

# Introduction to Inelastic Neutron Scattering

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With Thanks to:

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Roger Pynn - Indiana University

**HFIR**

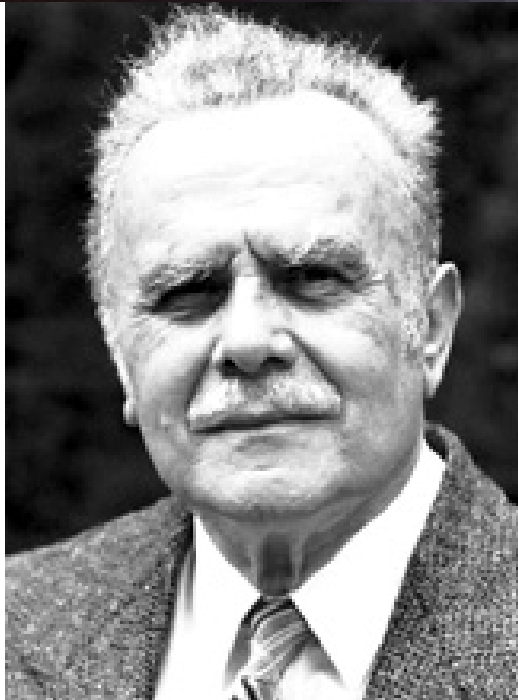


# The Nobel Prize in Physics 1994



"for pioneering contributions to the development of neutron scattering techniques for studies of condensed matter"

Neutrons show  
what atoms do



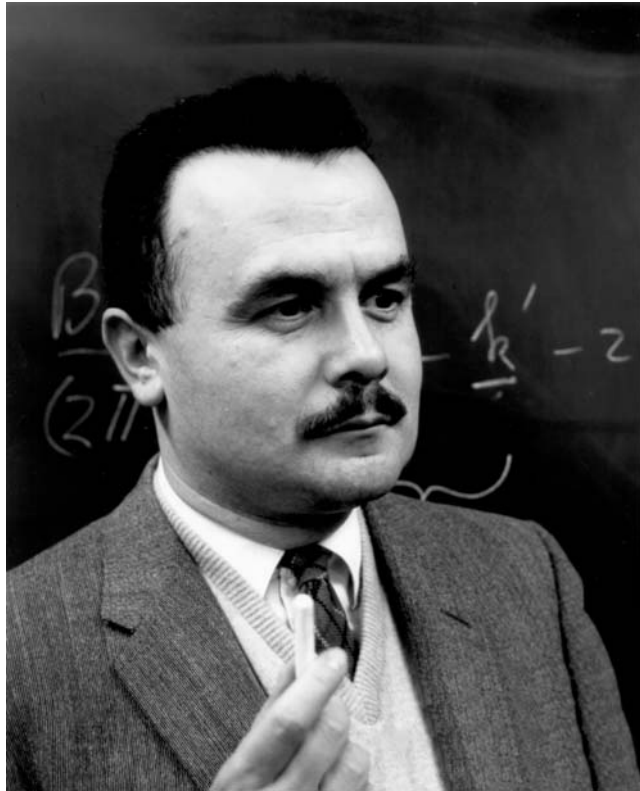
**Bertram N. Brockhouse**

Neutrons show  
where atoms are



**Clifford G. Shull**





*“If the neutron did not exist, it would need to be invented.”*  
- *B. Brockhouse*

# Why is the Neutron a Useful Particle?

- $m = 1.675\text{E-}27$  kg, no charge,  $S = \frac{1}{2}\hbar$ ,  $\mu_n = -1.913\mu_N$
- Typical de Broglie wavelengths,  $\lambda = 2\pi/k$ , are similar to lattice spacings.
- Energies of neutrons are tunable to that of excitations in condensed matter.
- Neutrons probe the entire sample, even inside complex sample environments.
- Neutrons interact with nuclei through nuclear force - lattice vibrations.
- Neutrons interact with unpaired electrons - magnetic excitations.



## Useful Formulae (non-relativistic neutrons)

$$p = mv = \hbar k = \frac{h}{\lambda}$$

$$E = \frac{1}{2}mv^2 = \frac{\hbar^2 k^2}{2m} = \frac{h^2}{\lambda^2 m}$$

$$E = 2.072k^2 = \frac{81.81}{\lambda^2}$$

Units:  $E$  (meV)    $k$  ( $\text{\AA}^{-1}$ )    $\lambda$  ( $\text{\AA}$ )



# Inelastic Scattering: scattering with energy and momentum transfer

Source

$E_i, \mathbf{k}_i$

Sample

$2\theta$

$E_f, \mathbf{k}_f$

$d\Omega$

Detector

Measure the inelastic (double differential) cross section:

$$\frac{d^2\sigma}{d\Omega dE_f} = \frac{k_f}{k_i} S(Q, \varepsilon)$$

$S(Q, \varepsilon)$  scattering function - physics

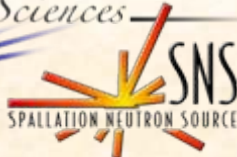
Conservation of energy

$$\hbar\omega = \varepsilon = E_i - E_f$$

Conservation of momentum

$$\vec{Q} = \vec{k}_i - \vec{k}_f$$

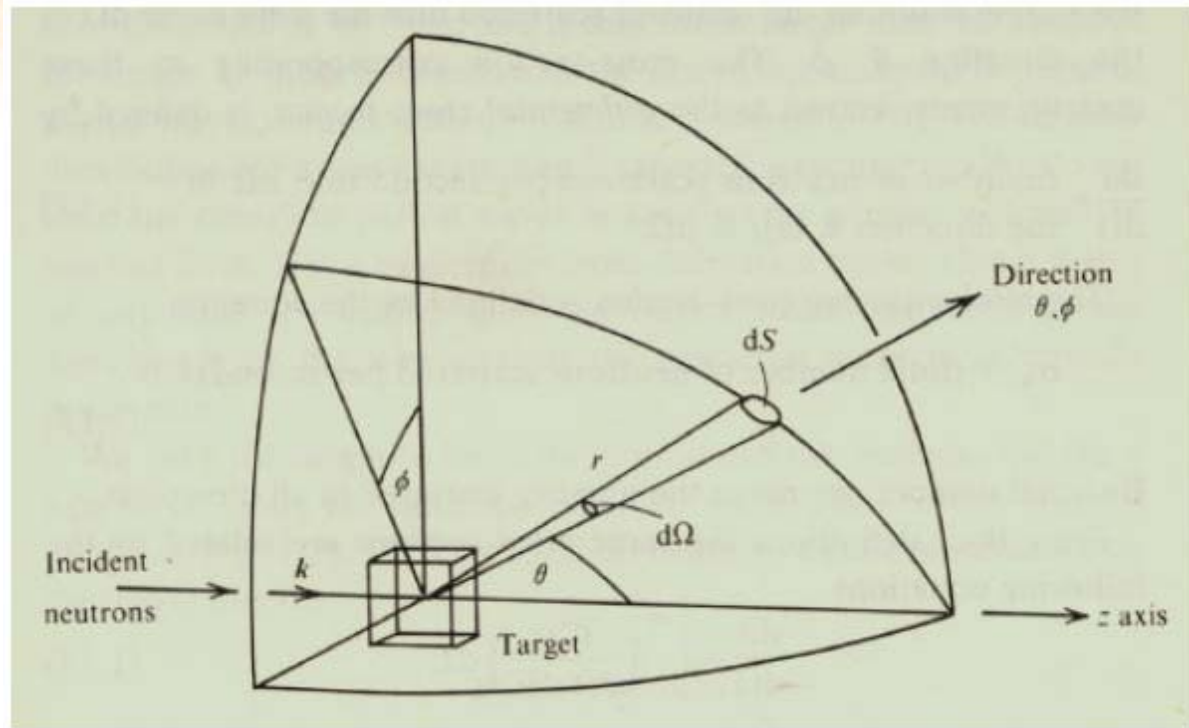
Neutron Sciences



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# Cross Sections



$\Phi$  = number of incident neutrons per  $\text{cm}^2$  per second

$\sigma$  = total number of neutrons scattered per second /  $\Phi$

$$\frac{d\sigma}{d\Omega} = \frac{\text{number of neutrons scattered per second into } d\Omega}{\Phi d\Omega}$$

$$\frac{d^2\sigma}{d\Omega dE} = \frac{\text{number of neutrons scattered per second into } d\Omega \text{ \& } dE}{\Phi d\Omega dE}$$

## Scattering Functions – $S(Q, \omega)$

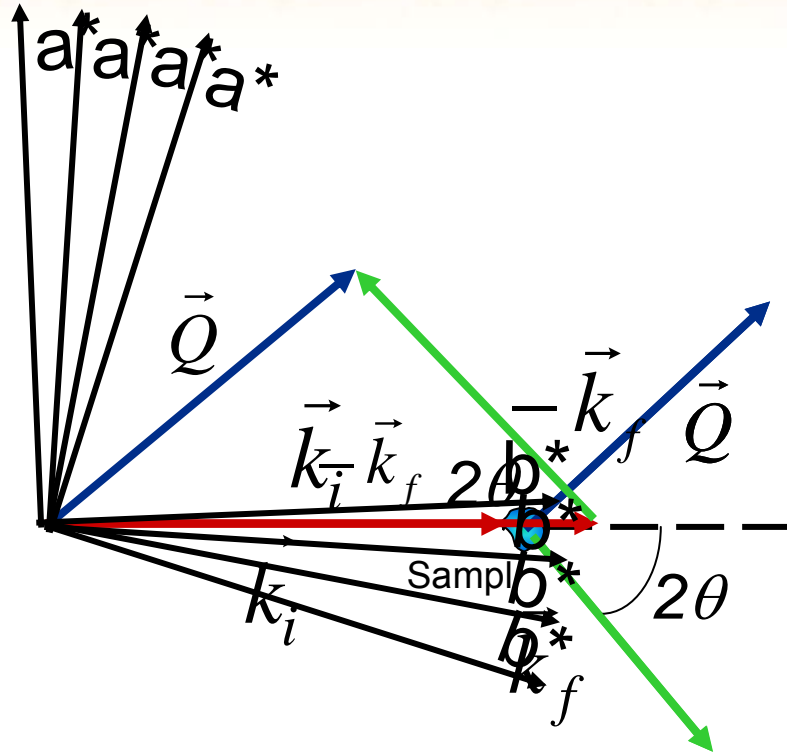
Expressions for  $S(Q, \omega)$  can be derived for a number of cases:

- Excitation or absorption of one quantum of lattice vibrational energy (phonon)
- Various models for atomic motions in liquids and glasses
- Various models of atomic & molecular translational & rotational diffusion
- Rotational tunneling of molecules
- Single particle motions at high momentum transfers
- Transitions between crystal field levels
- Magnons and other magnetic excitations such as spinons





# Reciprocal Space and the Scattering Triangle

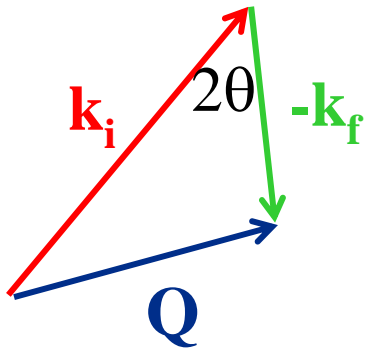


$$\vec{Q} = \vec{k}_i - \vec{k}_f$$

$$\hbar\omega = E_i - E_f$$

# Kinematic Constraints for Inelastic Scattering

Guideline I: one can not measure all Q and E



From scattering triangle and law of cosines

$$Q^2 = k_i^2 + k_f^2 - 2k_i k_f \cos(2\theta)$$

In terms of E and Q

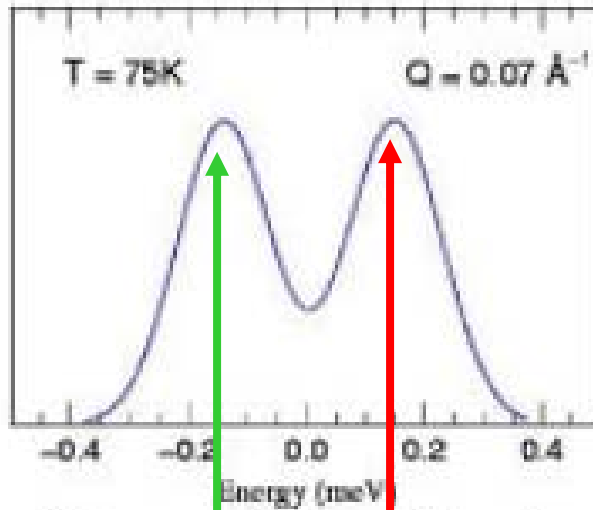
$$\frac{\hbar^2}{2m} Q^2 = E_i + E_f - 2\sqrt{E_i E_f} \cos(2\theta)$$

And substituting  $E_f = E_i - \varepsilon$

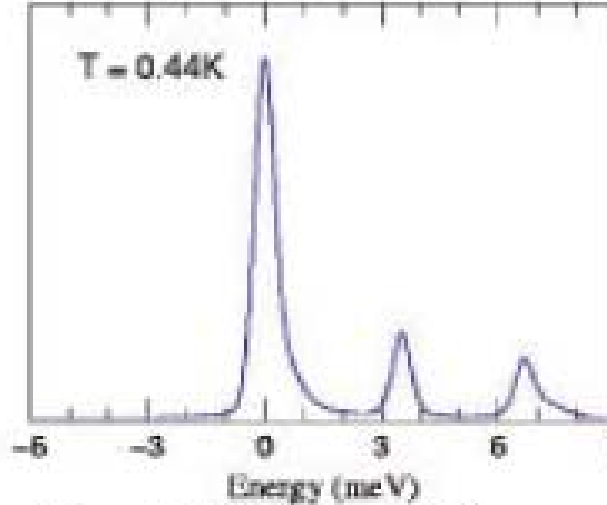
$$\frac{\hbar^2}{2m} Q^2 = 2E_i - \varepsilon - 2\sqrt{E_i(E_i - \varepsilon)} \cos(2\theta)$$

- For a given  $2\theta$  and  $E_i$ , one can only access a locus of Q and  $\varepsilon$  points. - limits your measurement

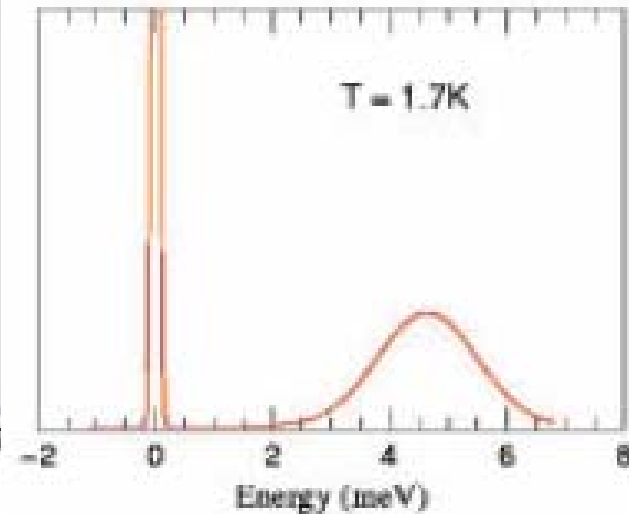
# Examples



Spin waves – collective excitations



Crystal Field splittings (HoPd<sub>2</sub>Sn) – local excitations

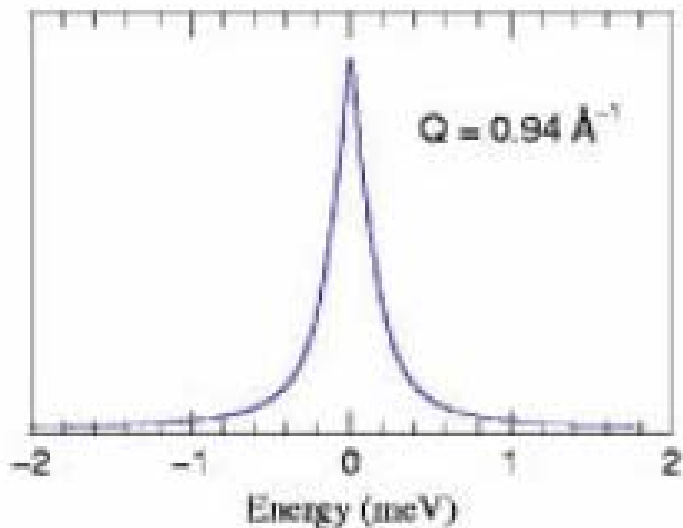


Local spin resonances (e.g. ZnCr<sub>2</sub>O<sub>4</sub>)

Neutron energy loss

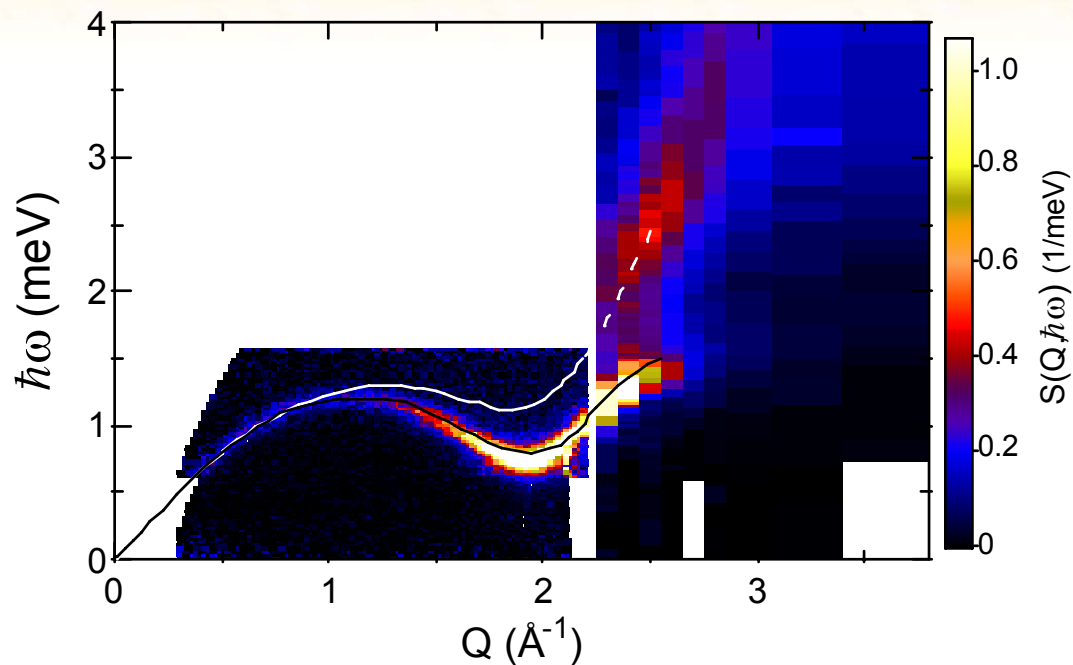
Neutron energy gain

## Examples II - Liquids



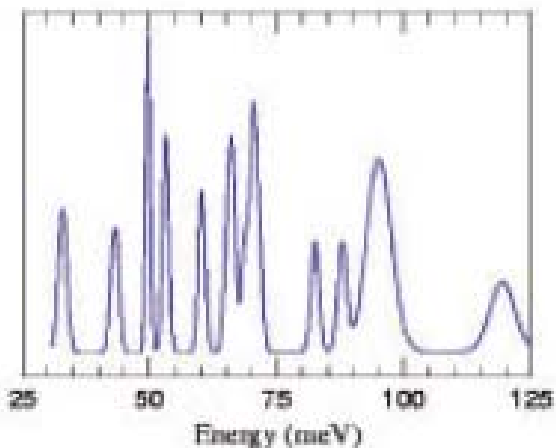
“Simple” liquids (e.g. water)

Diffuse scattering

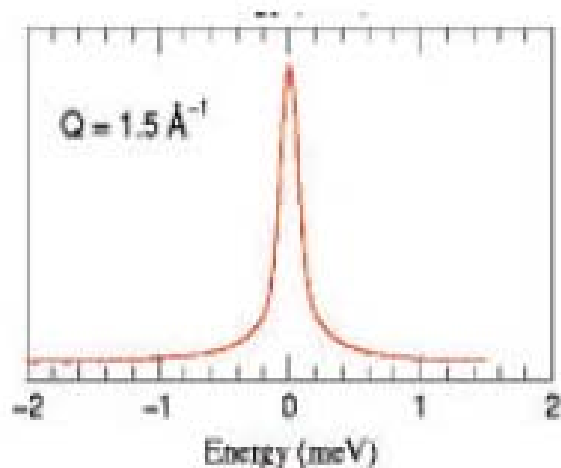


Quantum fluids  
Superfluid  $^4\text{He}$

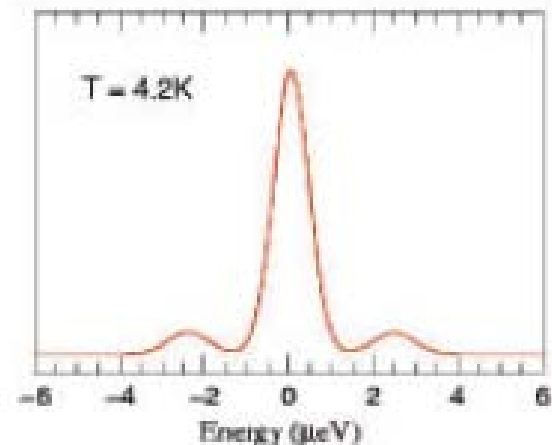
# Examples III: molecules ( $\mu\text{eV}$ to $\text{keV}$ )



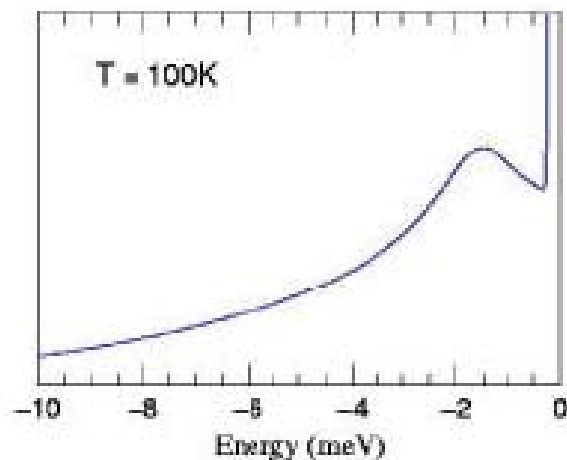
Vibrational spectroscopy  
(e.g.  $\text{C}_{60}$ )



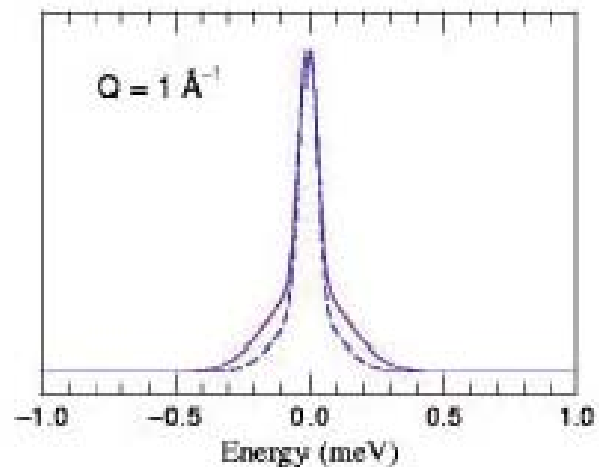
Molecular reorientation  
(e.g. pyrazine)



Rotational tunneling  
(e.g.  $\text{CH}_3\text{I}$ )



Polymers



Proteins



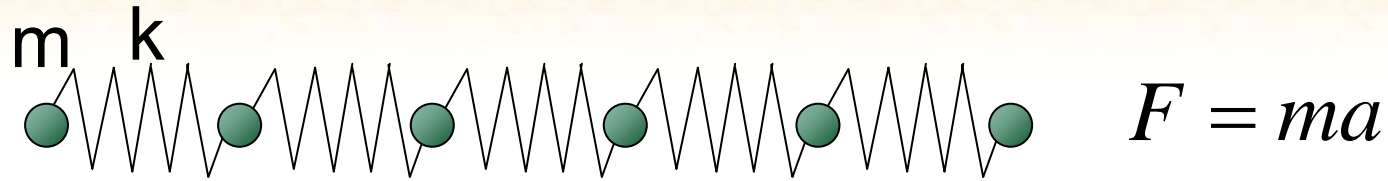
## Two short case studies

Lattice (phonon) excitations in  $\text{Ca}_{2-x}\text{Sr}_x\text{RuO}_4$   
and

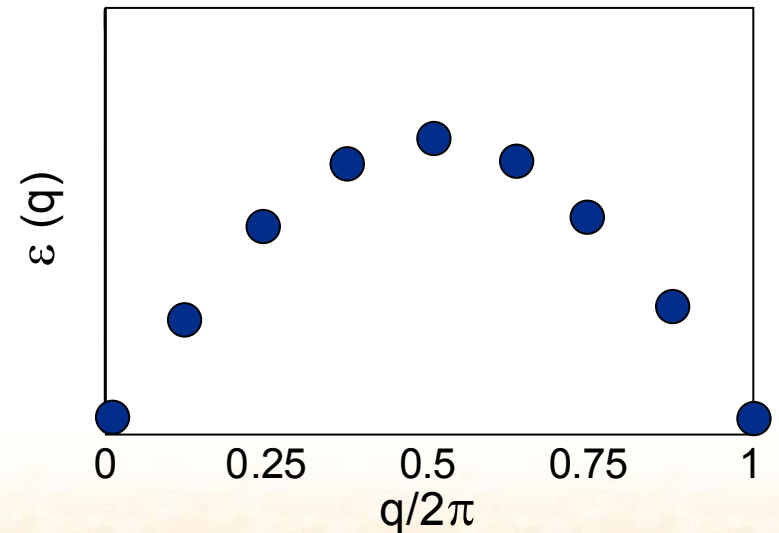
Magnetic (spin-wave) excitations in  $\text{CsMnCl}_3 \cdot 2(\text{D}_2\text{O})$



# Lattice Excitations in Condensed Matter

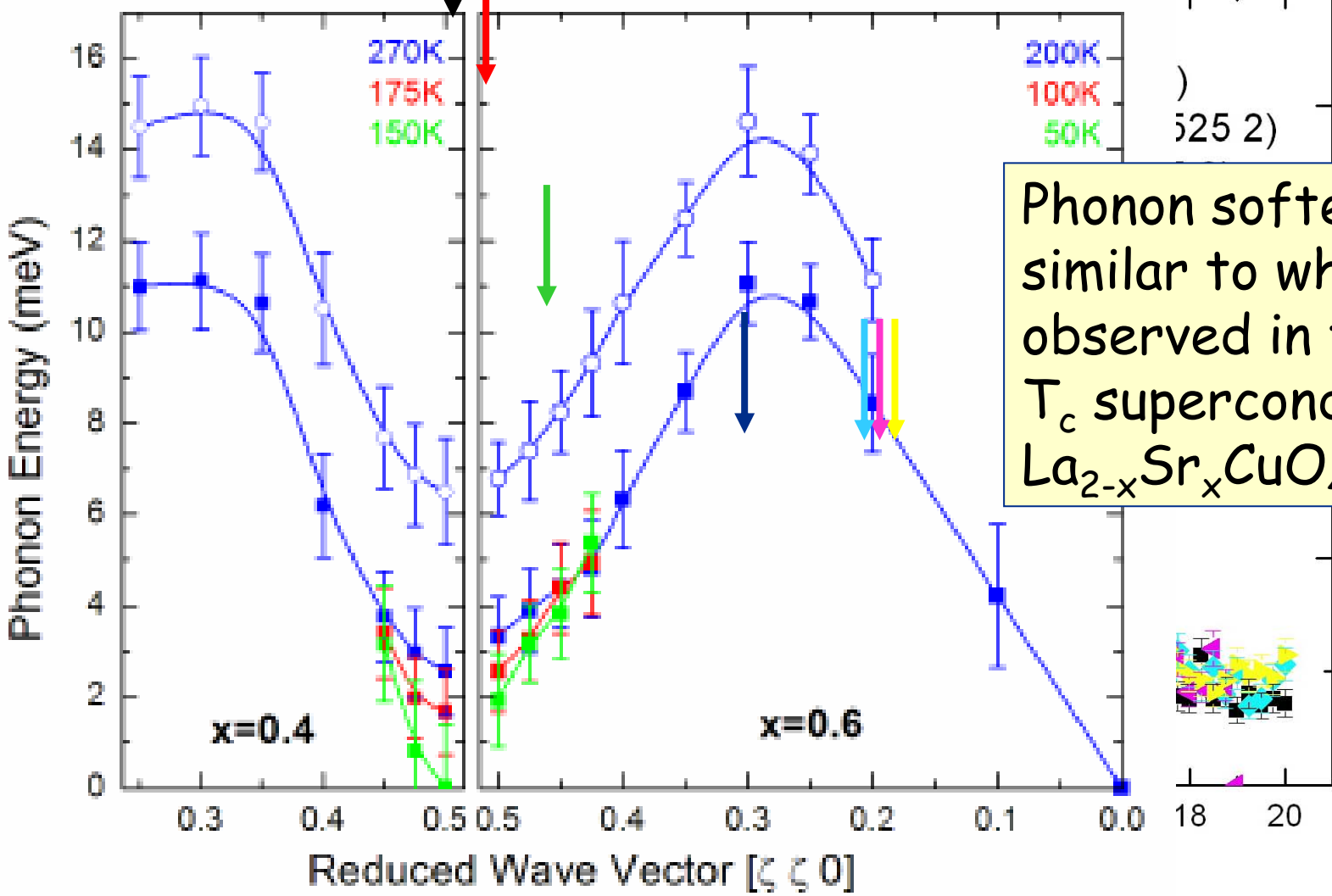


- Well defined excitations out of a long range ordered (LRO) state of atoms in crystal lattice
- Transmission of sound – Phonons
- Wave-vector dependent quasiparticles



# Phonons in $\text{Ca}_{2-x}\text{Sr}_x\text{RuO}_4$

R. G. Moore, M. D. Lumsden, M. B. Stone, J. Zhang, Y. Chen, J. W. Lynn, R. Jin, D. Mandrus, and E. W. Plummer



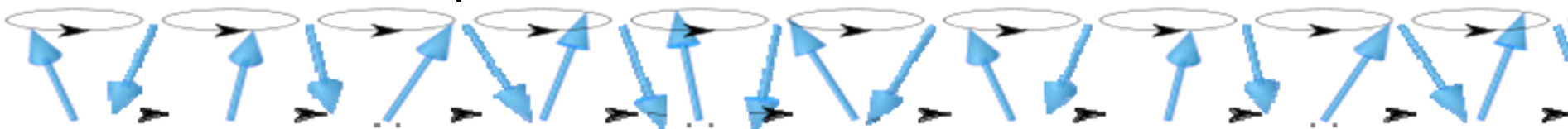
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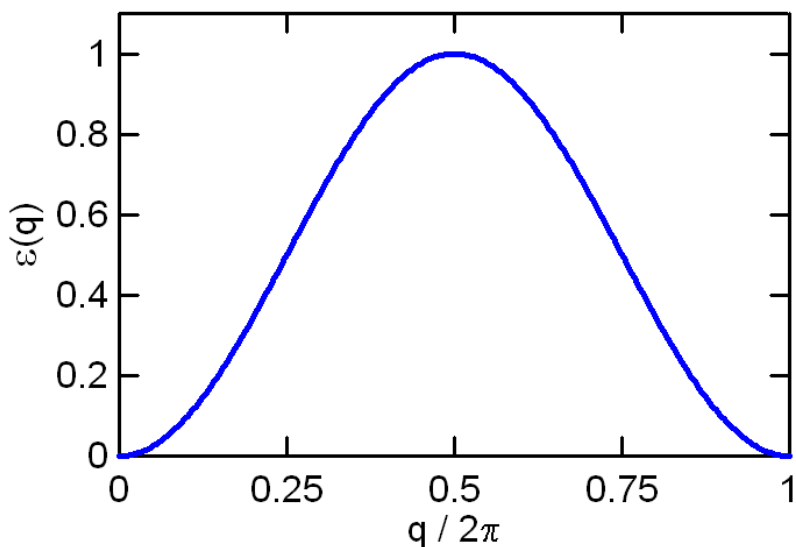
# Magnetic Excitations in an Ordered Magnet

$$H = \sum J(\mathbf{S} \cdot \mathbf{S})$$

- Exchange,  $J$ , in Hamiltonian sets energy scale
- ‘Spin-Solid’ – Long range ordered ground state
- Semi-classical spin-wave excitations – a doublet

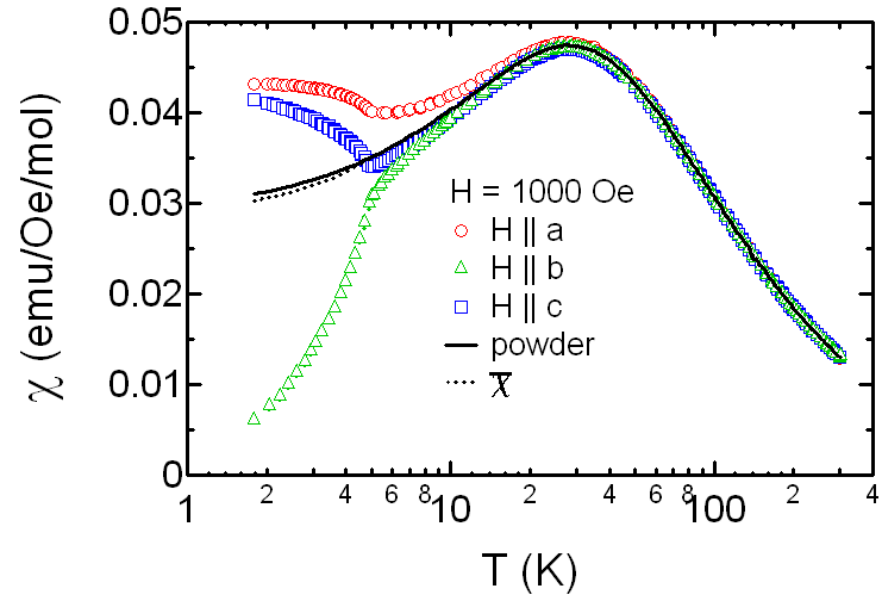
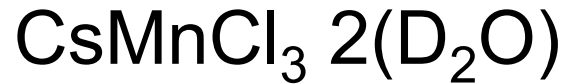


Antiferromagnetic  
Spin Waves



• Wave-vector  
dependent  
quasiparticles

# A Classic Antiferromagnet

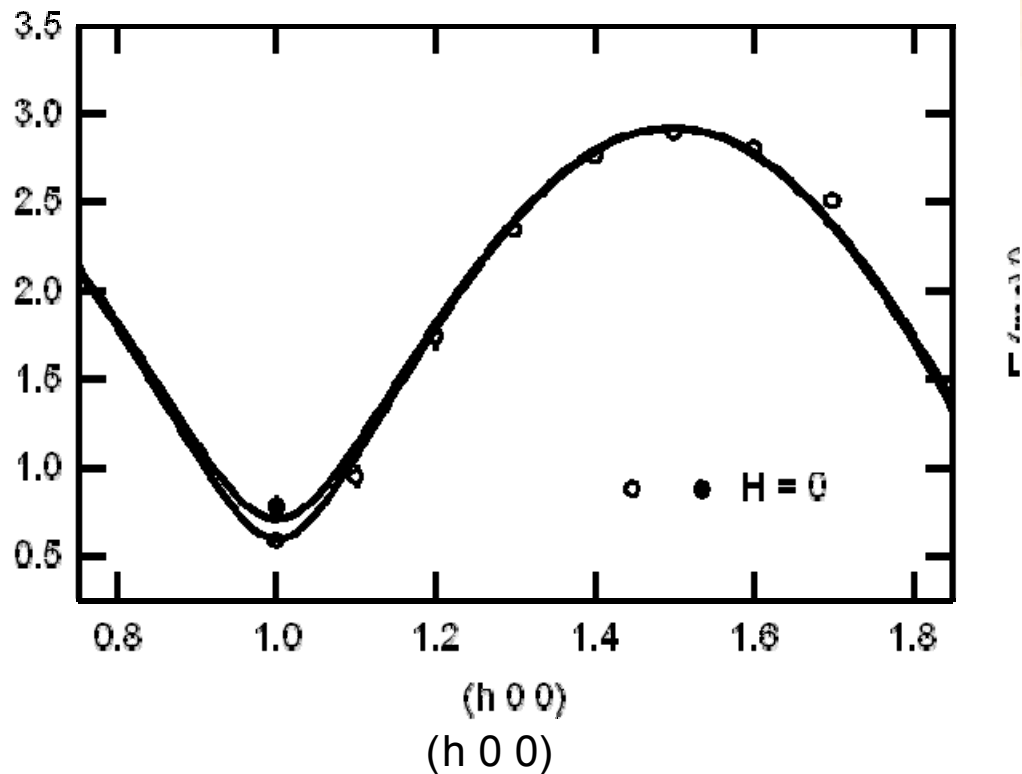
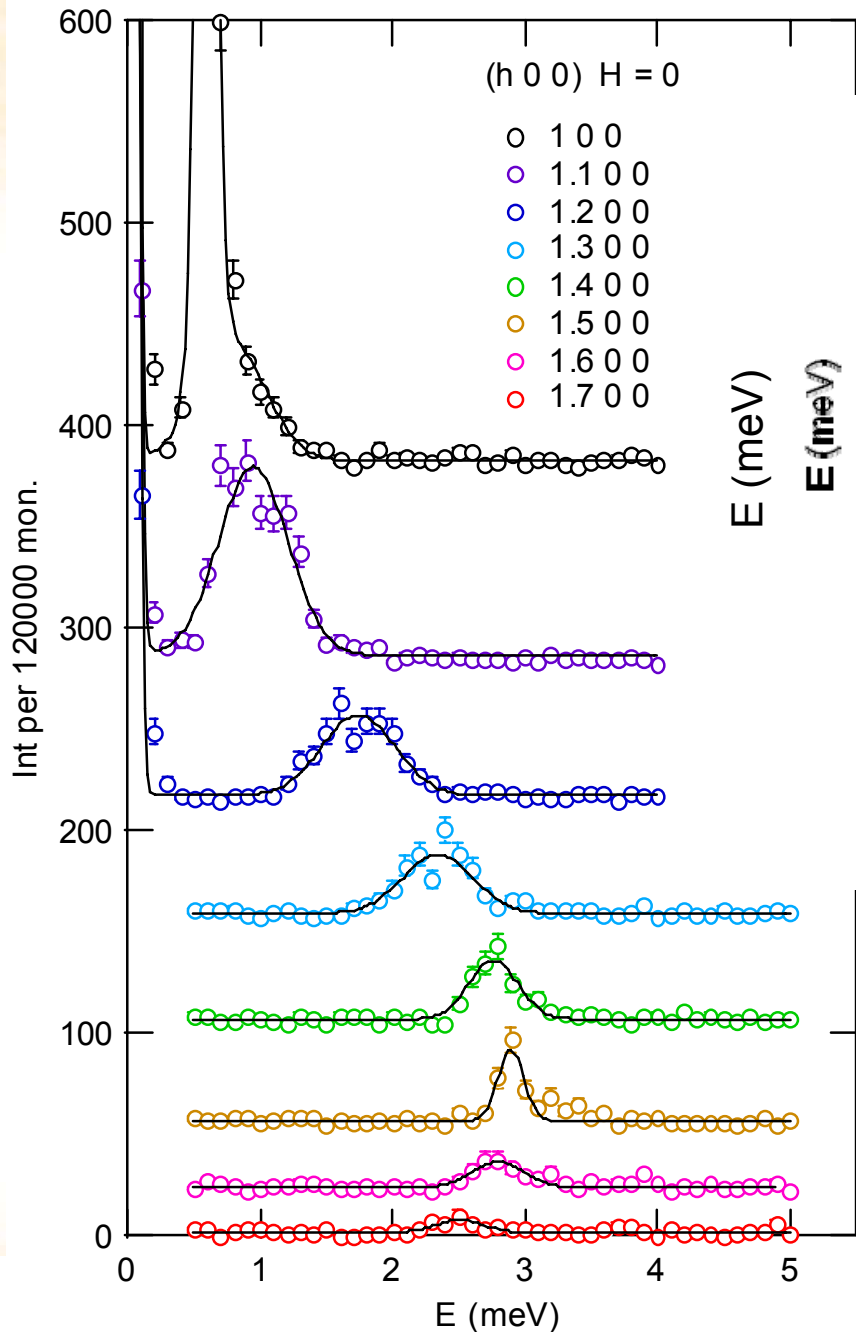


- $\text{Mn}^{2+}$ ,  $S=5/2 \approx \infty$ , semiclassical
- Gas of spins at high T  
Curie-Weiss law,  $\Theta = -48$  K
- Long range order "spin solid"  $T < T_N = 4.9$  K

*With A. Zheludev and I. Zaliznyak*

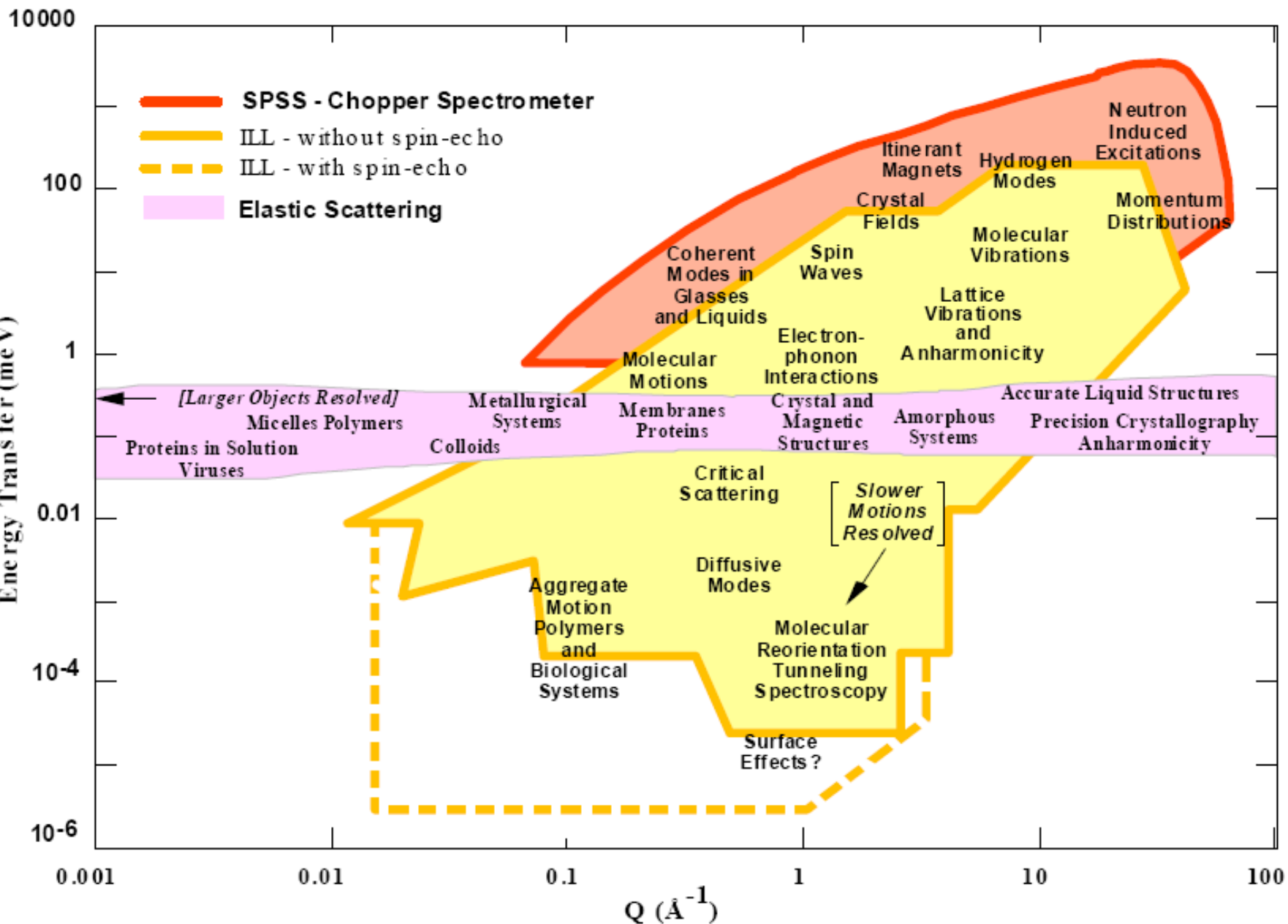


# CsMnCl<sub>3</sub> 2(D<sub>2</sub>O)



- $J_A = 0.58$  meV
- $J_A/J_B = 0.8\%$
- Semi-classical spin-waves
- $D \sim -0.007$  meV (anisotropy)

# Guideline II: There is no Universal Neutron Spectrometer



# Guideline IIb: Spallation versus Reactor Based Sources

Short Pulse Spallation Source	Reactor
Neutron spectrum is “slowing down” spectrum	Neutron spectrum is Maxwellian
<ul style="list-style-type: none"> <li>Choose your instrument based upon the physics being probed</li> </ul>	
<ul style="list-style-type: none"> <li>Instrument scientists can discuss possibilities and limitations prior to submission of beam-time proposal</li> </ul>	slow dynamics
Low background between pulses => good signal to noise	Pulse rate for TOF can be optimized independently for different spectrometers
Polarization possible, but difficult	Neutron polarization easier

## Guideline III: Neutron Flux is comparatively low

	<i>Brightness</i> ( $s^{-1} m^{-2} ster^{-1}$ )	<i>dE/E</i> (%)	<i>Divergence</i> ( $mrad^2$ )	<i>Flux</i> ( $s^{-1} m^{-2}$ )
Neutrons	$10^{15}$	2	10 x 10	$10^{11}$
Rotating Anode	$10^{16}$	3	0.5 x 10	$5 \times 10^{10}$
Bending Magnet	$10^{24}$	0.01	0.1 x 5	$5 \times 10^{17}$
Wiggler	$10^{26}$	0.01	0.1 x 1	$10^{19}$
Undulator (APS)	$10^{33}$	0.01	0.01 x 0.1	$10^{24}$

- Implies large single crystals or large quantities of powder.
- Calculate or Compare to prior measurements to determine appropriate sample size.
- Can use broad and tunable resolution to advantage!

# Closing thoughts

- **Inelastic neutron scattering allows for the measure of excitations over a range of energy and length scales.**
- **Need to choose the right instrument.**
- **Need appropriate samples (mass, powder, single crystal, isotopic purity, etc.).**
- **Often need to know what you are looking for.**
- **Don't hesitate to discuss experiments with instrument scientists prior to submission of proposal (feasibility, length of time, instrument considerations, etc.).**
- **Submit proposals to** <http://www.ornl.gov/sci/iuims/ipts/>  
**Or Google: ORNL neutron proposal**

