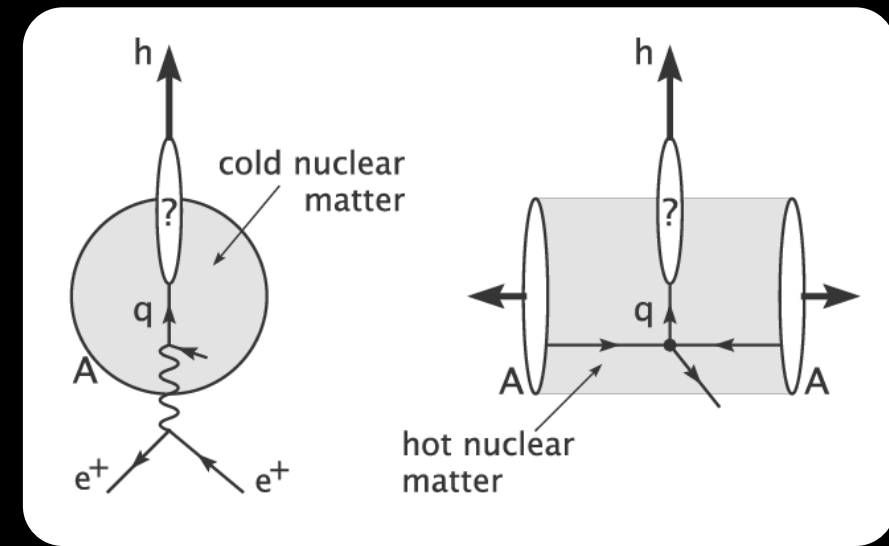
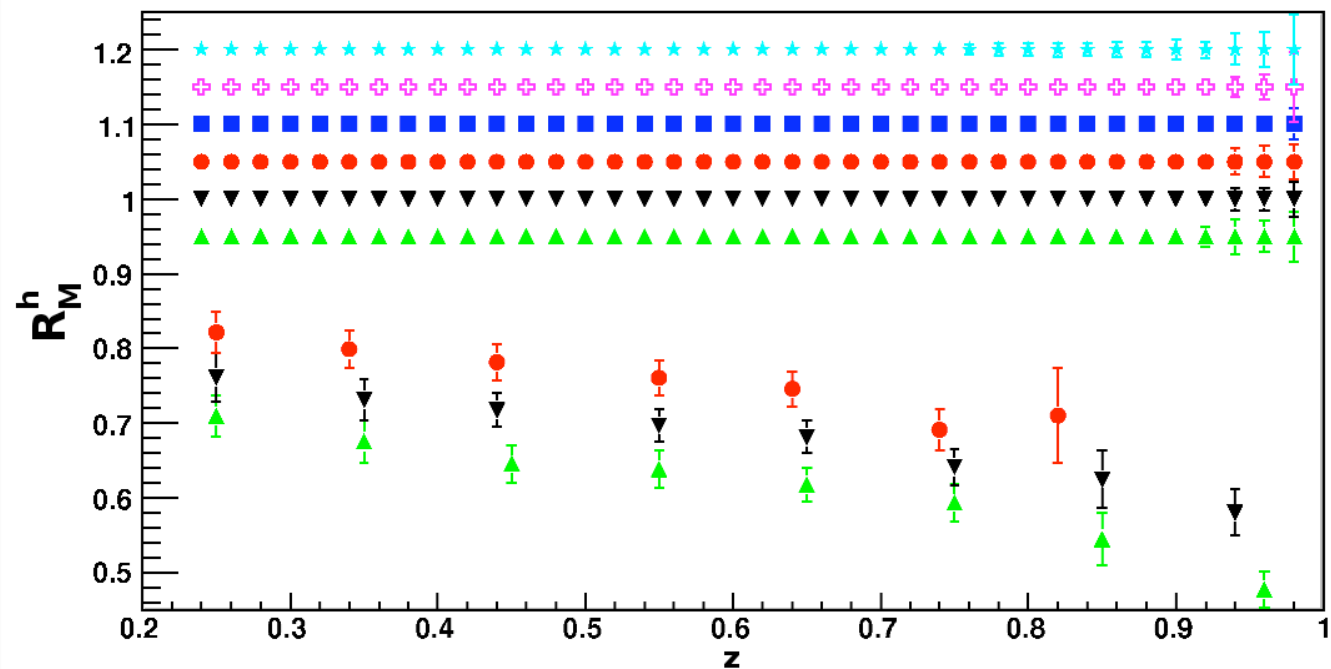
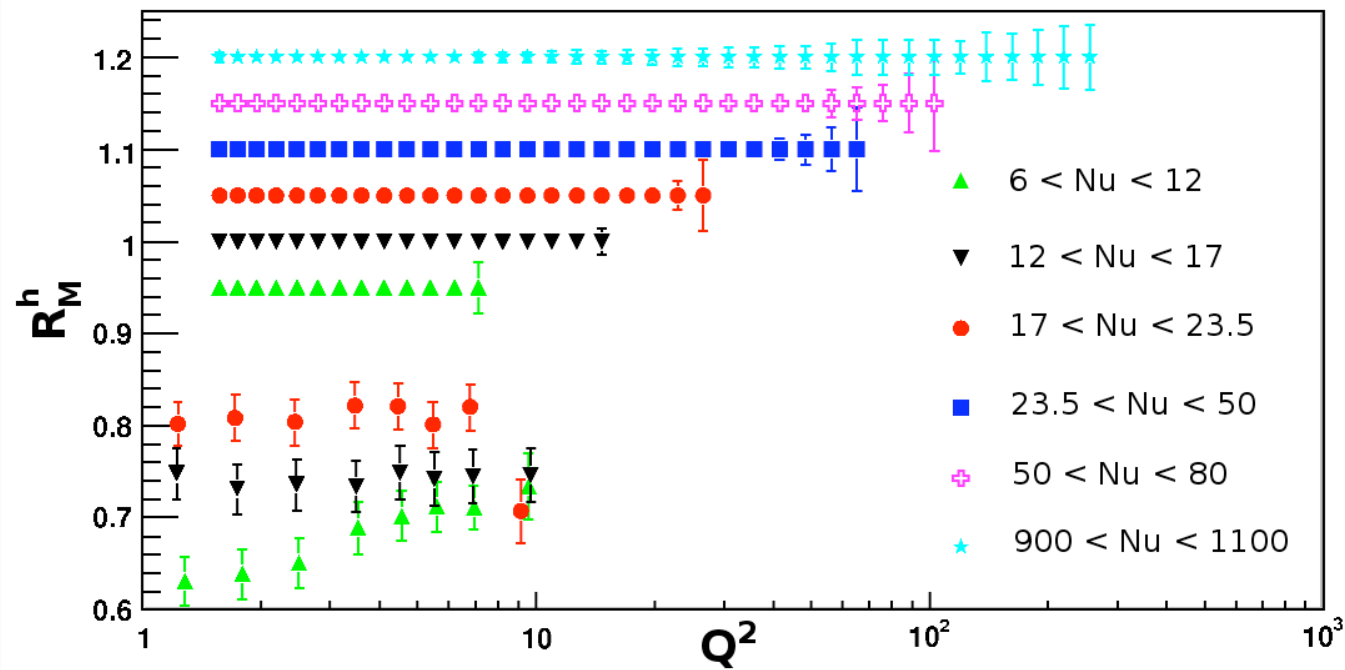
EIC:  $10 < \nu < 1600$  GeV

HERMES: 2-25 GeV

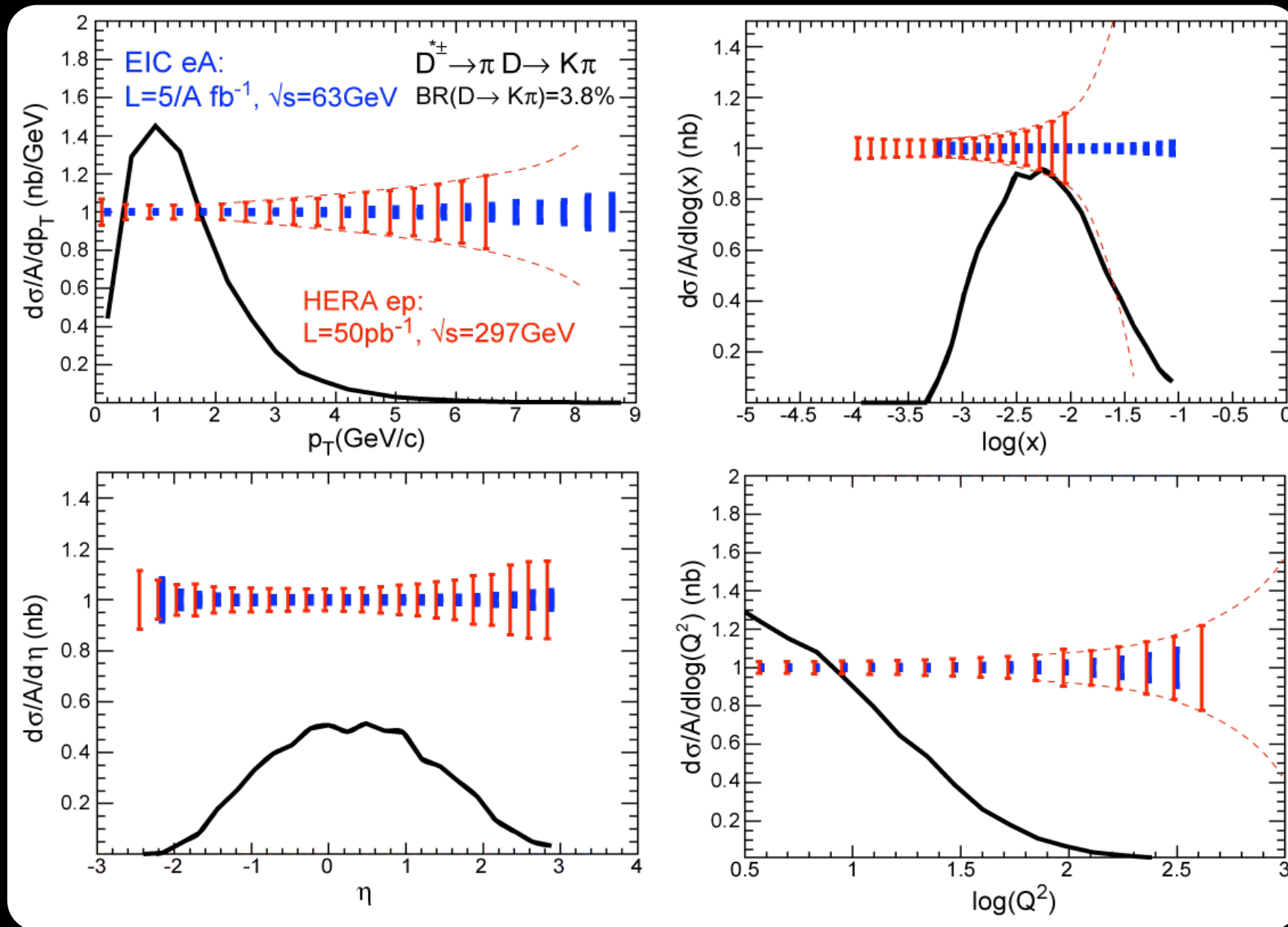
EIC: can measure *heavy flavour* energy loss



# Interaction of fast probes with gluonic medium

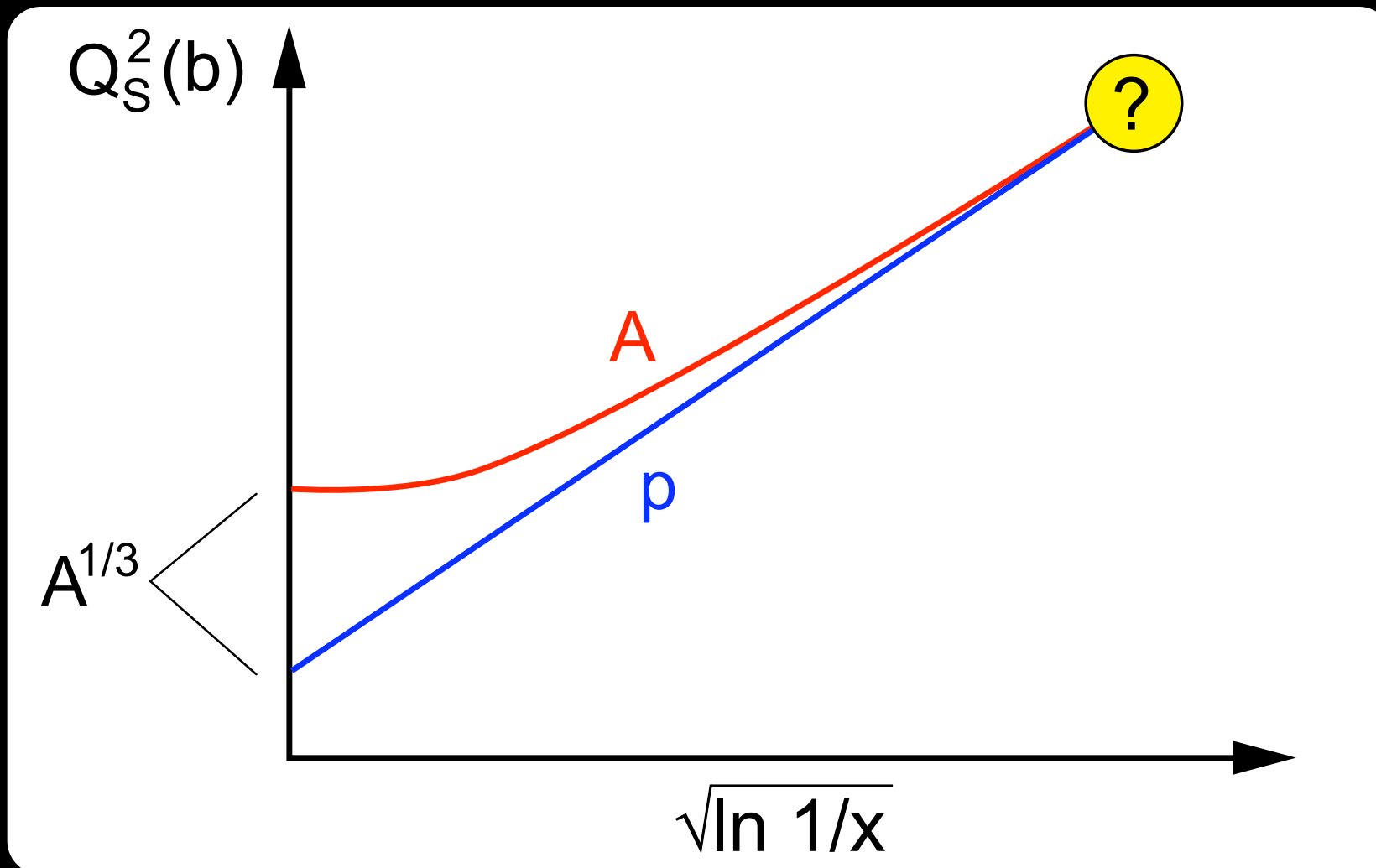


# Charm measurements at an EIC



- EIC: allows multi-differential measurements of heavy flavour
- covers and extends energy range of SLAC, EMC, HERA, and JLAB allowing for the study of wide range of formation lengths

# Matter at low- $x$ : A truly universal regime?



Small- $x$  QCD evolution predicts:

$Q_s$  approaches universal behavior for *all* hadrons and nuclei

⇒ Not only functional form  $f(Q_s)$  universal but even  $Q_s$  becomes the same

A.H. Mueller, hep-ph/0301109

- Nuclei and all hadrons have a component of their wave function with the *same* behaviour
- This is a conjecture! Needs to be tested

# Summary

- The study of  $e+A$  collisions at an EIC allow us to explore the physics of **Strong Colour Fields** and the nature of **non-linear QCD** and **saturation**. We can address the questions:
  - What are the momentum distributions of gluons in nuclei?
    - Measure  $F_2$ ,  $F_L$  distributions
  - What are the space-time distributions of gluons in nuclei?
    - Vector meson survival probability
  - What is the role of colour-neutral excitations (Pomerons)?
    - Diffractive physics
  - How do fast partons interact with cold nuclear matter?
    - Measure energy loss of fast-moving hadrons