

Many Body Lattice QCD

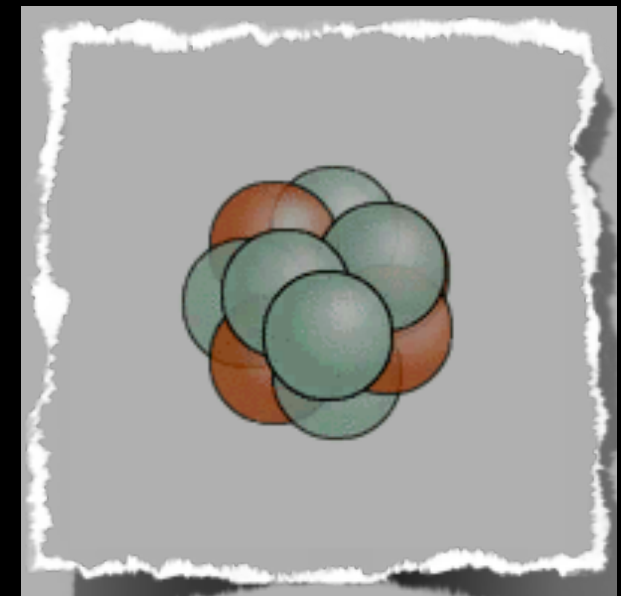
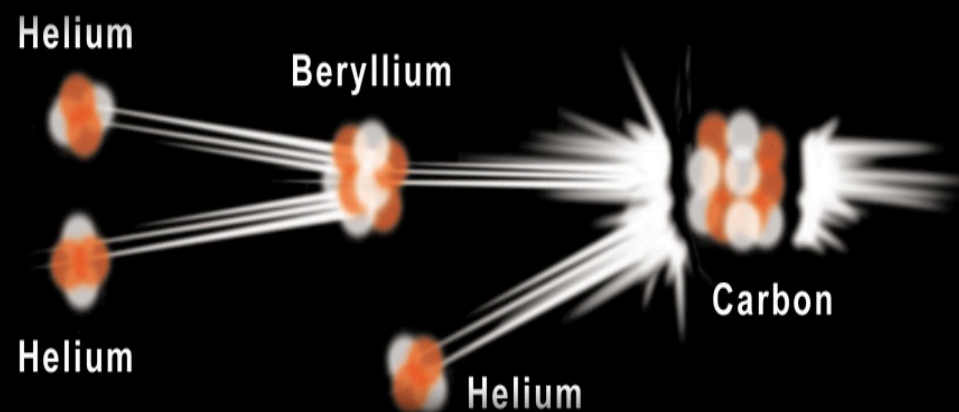
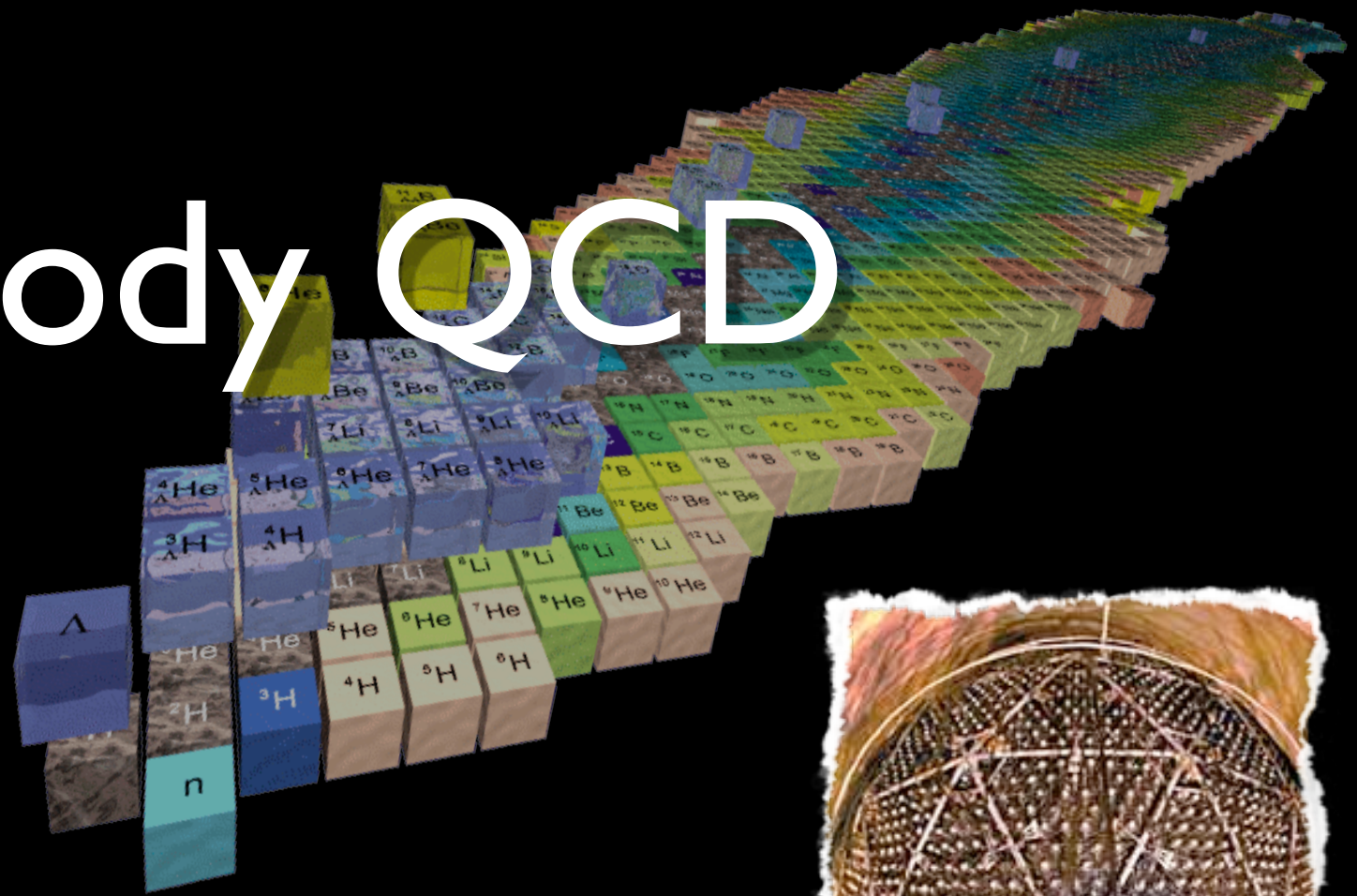
William Detmold

*The College of William and Mary &
Thomas Jefferson National Accelerator Facility*

- Multi-hadron systems in QCD
- $n < 13$ pions and kaons: meson condensates
 - Two and three body interactions
 - Pion and kaon condensates
 - In-medium screening of the $Q\bar{Q}$ potential

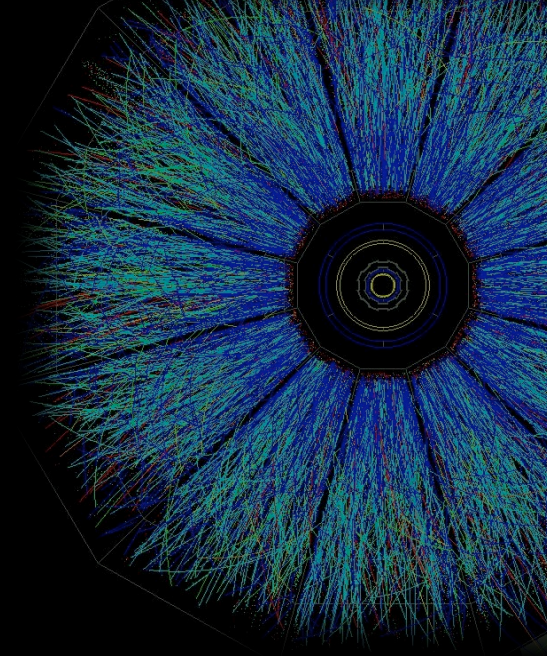
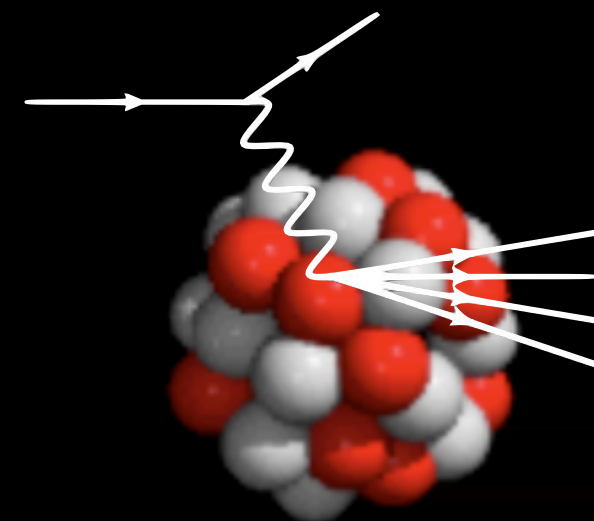
Many body QCD

- Nuclei
 - Spectra
 - Interactions and decays
 - m_q dependence



Many body QCD

- π and K condensation
- EOS in n stars (kaons)
- Critical point at RHIC
- Three π (and K) interferometry (HBT) at SPS and RHIC
- Medium effects



Nuclear physics from LQCD

- Can we compute the mass of ^{208}Pb in QCD?

$$\langle 0|T q_1(t) \dots q_{624}(t) \bar{q}_1(0) \dots \bar{q}_{624}(0)|0\rangle$$



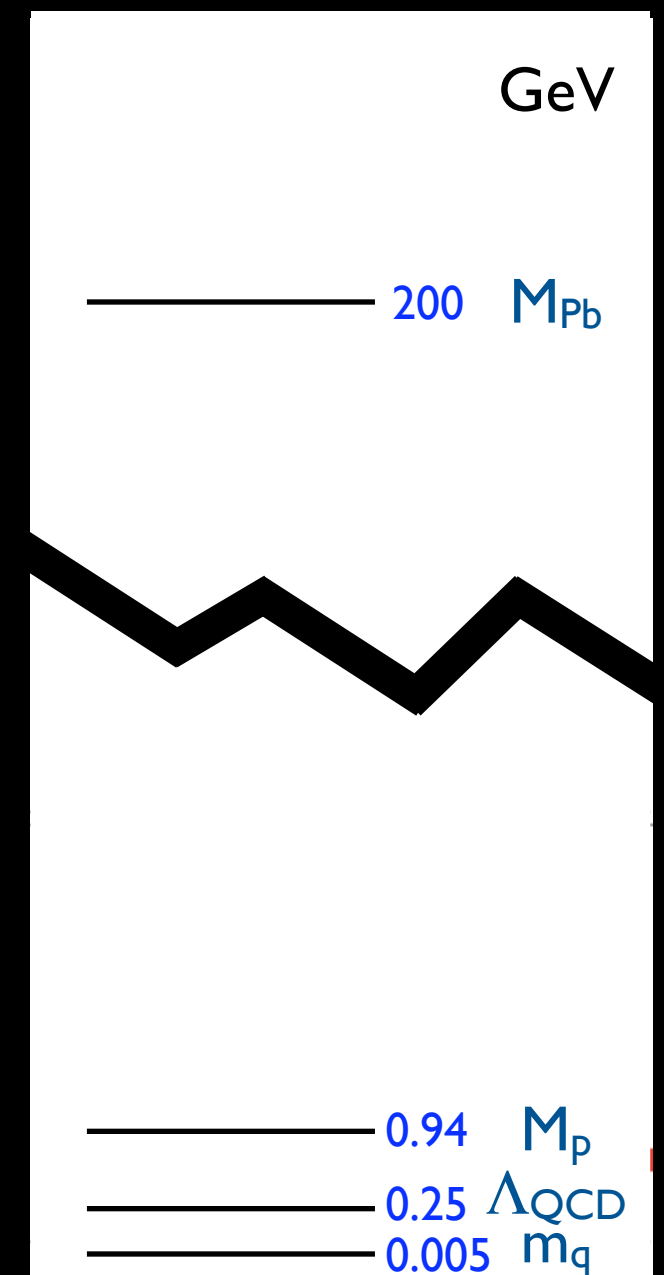
- Long time behaviour gives GS energy

$$\xrightarrow{t \rightarrow \infty} \# \exp(-M_{Pb}t)$$

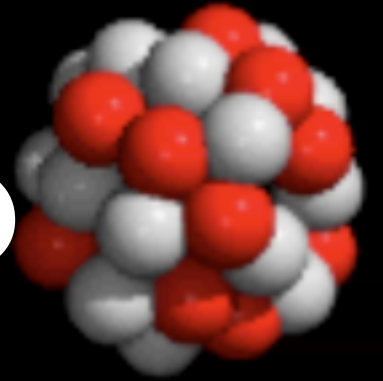
- But...

Nuclear physics is hard!

- Contractions: $(A+Z)!(2A-Z)!$
- Signals for very massive states
- Small energy splittings
- Numerical precision
- Statistical noise: exponentially increases with A



Signal-to-noise ratio



- QCD functional integrals done by importance sampling: propagators
- Variance in correlator determined by

$$\sigma^2 \langle C \rangle = \langle C C^\dagger \rangle - |\langle C \rangle|^2$$

- For nucleon:

$$s/n \sim \exp[-(M_N - 3/2m_\pi)t]$$

- For nucleus A:

$$s/n \sim \exp[-A(M_N - 3/2m_\pi)t]$$



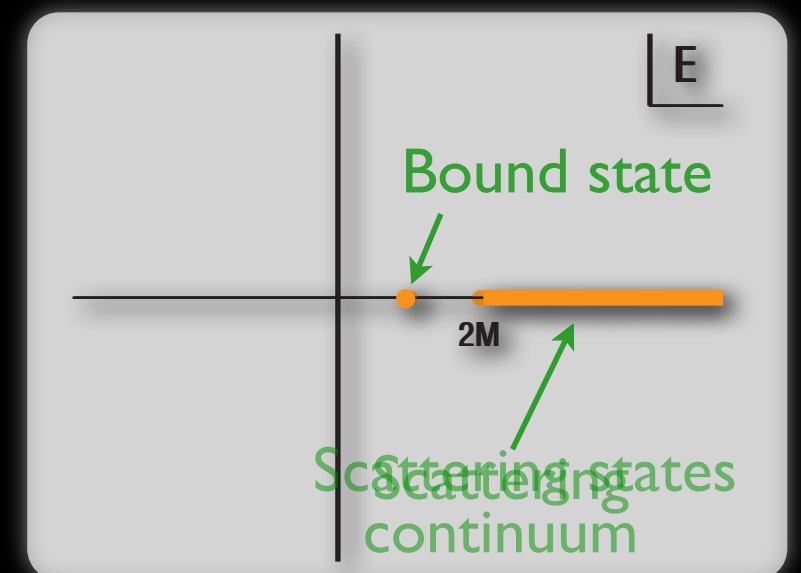


Phys Rev D76:074507, 2007
Phys. Rev. Lett. 100:082004, 2008
Phys Rev D77:057502, 2008
Phys Rev D78:014507, 2008
Phys Rev D78:054514, 2008
Phys. Rev. Lett. 102:032004, 2009
+ ...

Multi-meson systems

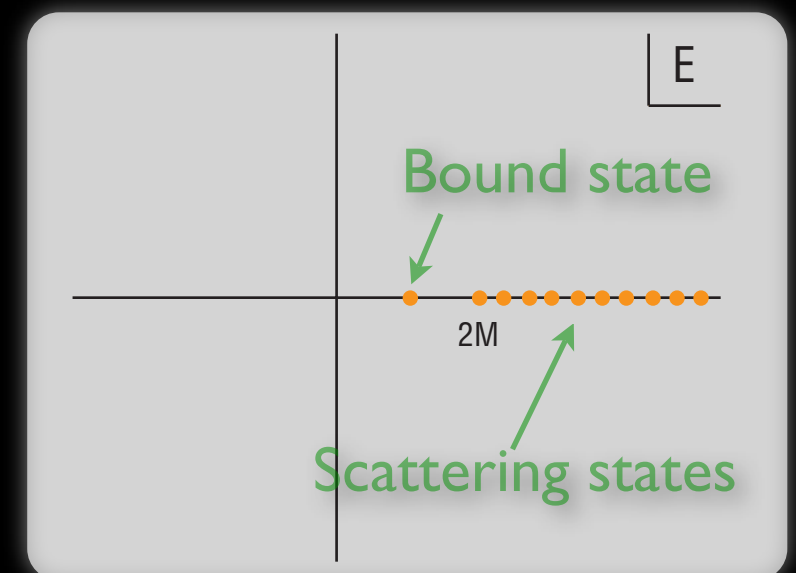
Hadron scattering

- Maiani-Testa: *extracting multi-hadron S-matrix elements from Euclidean lattice calculations of corresponding Green functions is impossible*
- Lüscher: volume dependence of two-particle energy levels \Rightarrow scattering phase-shift up to inelastic threshold



Hadron scattering

- Maiani-Testa: *extracting multi-hadron S-matrix elements from Euclidean lattice calculations of corresponding Green functions is impossible*
- Lüscher: volume dependence of two-particle energy levels \Rightarrow scattering phase-shift up to inelastic threshold
- Exact relation provided $r \ll L$
- Used for $\pi\pi$, KK , NN , ΛN



Bosons in a box

[Beane, WD & Savage; WD & Savage]



- Large volume expansion of GS energy of n meson system to $1/L^7$
- 2 & 3 body interactions (N body: $L^{-3(N-1)}$)
- Relativistic up to particle production
- $n=2$: reproduces expansion of Lüscher
- Can include higher PW, higher body, excited states

Bosons in a box

[WD+Savage arXiv:0801.0763]

$$\Delta E_n = \frac{4\pi\bar{a}}{M L^3} {}^n C_2 \left\{ 1 - \left(\frac{\bar{a}}{\pi L}\right) \mathcal{I} + \left(\frac{\bar{a}}{\pi L}\right)^2 [\mathcal{I}^2 + (2n-5)\mathcal{J}] \right. \\ \left. - \left(\frac{\bar{a}}{\pi L}\right)^3 [\mathcal{I}^3 + (2n-7)\mathcal{I}\mathcal{J} + (5n^2 - 41n + 63)\mathcal{K}] \right. \\ \left. + \left(\frac{\bar{a}}{\pi L}\right)^4 [\mathcal{I}^4 - 6\mathcal{I}^2\mathcal{J} + (4+n-n^2)\mathcal{J}^2 + 4(27-15n+n^2)\mathcal{I}\mathcal{K} \right. \\ \left. + (14n^3 - 227n^2 + 919n - 1043)\mathcal{L} + 16(n-2)(\mathcal{T}_0 + n\mathcal{T}_1)] \right\} \\ + {}^n C_3 \frac{1}{L^6} \hat{\eta}_3^L + {}^n C_3 \frac{6\pi\bar{a}^3}{M^3 L^7} (n+3) \mathcal{I} + \mathcal{O}(L^{-8})$$

$$\bar{a} = a + \frac{2\pi}{L^3} a^3 r$$

Scattering length

$$\hat{\eta}_3^L = \bar{\eta}_3^L \left[1 - \frac{6\bar{a}}{\pi L} \mathcal{I} \right] + \frac{72\pi\bar{a}^4 r}{ML} \mathcal{I}$$

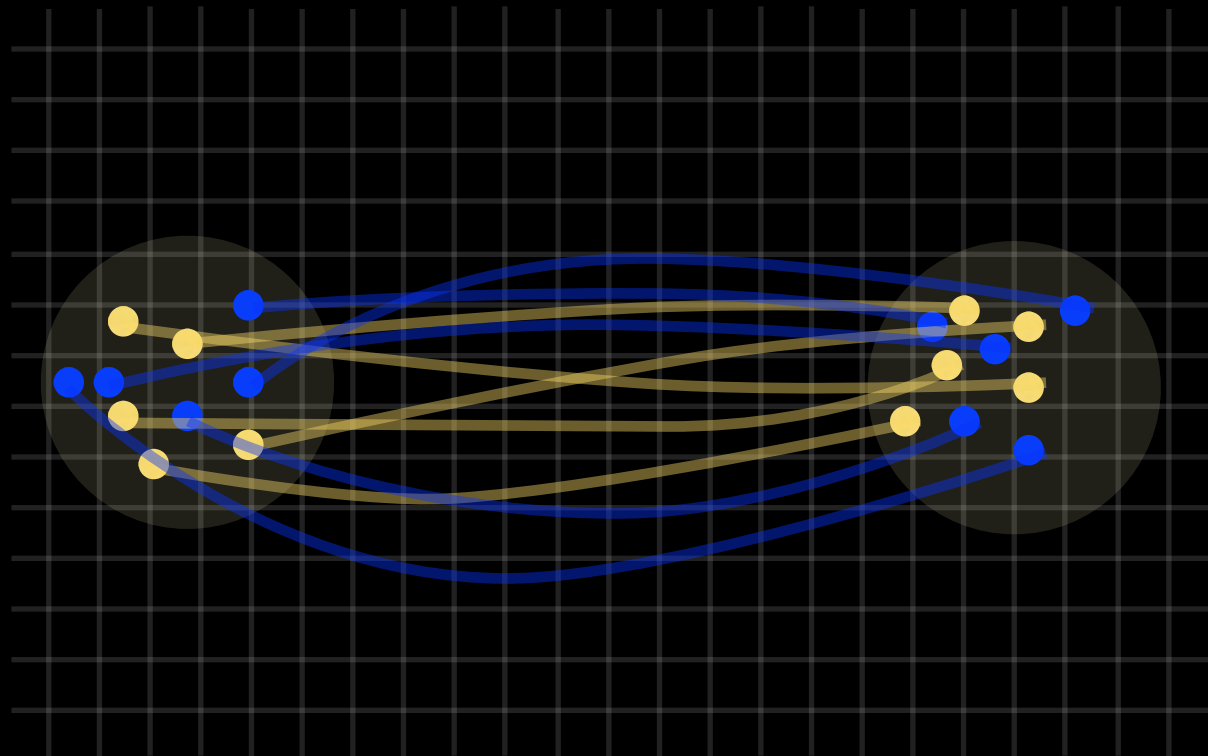
Three body interaction

$\mathcal{I}, \mathcal{J}, \dots$
geometric
constants

Many mesons in LQCD

- Consider n π^+ correlator ($m_u=m_d$)

$$C_n(t) = \left\langle 0 \left| \left[\sum_{\mathbf{x}} \bar{d}\gamma_5 u(\mathbf{x}, t) \bar{u}\gamma_5 d(0, 0) \right]^n \right| 0 \right\rangle$$
$$\rightarrow A e^{-E_n t}$$



Many mesons in LQCD

- Consider n π^+ correlator ($m_u=m_d$)

$$C_n(t) = \left\langle 0 \left| \left[\sum_{\mathbf{x}} \bar{d} \gamma_5 u(\mathbf{x}, t) \bar{u} \gamma_5 d(0, 0) \right]^n \right| 0 \right\rangle$$
$$\rightarrow A e^{-E_n t}$$

- $n!^2$ Wick contractions: $(12!)^2 \sim 10^{17}$

$$C_3(t) = \text{tr} [\Pi]^3 - 3 \text{tr} [\Pi] \text{tr} [\Pi^2] + 2 \text{tr} [\Pi^3]$$

$$\Pi = \sum_{\mathbf{x}} \gamma_5 S(\mathbf{x}, t; 0) \gamma_5 S^\dagger(\mathbf{x}, t; 0)$$



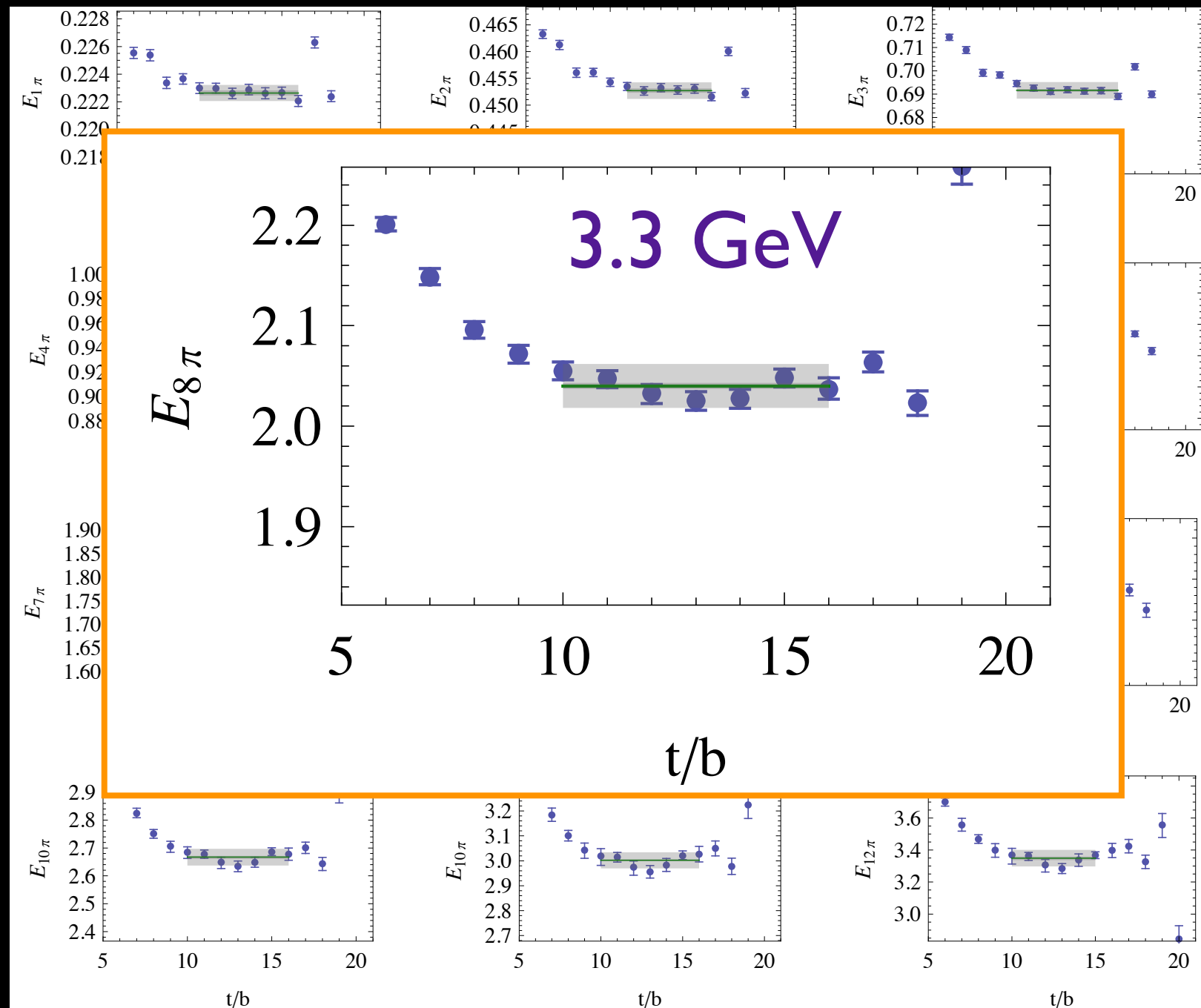
- Maximal isospin: only a single quark propagator

Lattice details

- Calculations use MILC gauge configurations
 - $L=2.5$ fm, $a=0.12$ fm, *rooted* staggered
 - also $L=3.5$ fm and $a=0.09$ fm
- NPLQCD: domain-wall quark propagators
 - $m_\pi \sim 291, 318, 352, 358, 491, 591$ MeV
 - 24 propagators / lattice in best case
- $I_z=n=1, \dots, 12$ pion and ($S=n$) kaon systems

n-meson energies

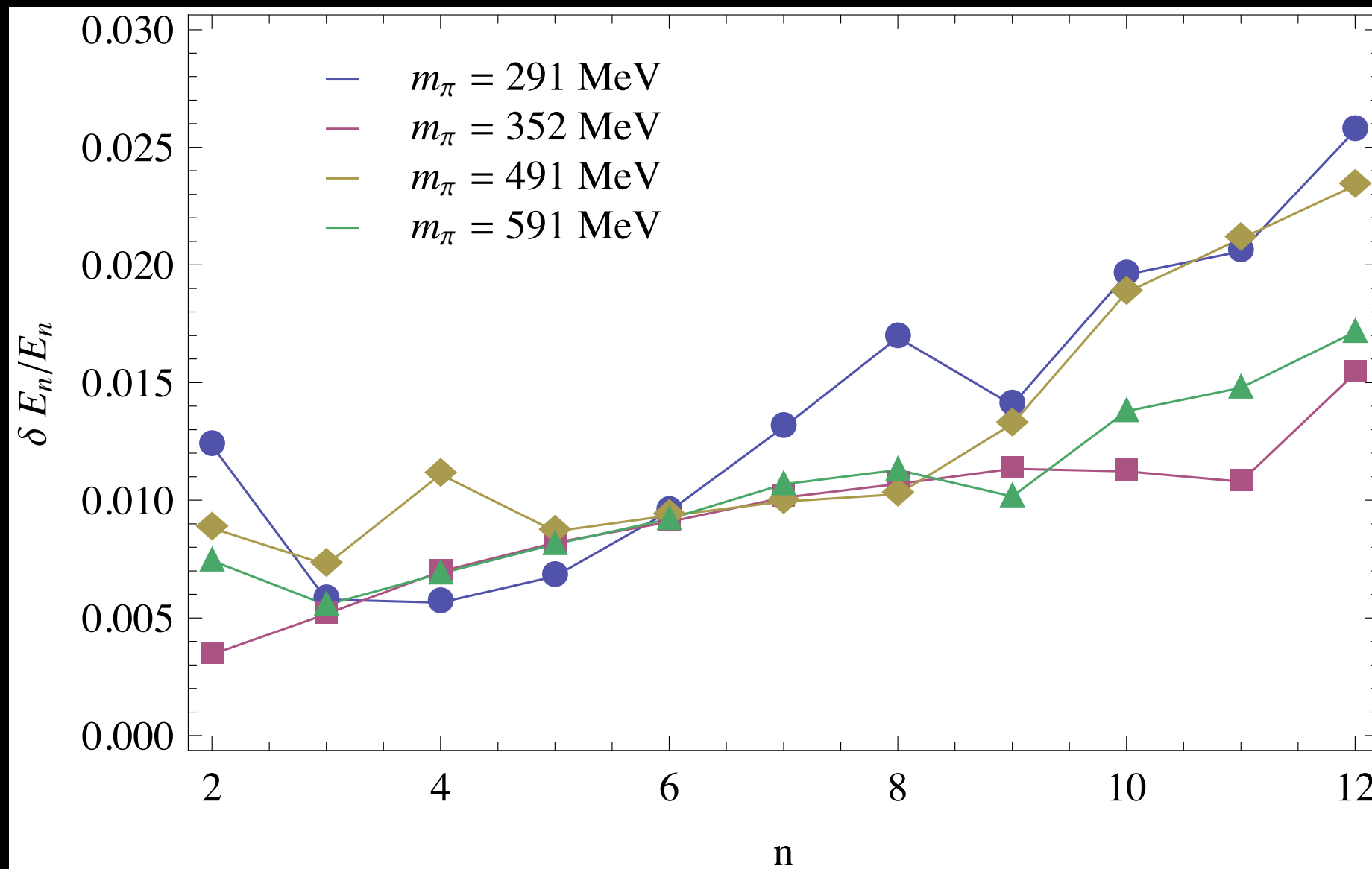
- Effective energy plots: $\log[C_n(t)/C_n(t+1)]$



$m_\pi = 352 \text{ MeV}$

n-meson energies

- Clean signals for $n=1, \dots, 12$



Bosons in a box



- Multiple extractions of \bar{a} and $\hat{\eta}_3^L$

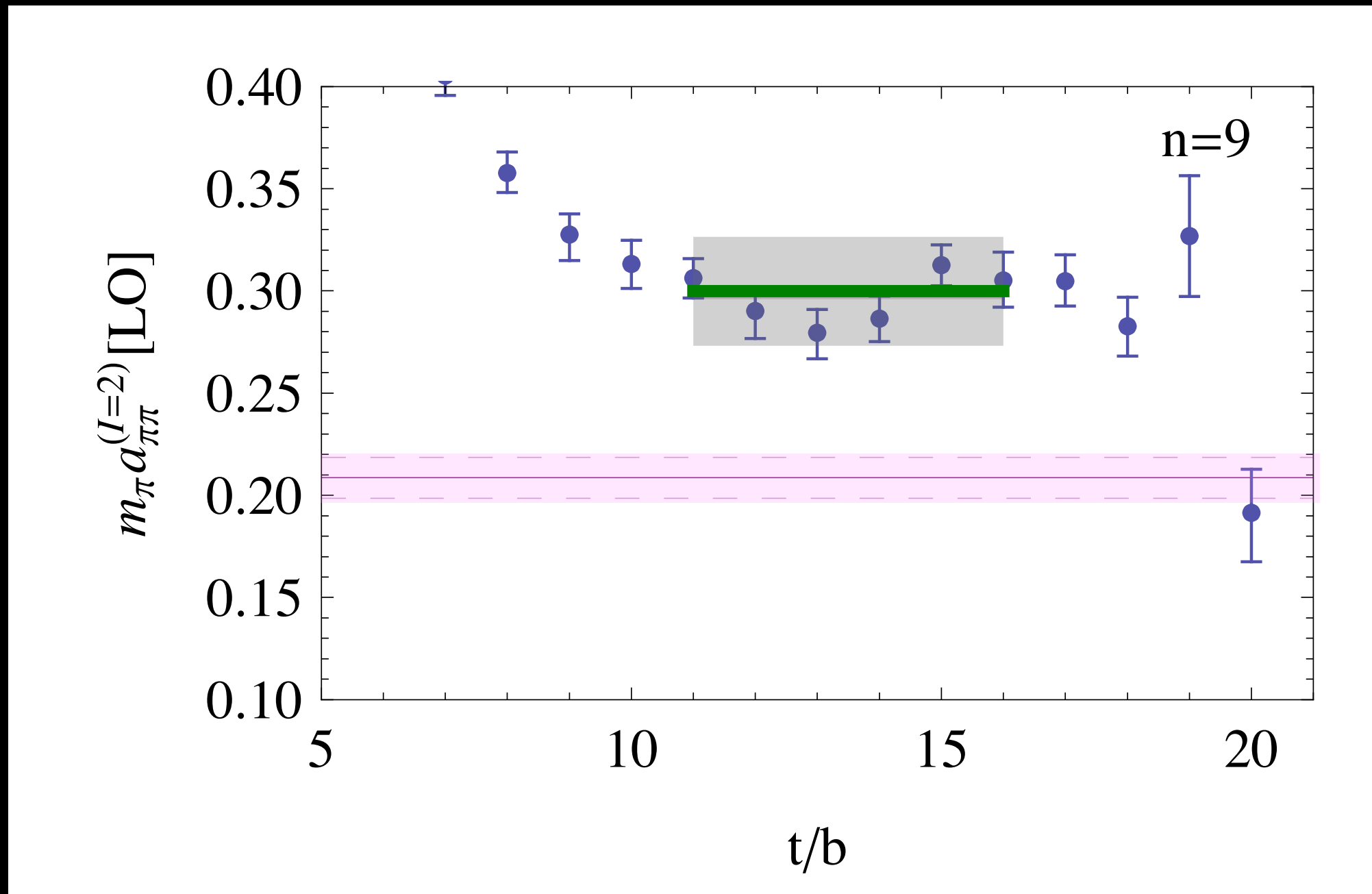
$$\Delta E_n = \frac{4\pi\bar{a}}{M L^3} {}^n C_2 \left\{ 1 - \left(\frac{\bar{a}}{\pi L}\right) \mathcal{I} + \left(\frac{\bar{a}}{\pi L}\right)^2 [\mathcal{I}^2 + (2n-5)\mathcal{J}] \right. \\ \left. - \left(\frac{\bar{a}}{\pi L}\right)^3 [\mathcal{I}^3 + (2n-7)\mathcal{I}\mathcal{J} + (5n^2 - 41n + 63)\mathcal{K}] \right. \\ \left. + \left(\frac{\bar{a}}{\pi L}\right)^4 [\mathcal{I}^4 - 6\mathcal{I}^2\mathcal{J} + (4+n-n^2)\mathcal{J}^2 + 4(27-15n+n^2)\mathcal{I}\mathcal{K} \right. \\ \left. + (14n^3 - 227n^2 + 919n - 1043)\mathcal{L} + 16(n-2)(\mathcal{T}_0 + n\mathcal{T}_1)] \right\} \\ + {}^n C_3 \frac{1}{L^6} \hat{\eta}_3^L + {}^n C_3 \frac{6\pi\bar{a}^3}{M^3 L^7} (n+3) \mathcal{I} + \mathcal{O}(L^{-8})$$

$$\bar{a} = a + \frac{2\pi}{L^3} a^3 r$$

$$\hat{\eta}_3^L = \bar{\eta}_3^L \left[1 - \frac{6\bar{a}}{\pi L} \mathcal{I} \right] + \frac{72\pi\bar{a}^4 r}{ML} \mathcal{I}$$

Pion scattering

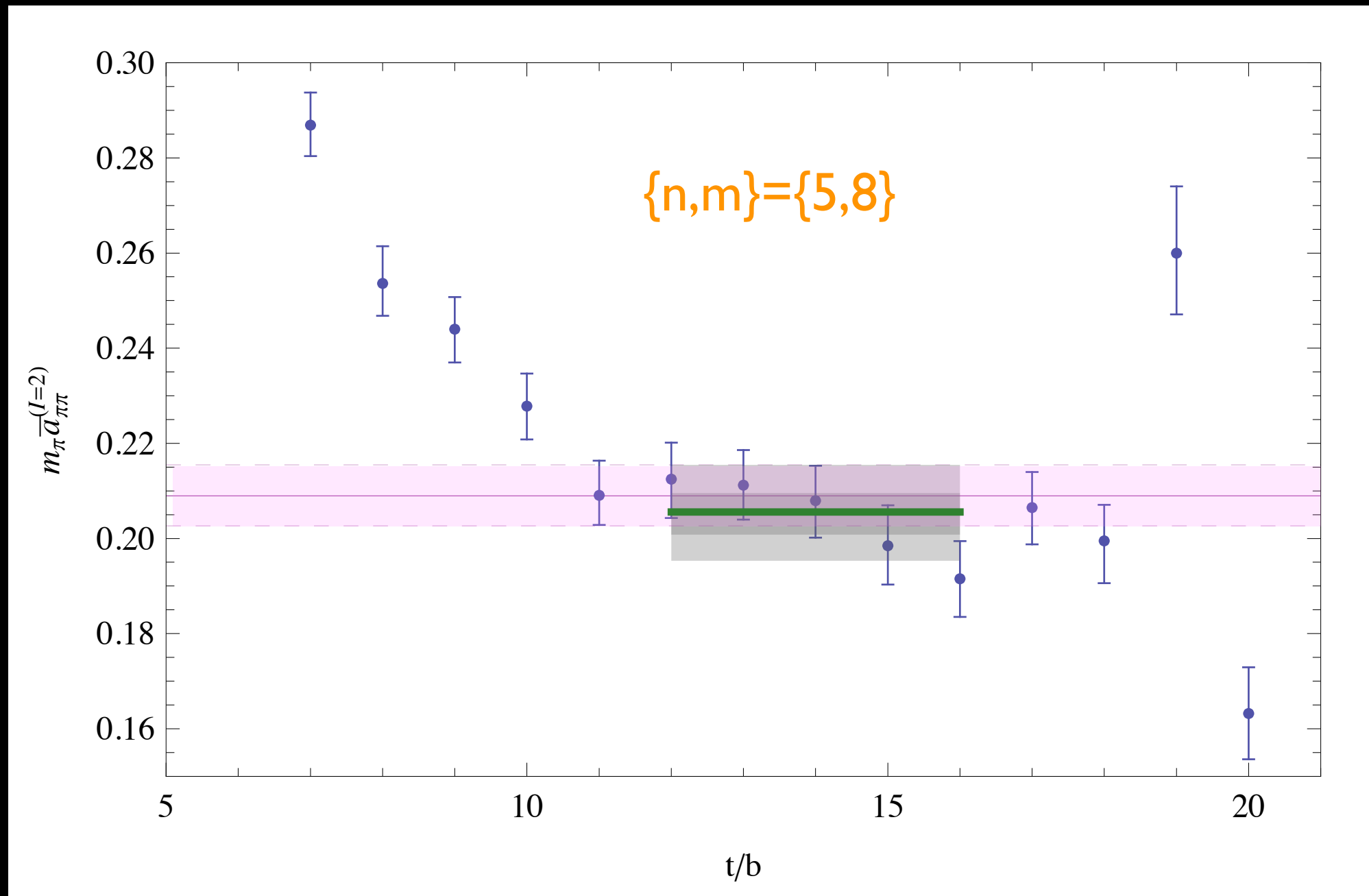
- Extractions of $m_\pi a$ from four orders in L



Lüscher
exact
two-body

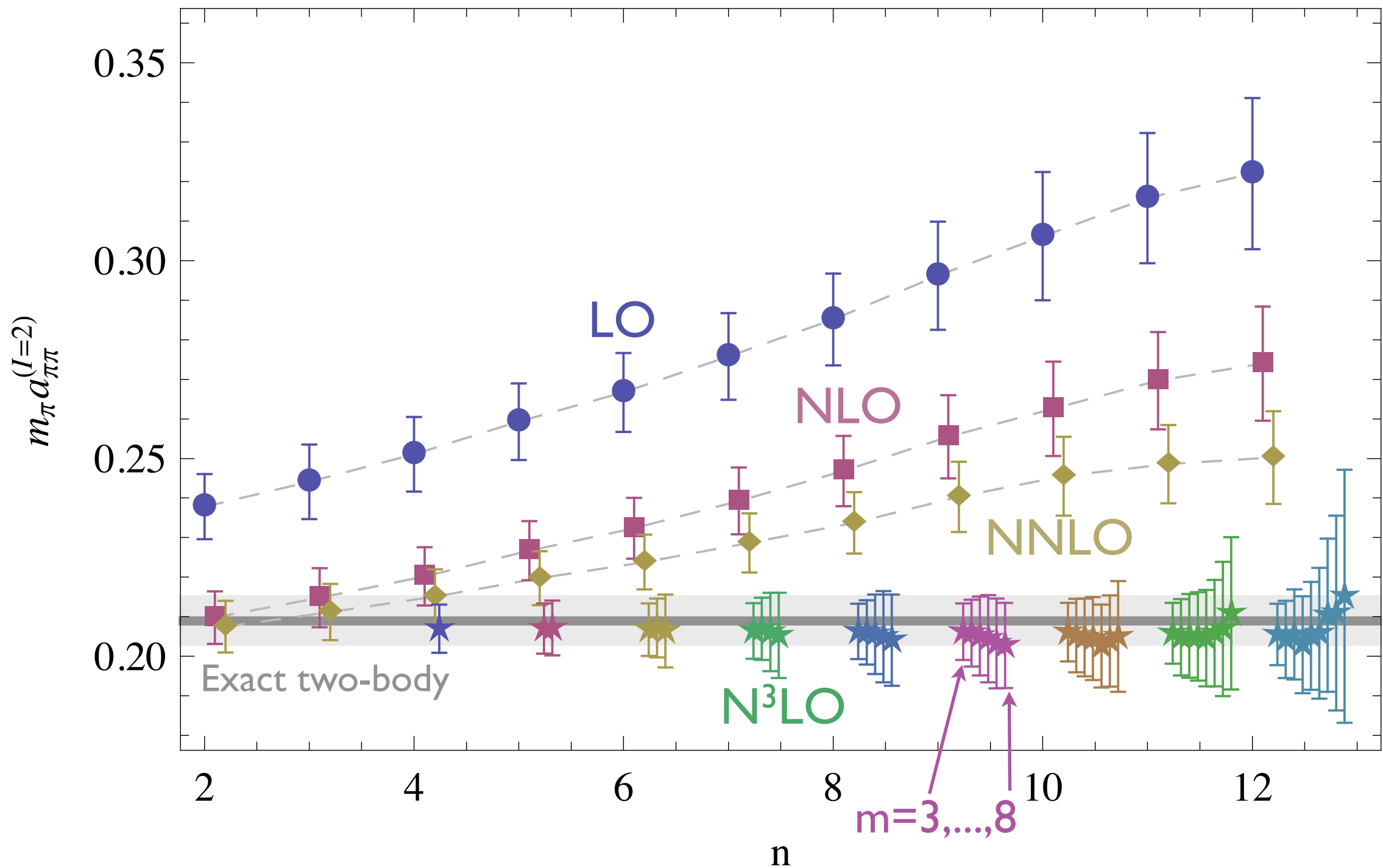
N³LO

- Two energies to cancel 3-body: 45 combinations



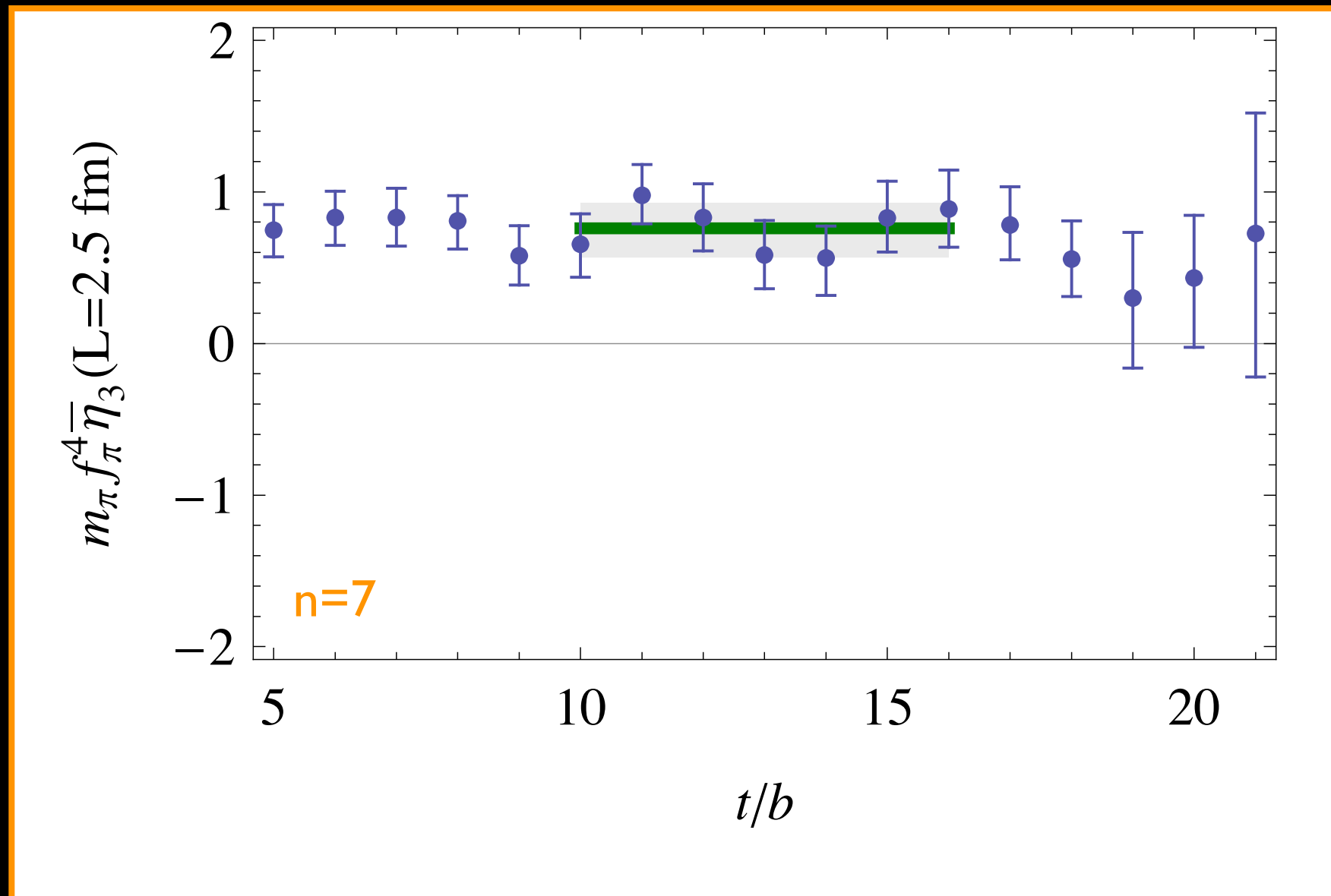
Lüscher
exact
two-body

Pion scattering



Expansion shows no sign of breakdown?

$\pi^+\pi^+\pi^+$ interaction



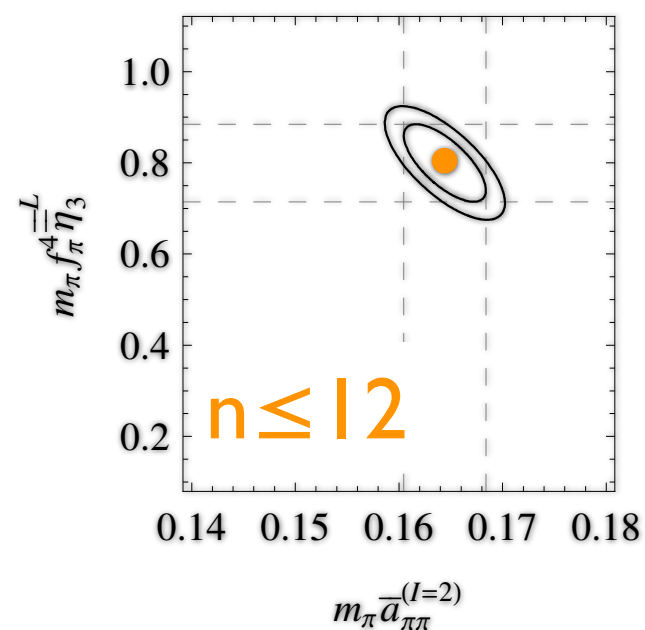
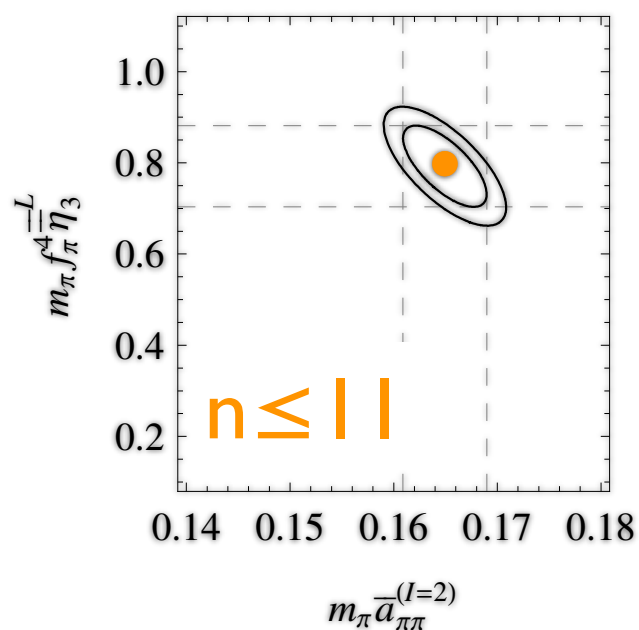
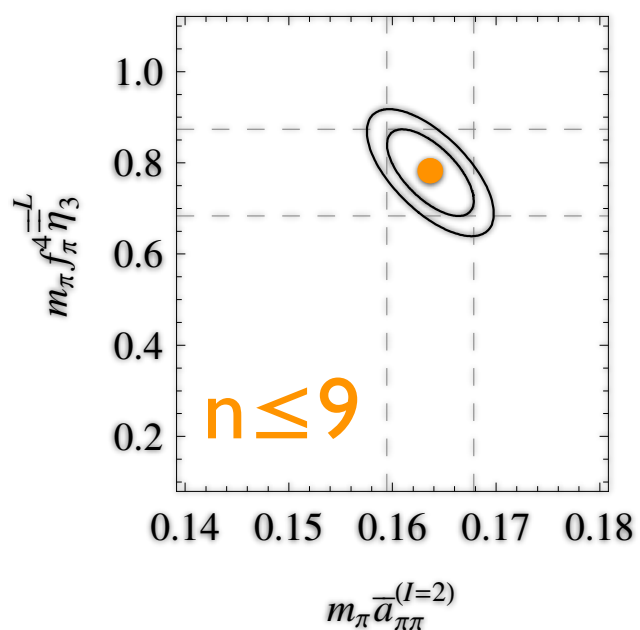
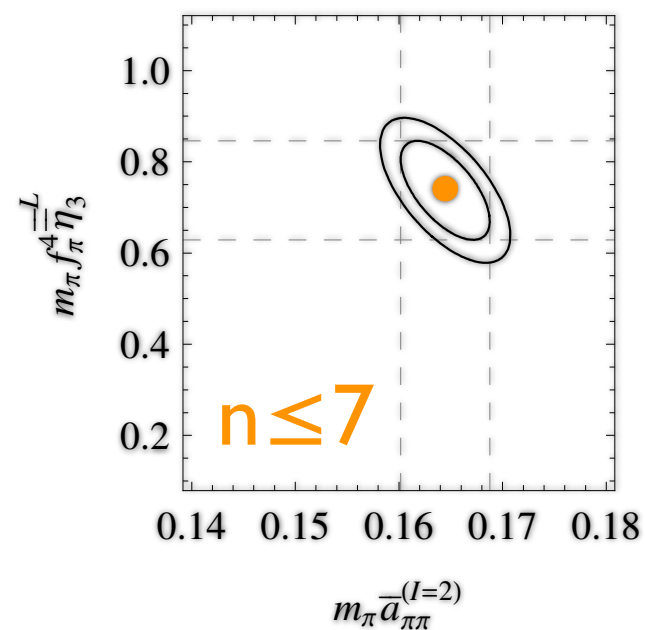
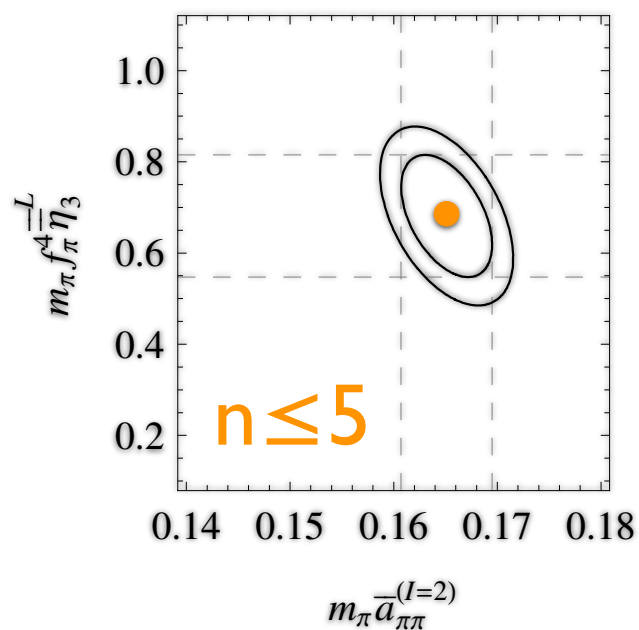
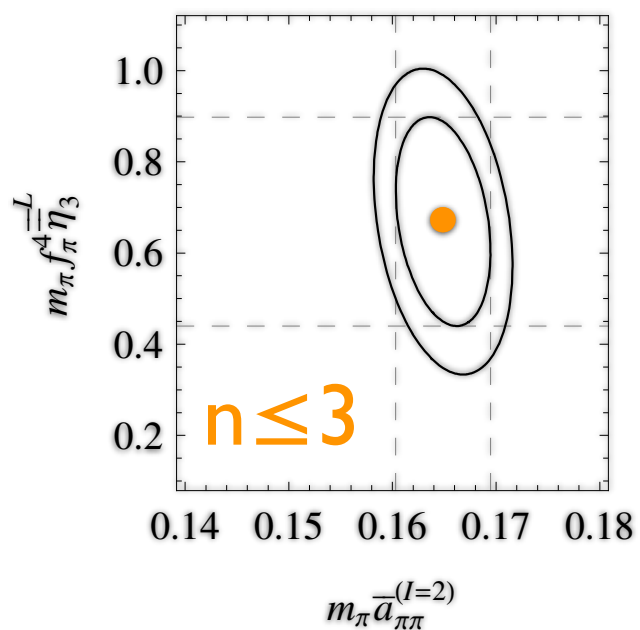
$$m_\pi = 352 \text{ MeV}$$

n correlations

- Fit effective energies to extract two parameters: a , η_3 from $1/L$ expansion
- Use 12 eff. energies in n - t -correlated analysis
 - Large correlation matrix: correlated χ^2
 - Reduces uncertainties as n pion correlators “explore more of the lattice”

n correlations

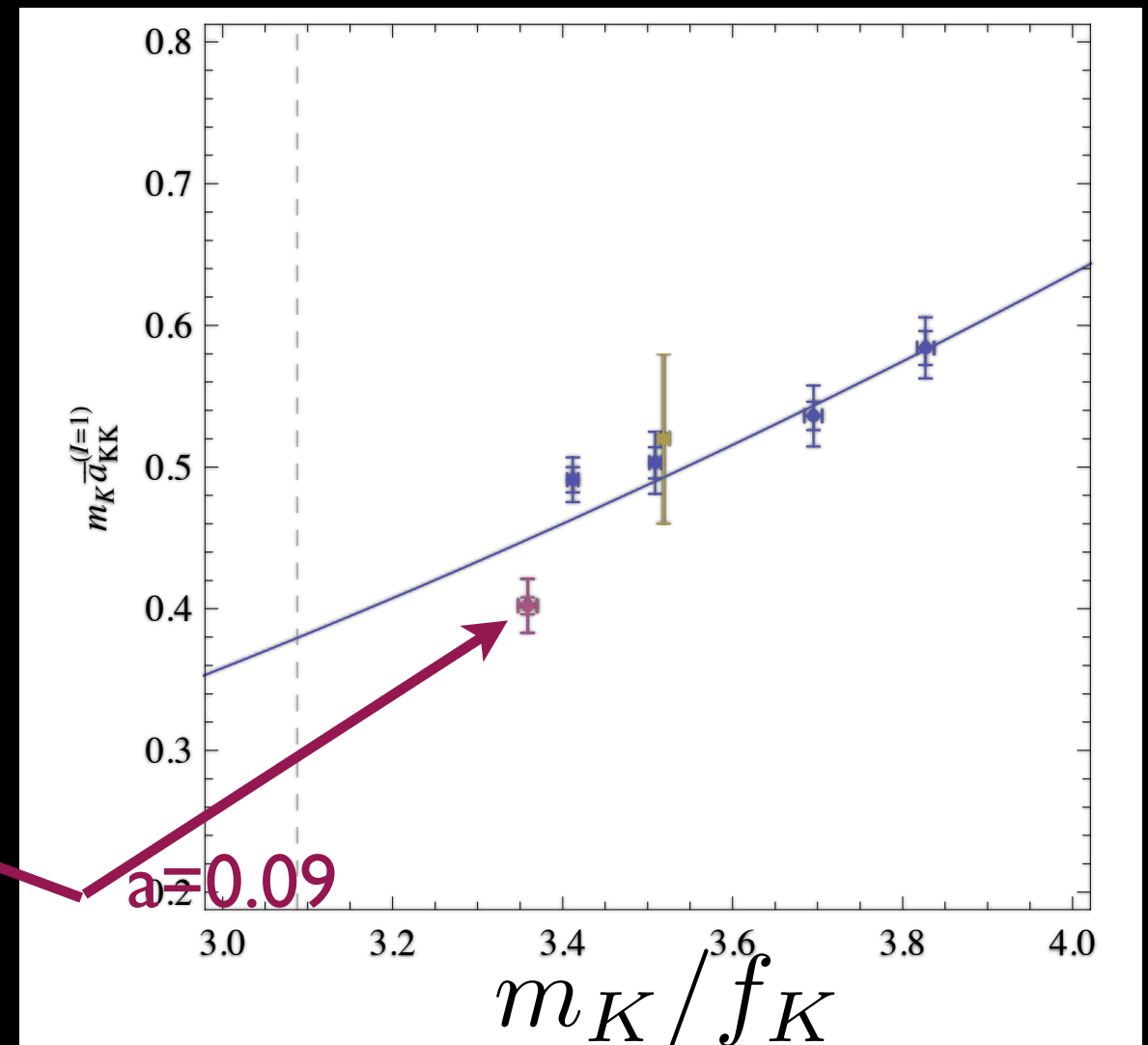
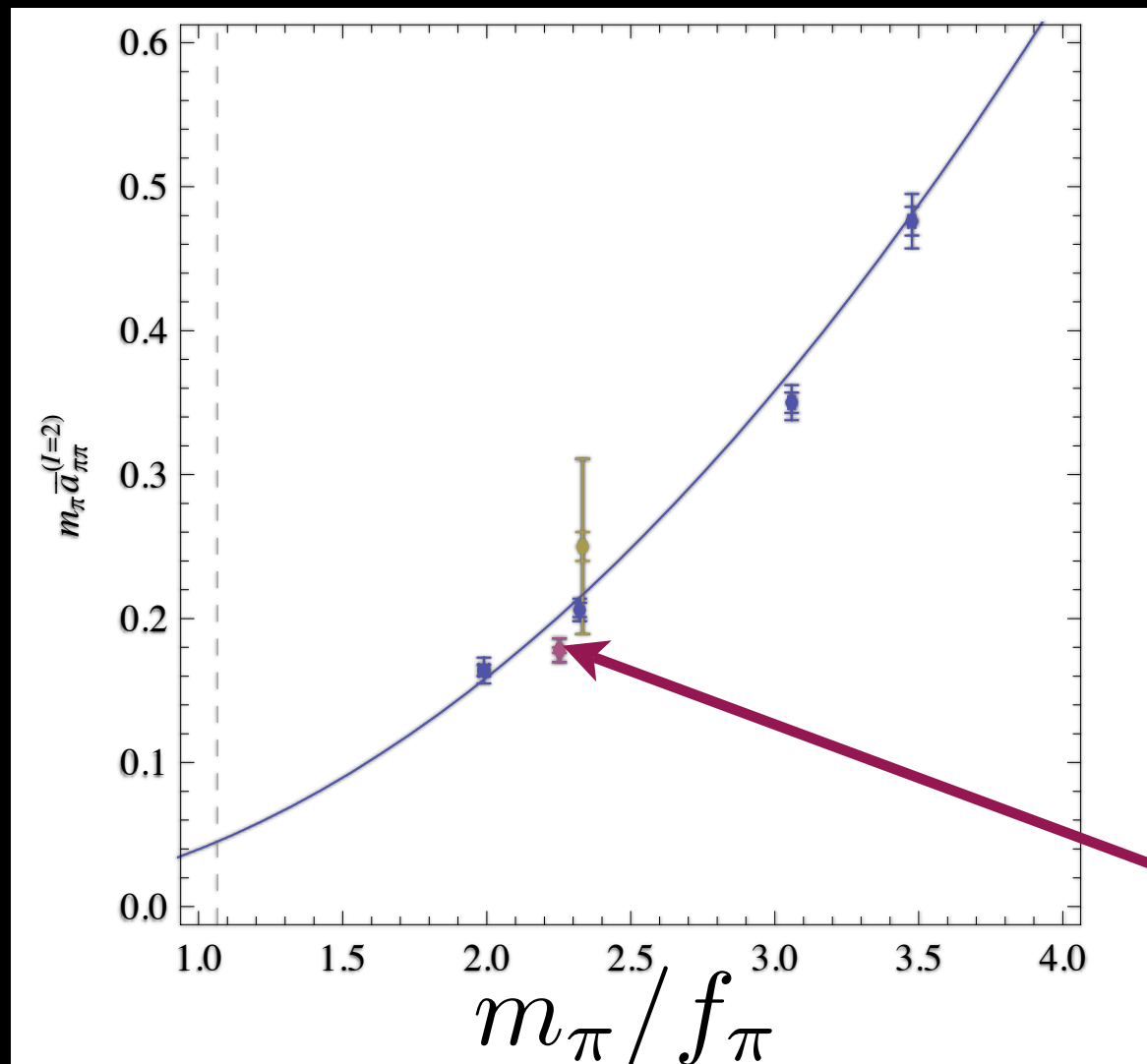
Three-body



Two body

$2\pi^+$ and $2K^-$ interaction

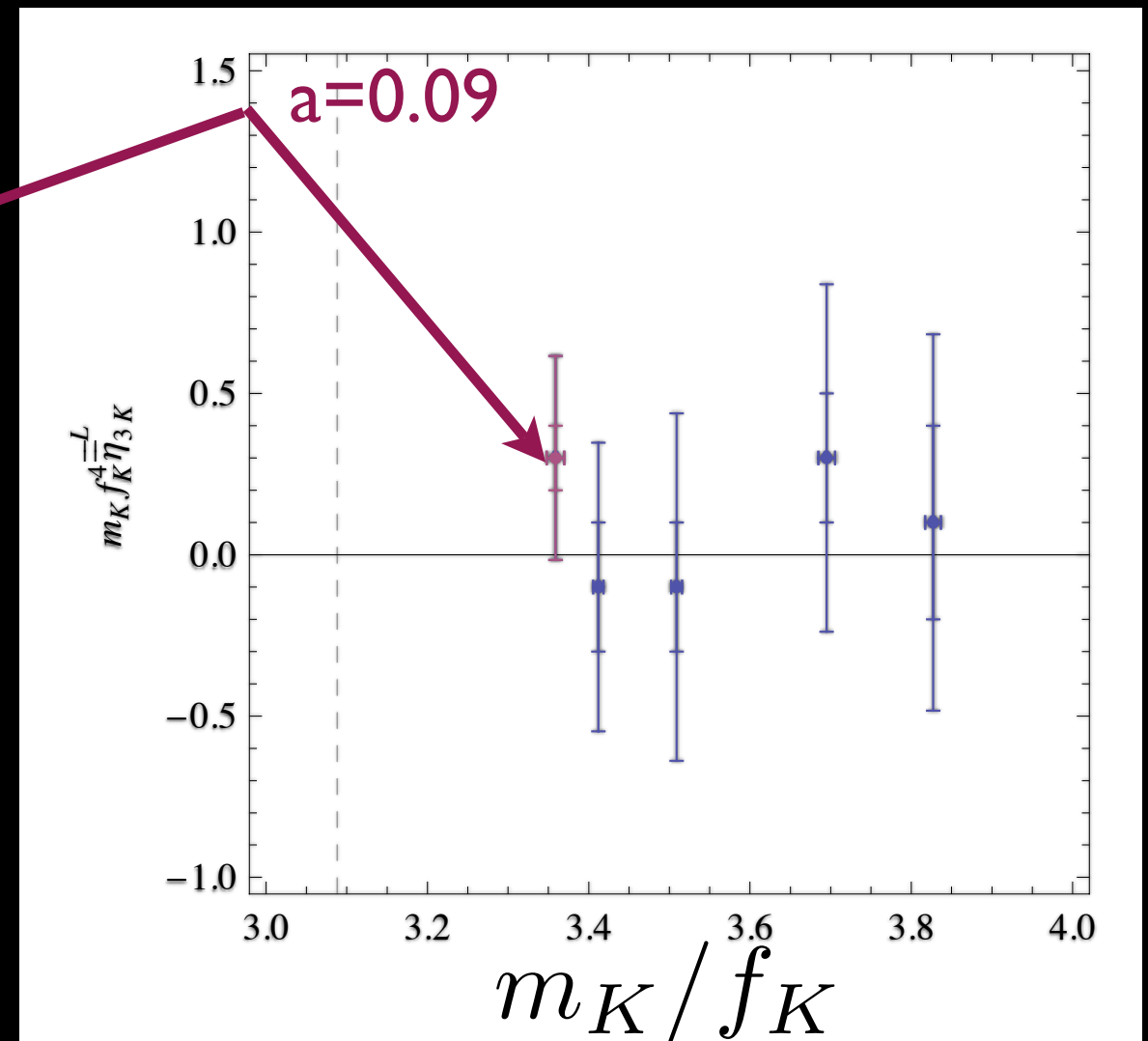
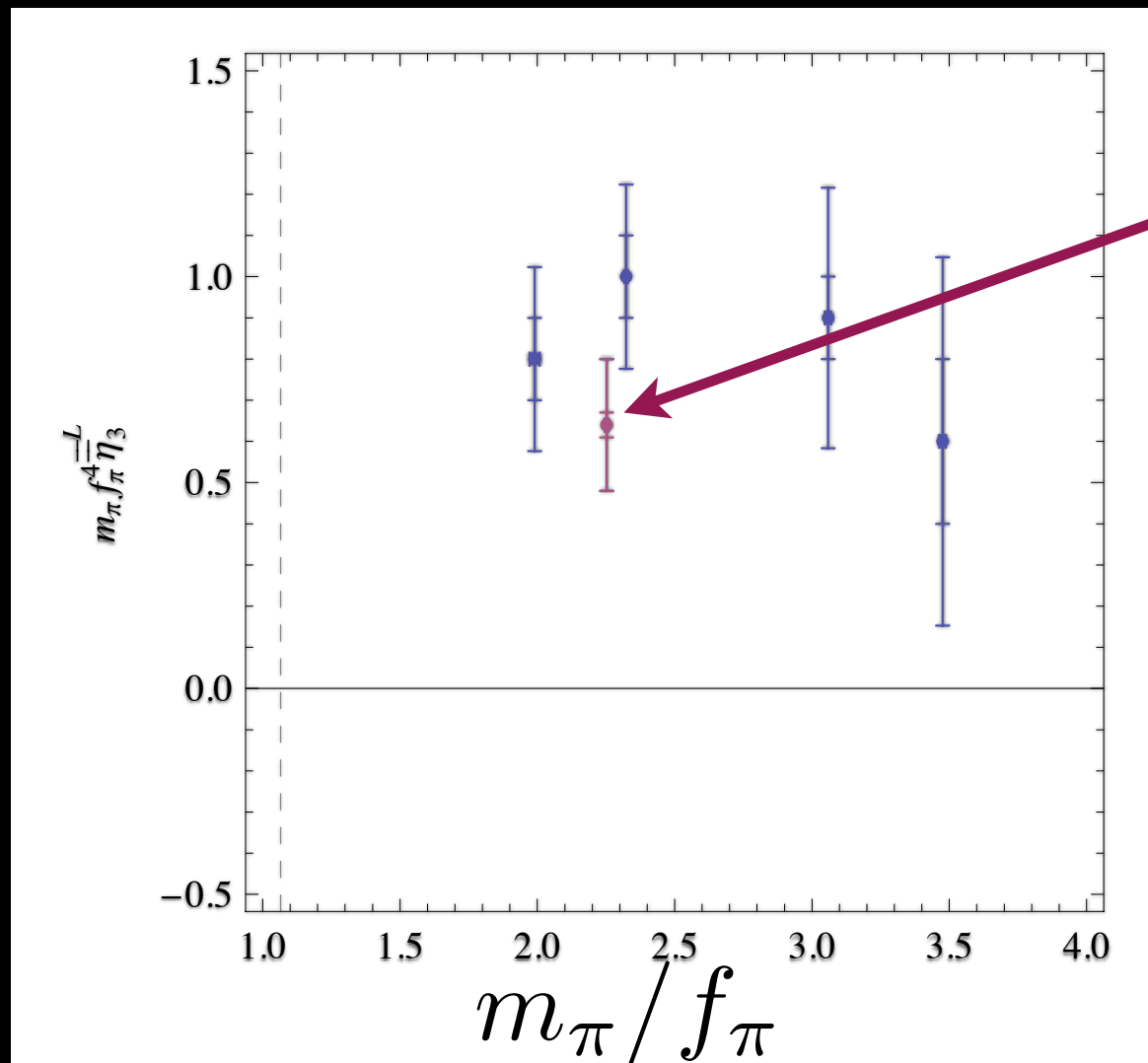
- Scattering lengths



curves: Weinberg

$3\pi^+$ and $3K^-$ interaction

- First QCD three body interaction



Naïve dimension analysis: 1

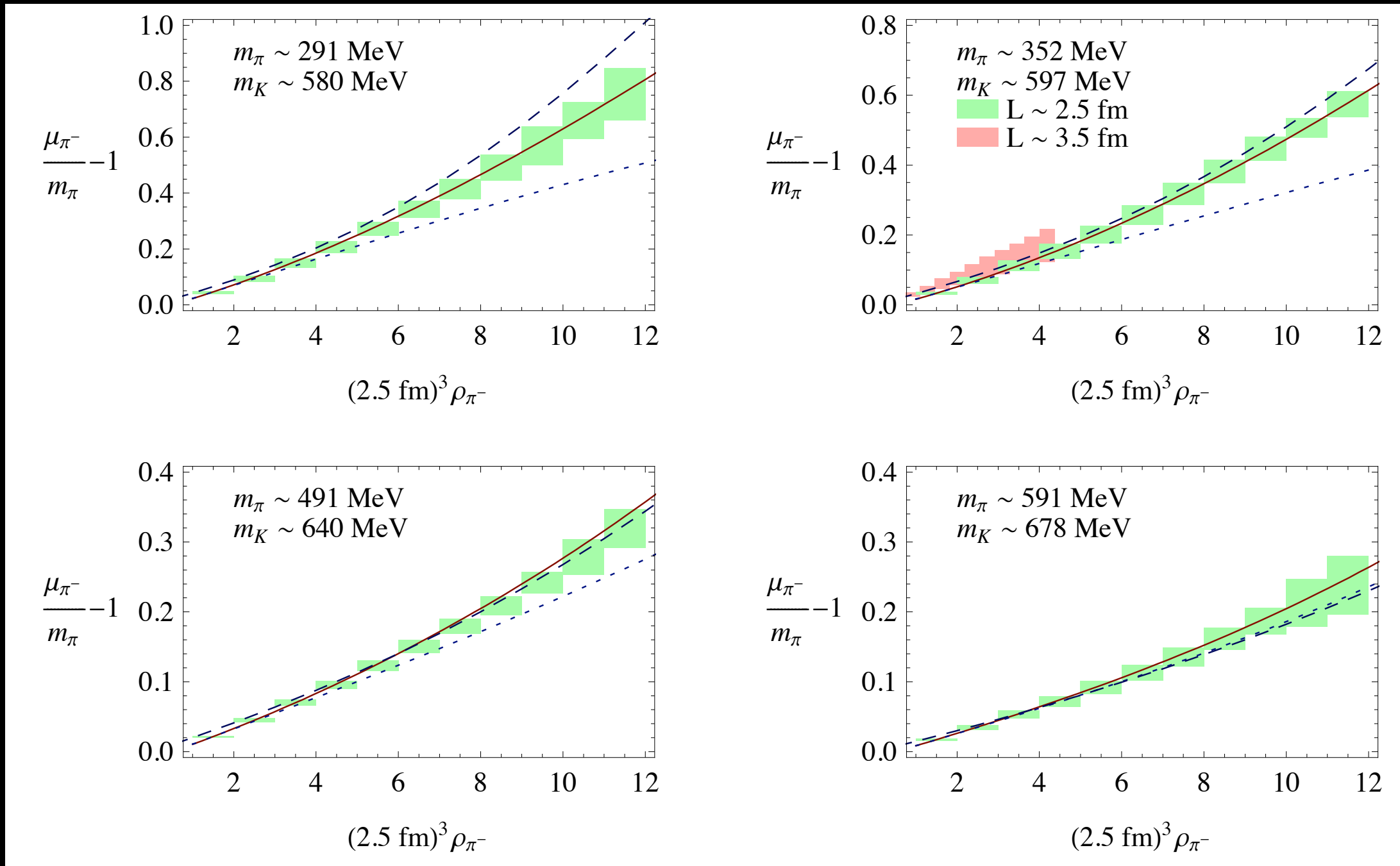
Equation of State

- I/L expansion: analytic form of EOS
- Chemical potential

$$\mu = \left. \frac{d E}{d n} \right|_{V_{const}}$$

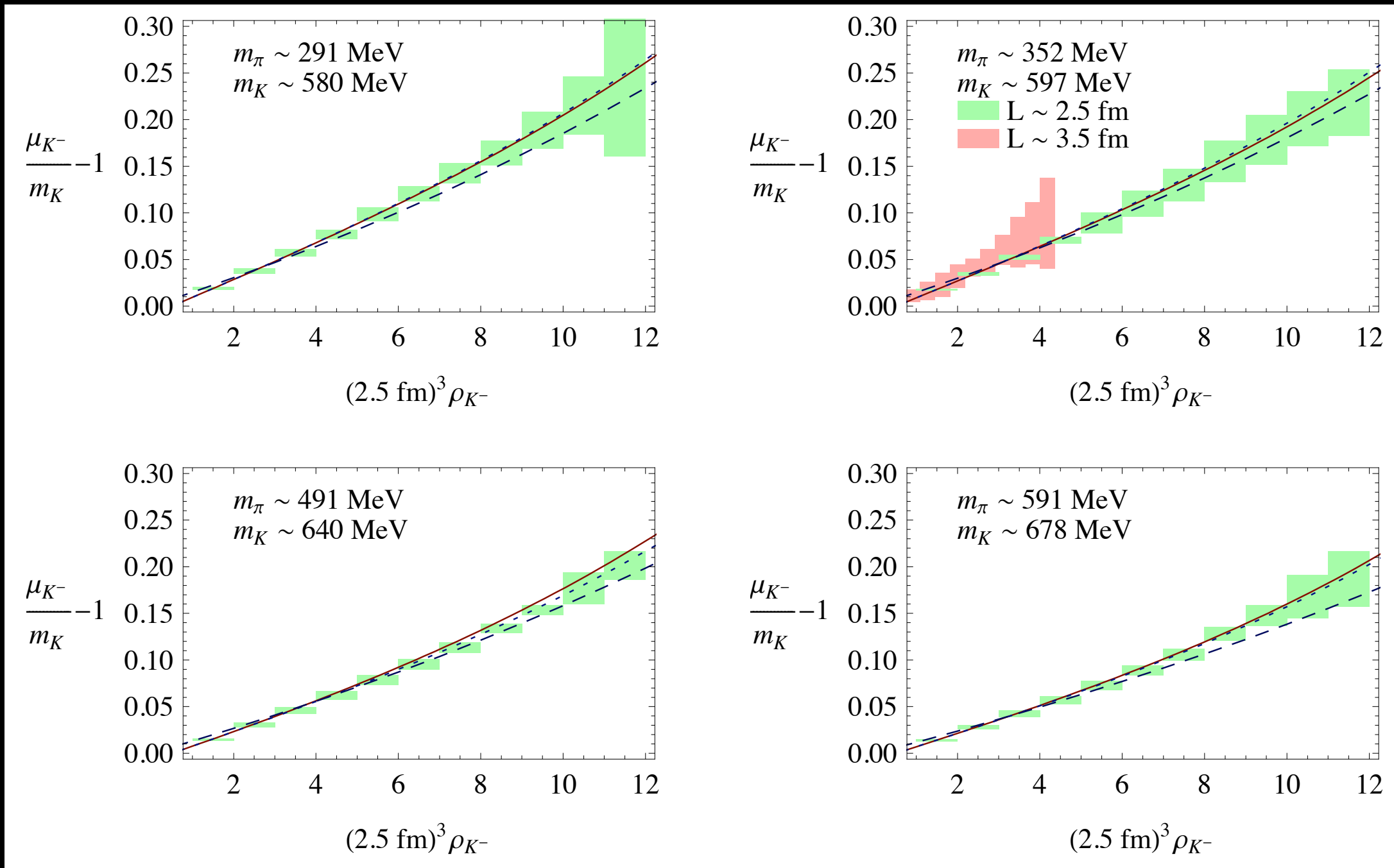
- $\mu(\rho)$ numerically using finite difference
- Compare with LO χ PT [Son & Stephanov]

Isospin Chemical Potential



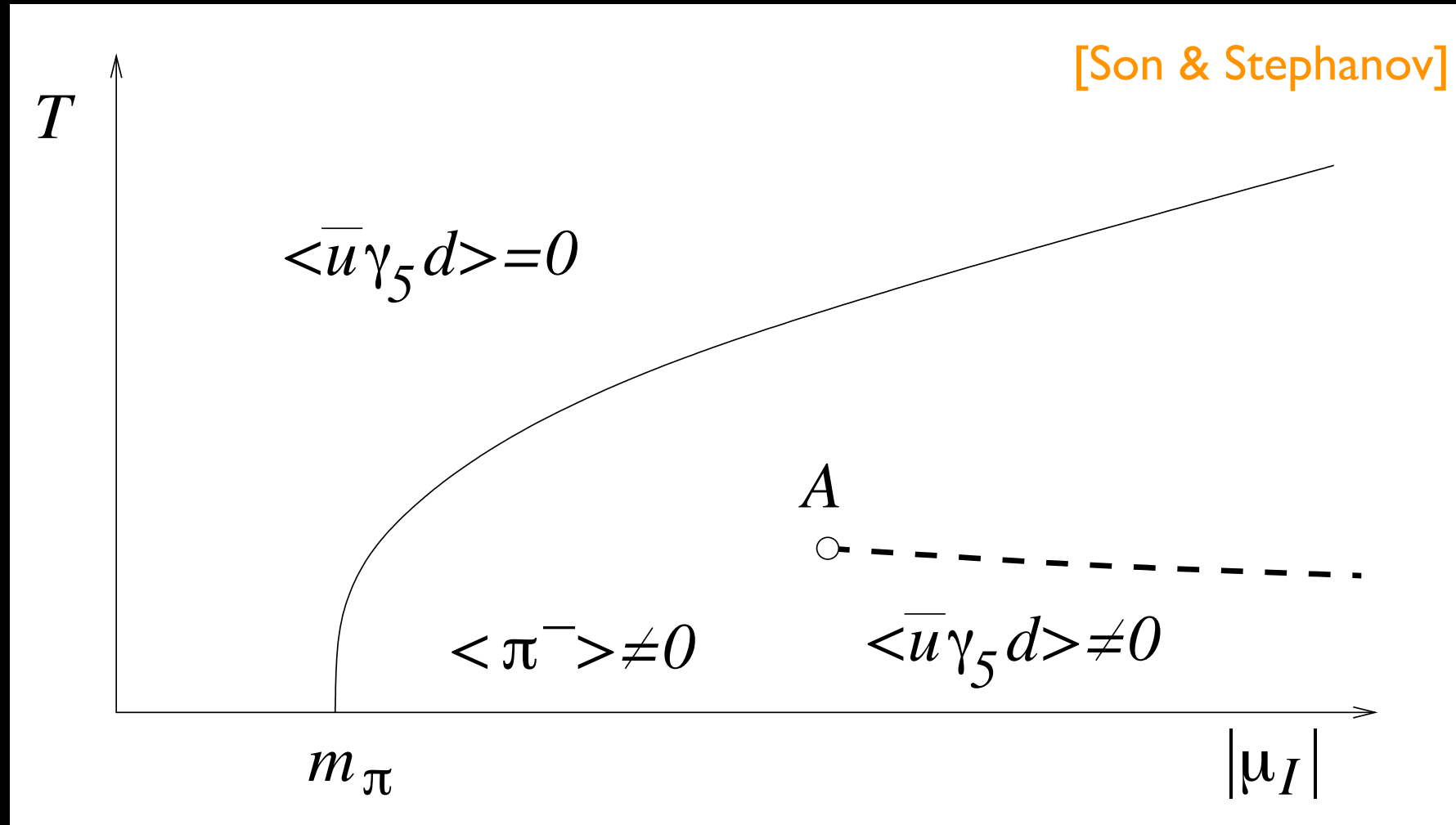
— 2+3 body fit
 ⋯ No 3 body
 - - - LO χ PT

Kaon Chemical Potential



- supports KN analysis of K^- condensation in n-stars

μ_I - T phase diagram

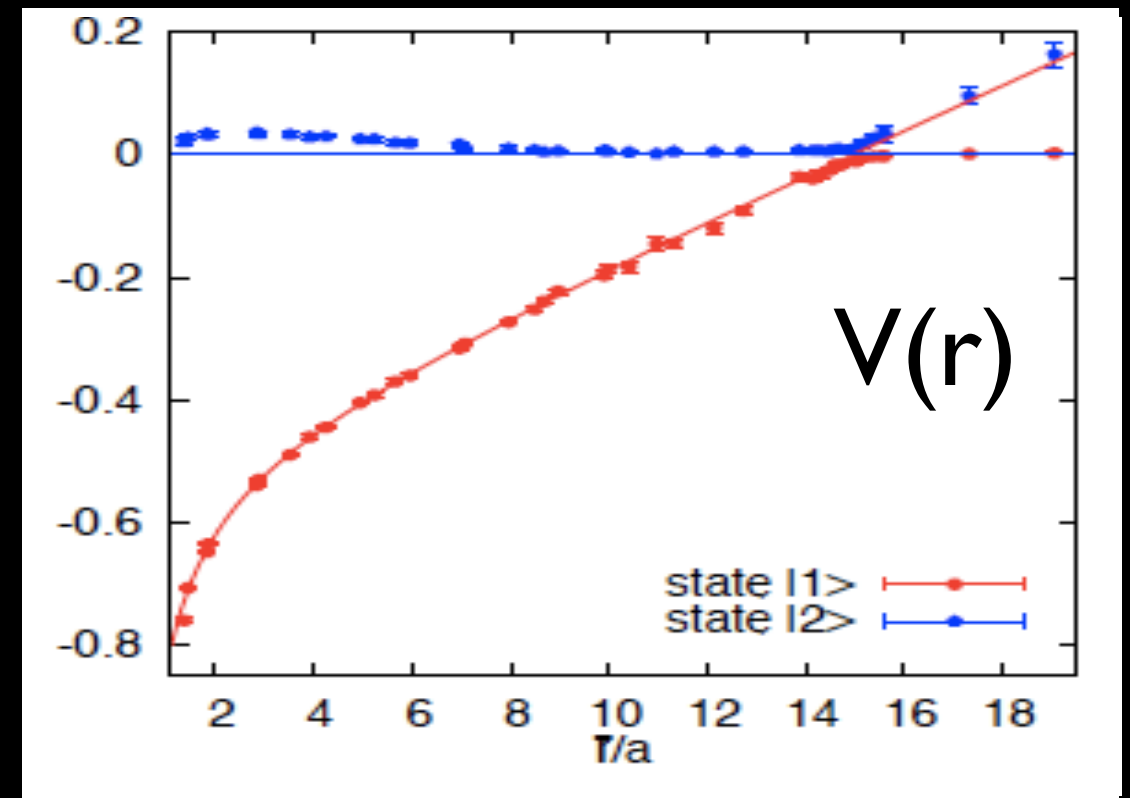
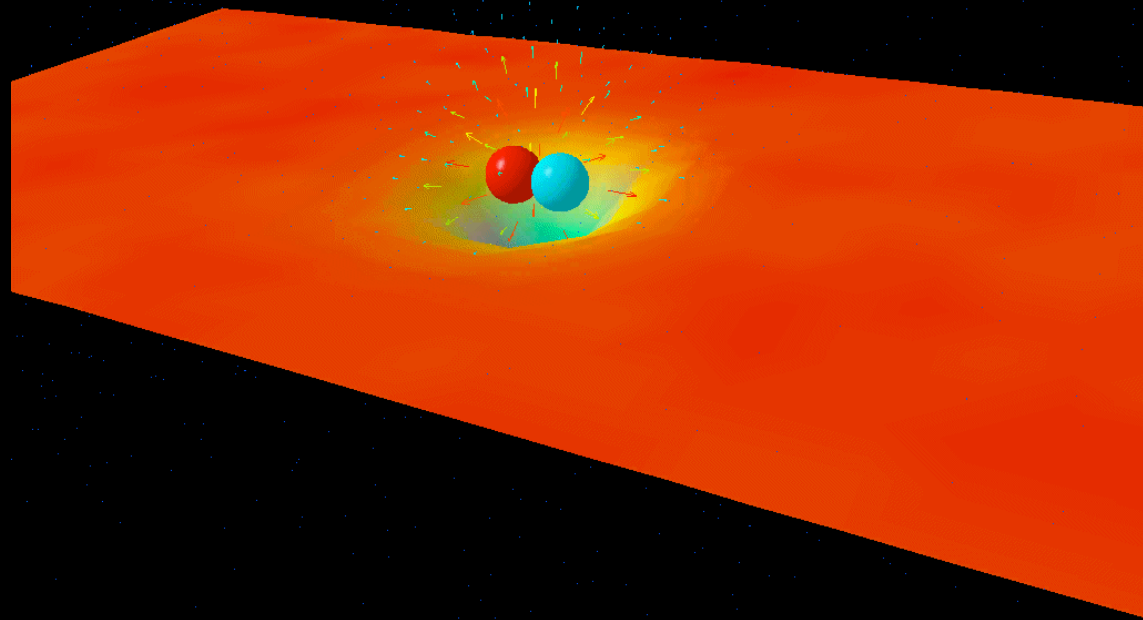


- BEC/BCS transition ?
- Vector meson condensate ?

Colour screening by pions

[WD+ M Savage PRL 09]

- Static quark potential

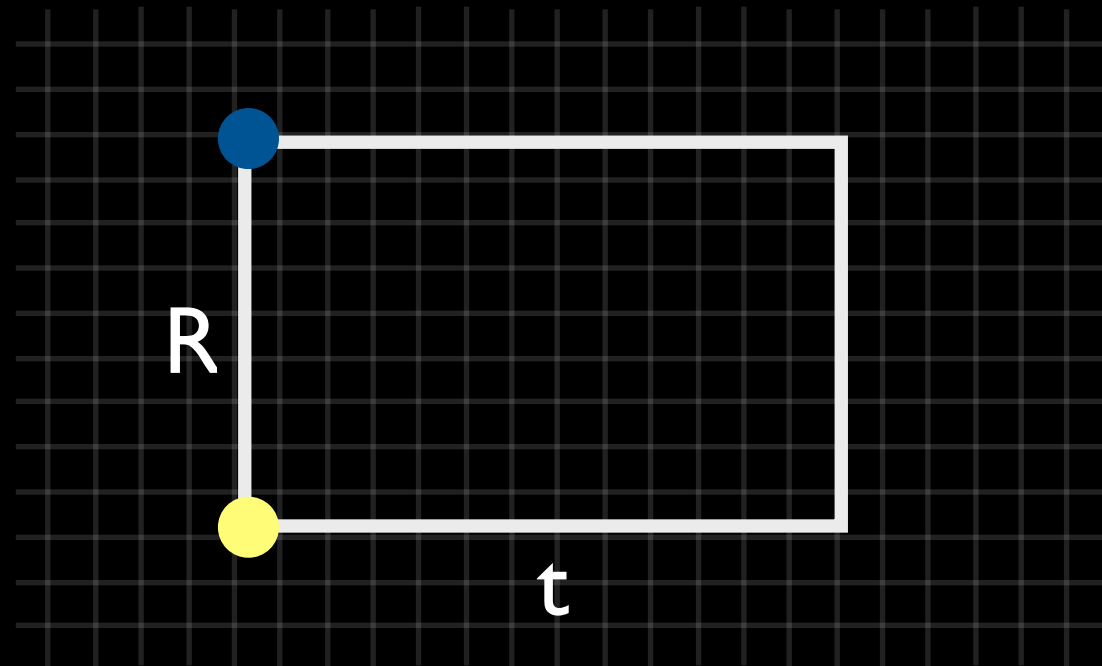


- Screening: evidence for quark-gluon plasma

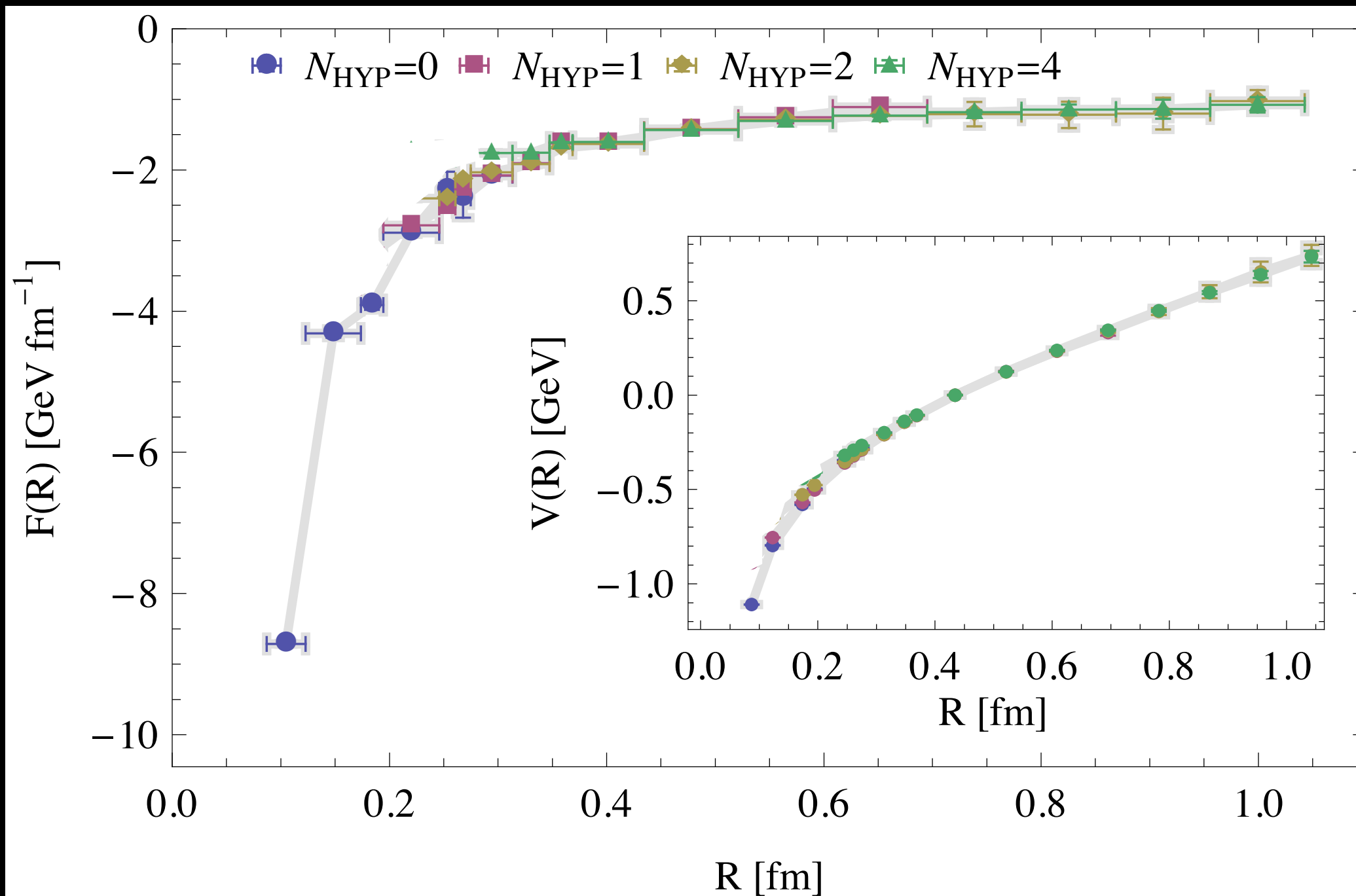
Correlation functions

- Static quark potential

$$C_W(R, t_w, t) = \langle 0 | \sum_{\mathbf{y}, |\mathbf{r}|=R} \mathcal{W}(\mathbf{y} + \mathbf{r}, t; \mathbf{y}, t_w) | 0 \rangle$$
$$\longrightarrow Z \exp[-V(r)(t - t_0)]$$

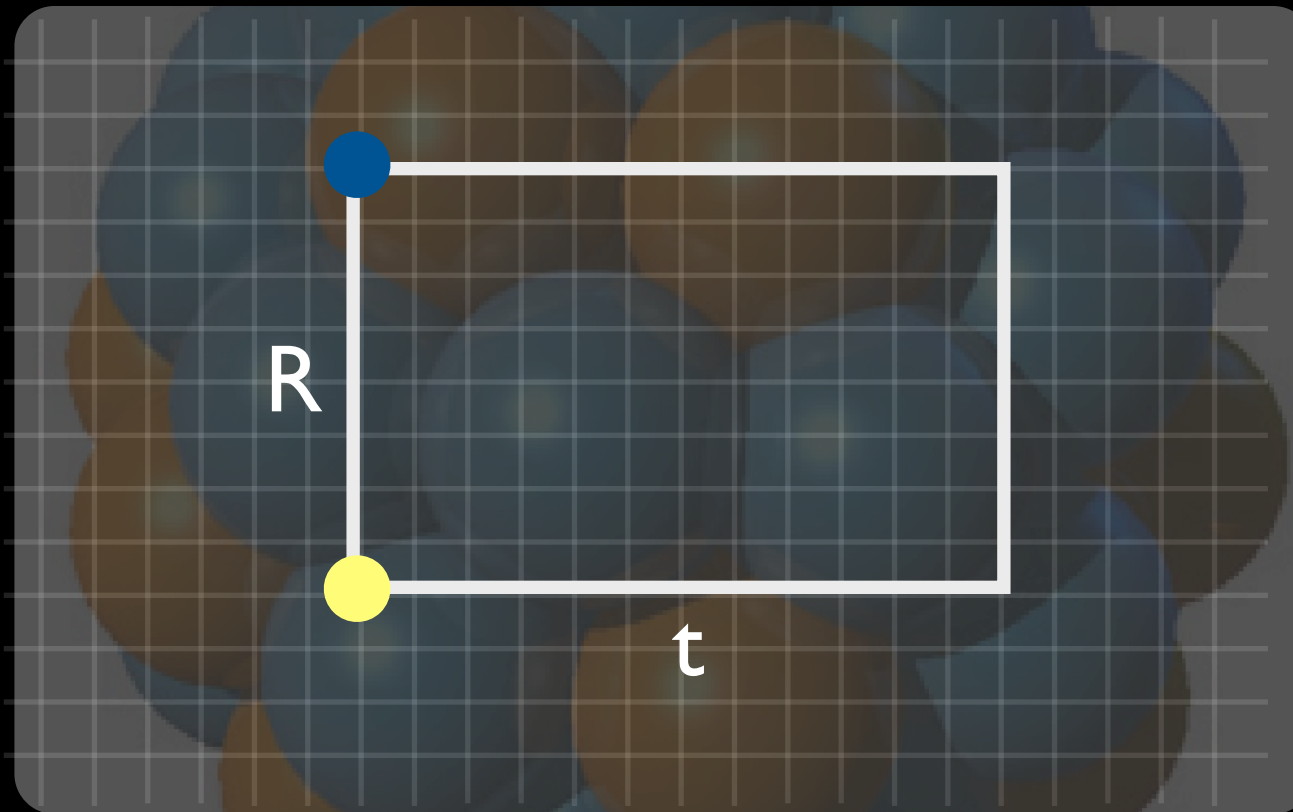


Vacuum $Q\bar{Q}$ force



Color screening

- Static quark potential



- Modified by condensate? Hadronic screening?

Correlation functions

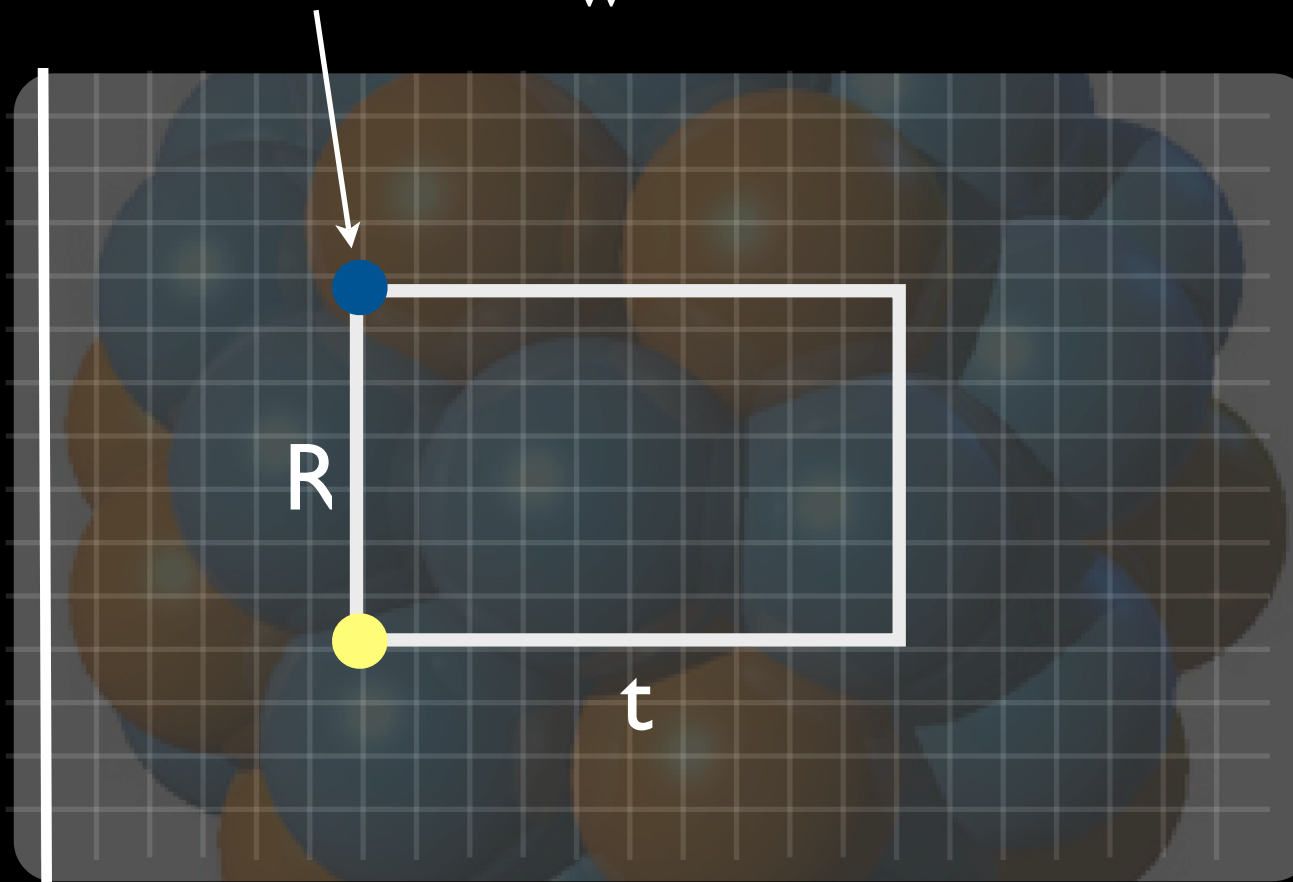
$$C_W(R, t_w, t) = \langle 0 | \sum_{\mathbf{y}, |\mathbf{r}|=R} \mathcal{W}(\mathbf{y} + \mathbf{r}, t; \mathbf{y}, t_w) | 0 \rangle$$
$$\longrightarrow Z \exp[-V(r)(t - t_0)]$$

$$C_n(t_\pi, t) = \langle 0 | \left[\sum_{\mathbf{x}} \chi_{\pi+}(\mathbf{x}, t) \chi_{\pi+}^\dagger(0, t_\pi) \right]^n | 0 \rangle$$
$$\longrightarrow Z' \exp[-E_{n\pi}(t - t_\pi)]$$

$$C_{n,W}(R, t_\pi, t_w, t) = \langle 0 | \left[\sum_{\mathbf{x}} \chi_{\pi+}(\mathbf{x}, t) \chi_{\pi+}^\dagger(0, t_\pi) \right]^n$$
$$\times \sum_{\mathbf{y}, |\mathbf{r}|=R} \mathcal{W}(\mathbf{y} + \mathbf{r}, t; \mathbf{y}, t_w) | 0 \rangle$$

In pictures

Wilson line
source $t=t_w$



$n-\pi$
source $t=t_\pi$

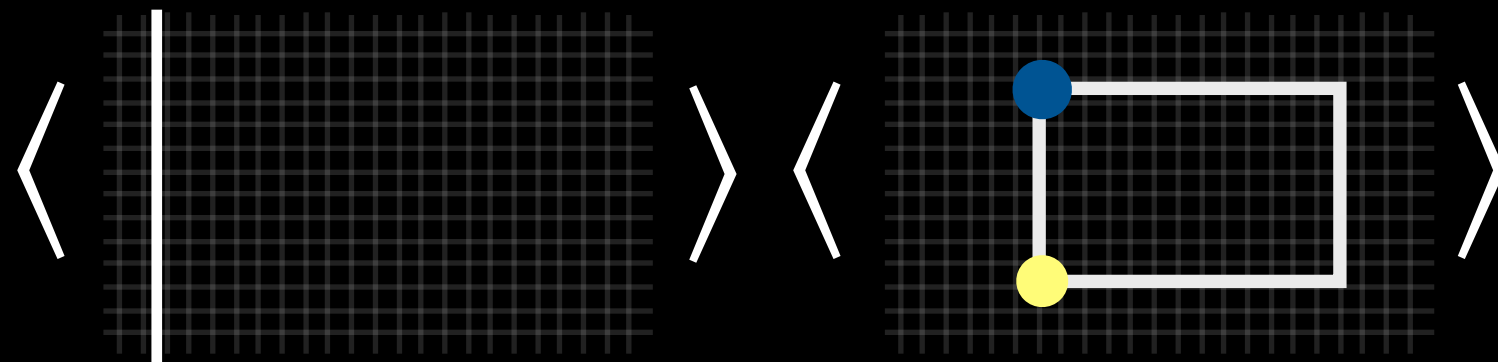
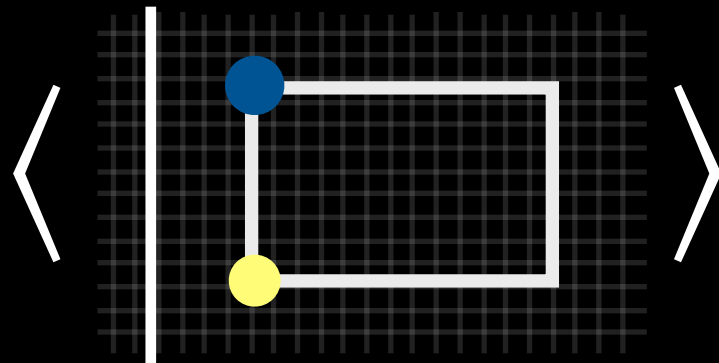
Euclidean time



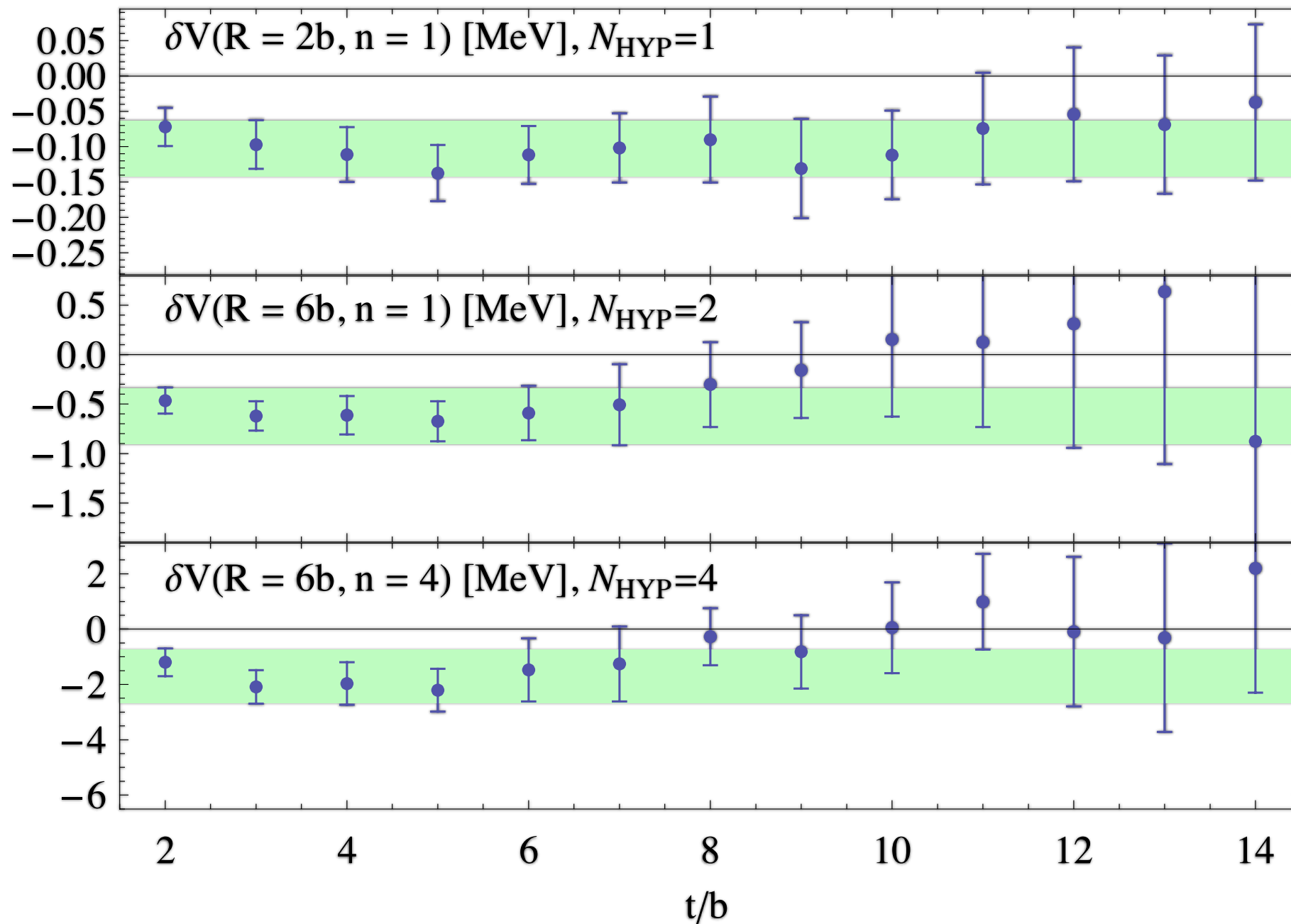
In medium effects

$$G_{n,W}(R, t_\pi, t_w, t) = \frac{C_{n,W}(R, t_\pi, t_w, t)}{C_n(t_\pi, t) C_W(R, t_w, t)}$$

$\longrightarrow \# \exp[-\delta V(r, n)(t - t_0)]$

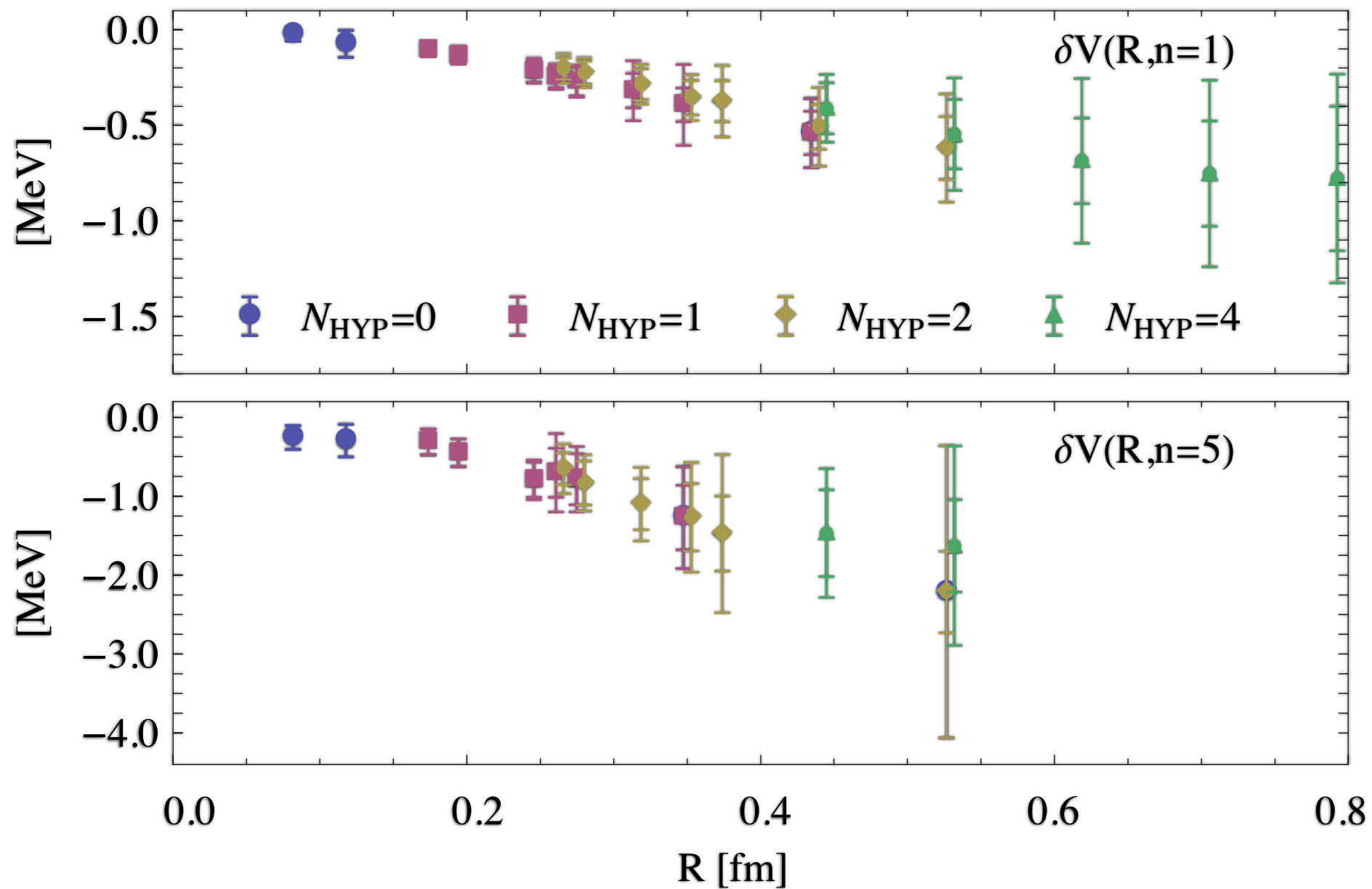


Effective δV plots



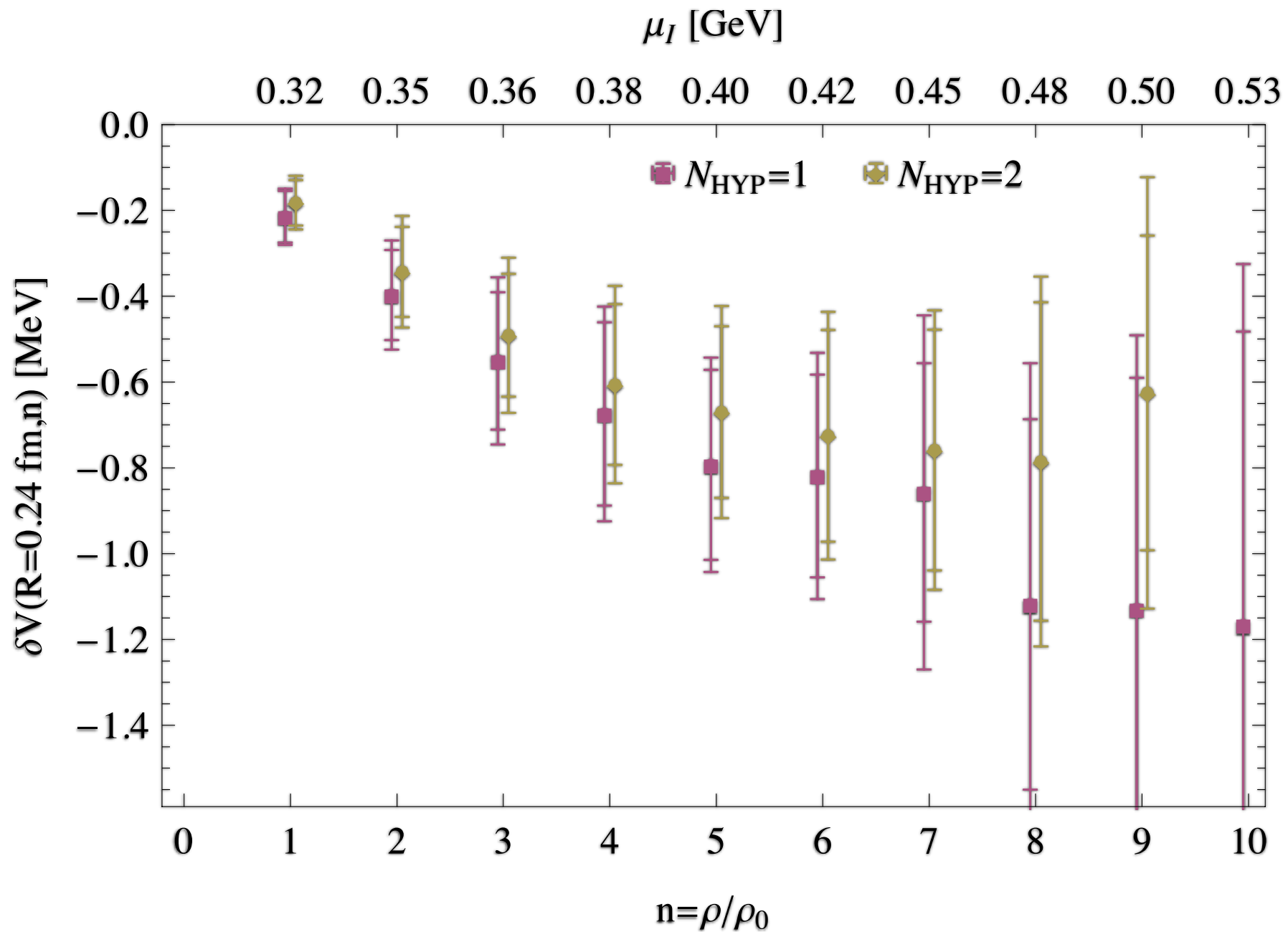
- DWF on MILC: $a=0.09$ fm, $28^3 \times 96$, $m_\pi=318$ MeV

$\delta V(R, n=1 \text{ \& } 5)$



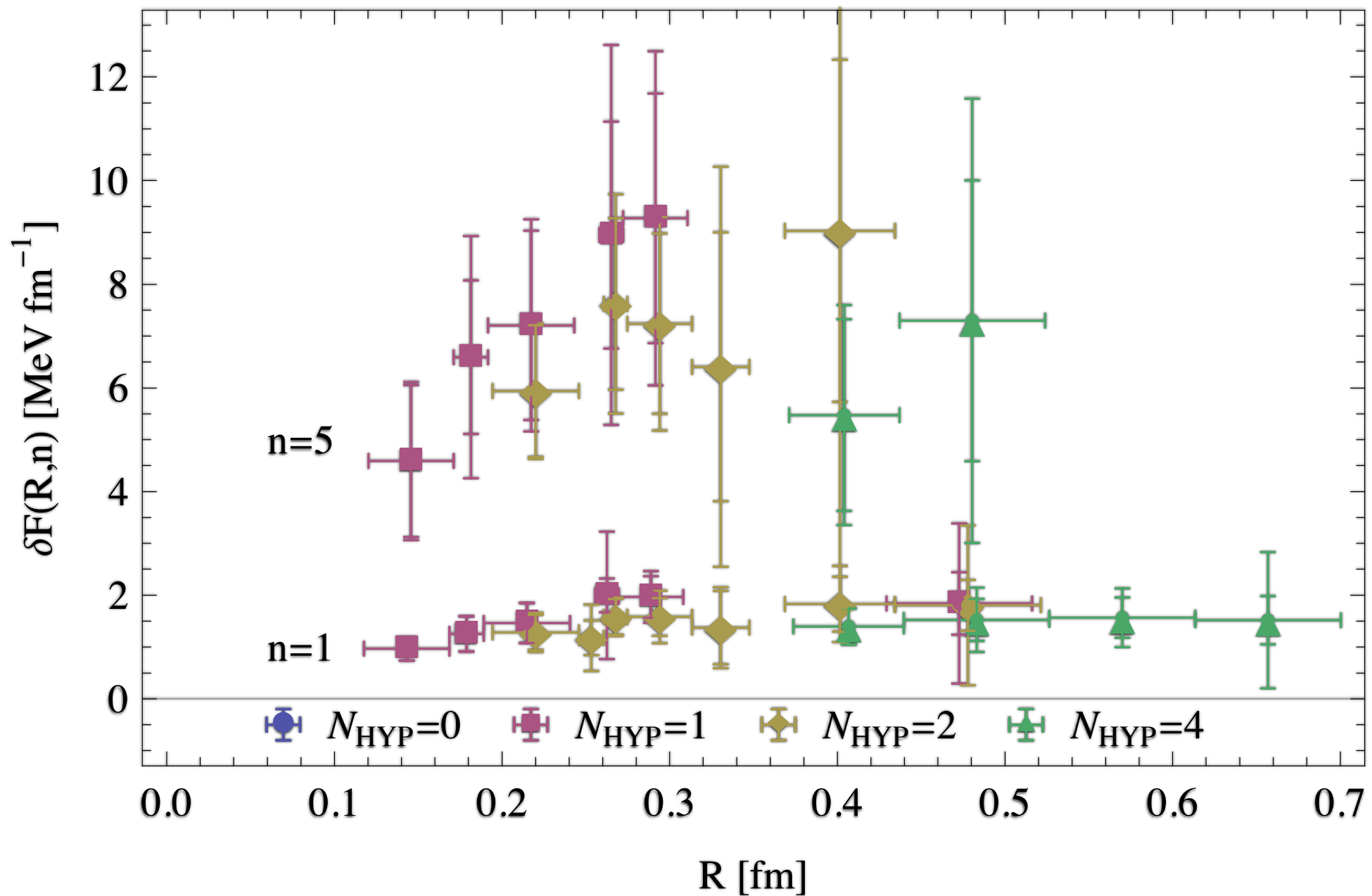
- 100 keV effect at small n , small R

$\delta V(R=0.24\text{fm}, n)$



- Approximately linear in isospin density

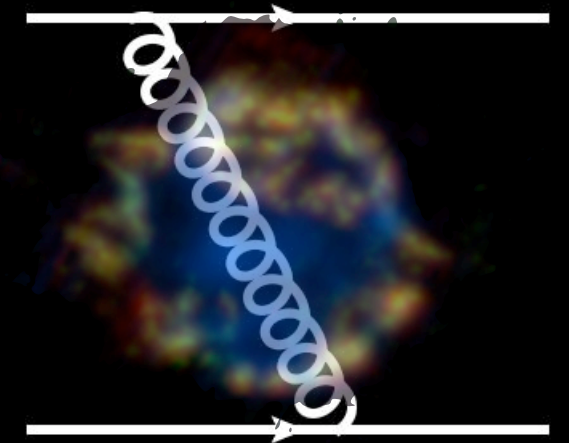
$\delta F(R, n=1 \text{ \& } 5)$



- Constant at large R , $\sim R^2$ at small R ??

Pion screening

- r independent shift in $Q\bar{Q}$ force
- Dielectric medium inside flux tube
- Small effect: $\delta F(n=1)/F = 0.002$ at large R



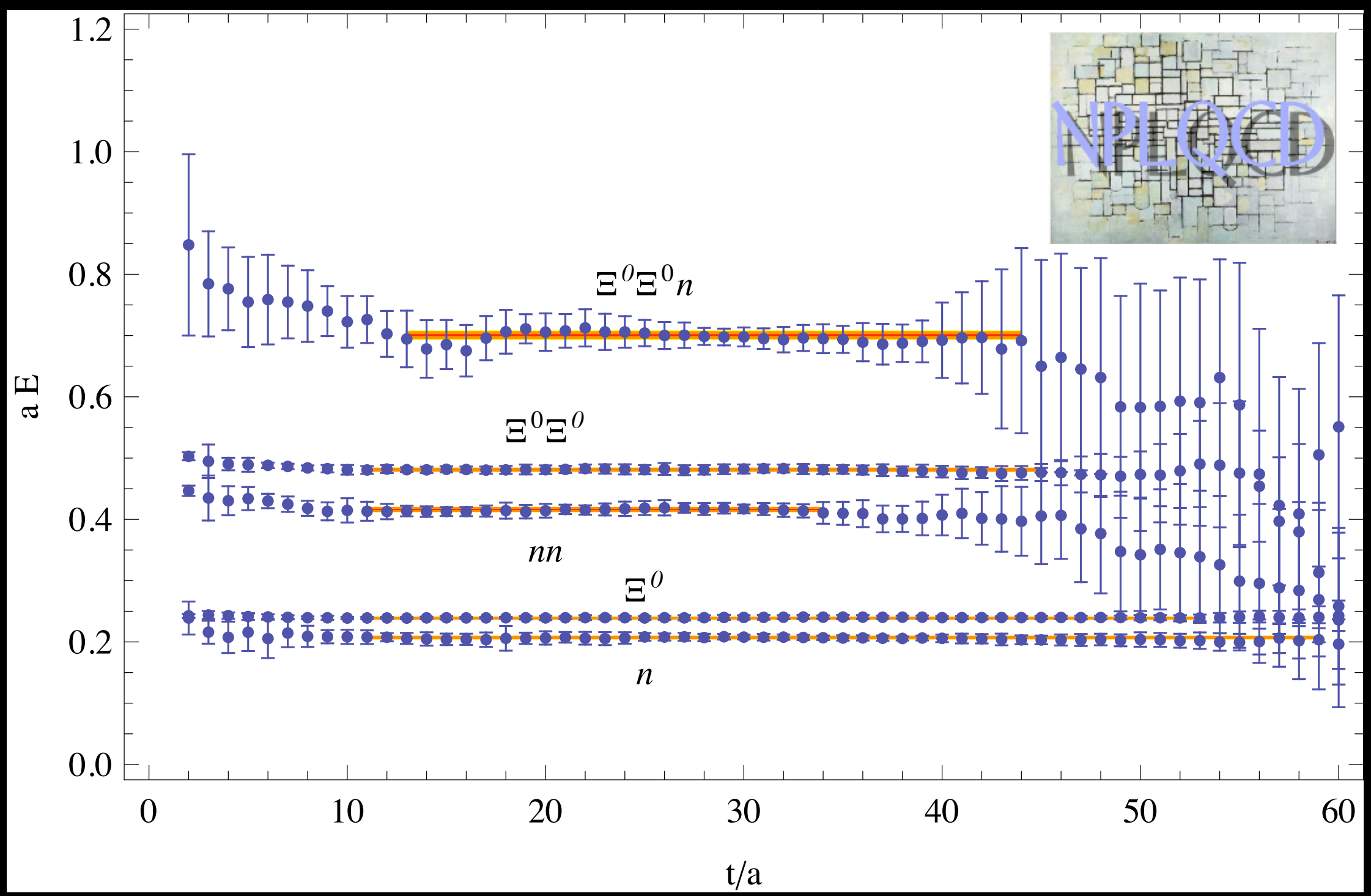
- Relevance to J/ψ suppression @ SPS/RHIC?

Baryons Strike Back

- Just how hard are nuclei?
- Simple three-baryon system: $\Xi^0 \Xi^0 n$
- Unphysical mass, single L and a
- Extreme statistics = 300K measurements

PRELIMINARY

$\Xi^0 \Xi^0 n$: effective energy



Summary



- Numerical meson condensates
 - Two and three meson interactions
 - Equation of state of meson gases
 - Effects: color screening
- Progress can be made in multi-baryon systems



FIN