

THE JET STRUCTURE
of

BAMA-RAY GLAST'S

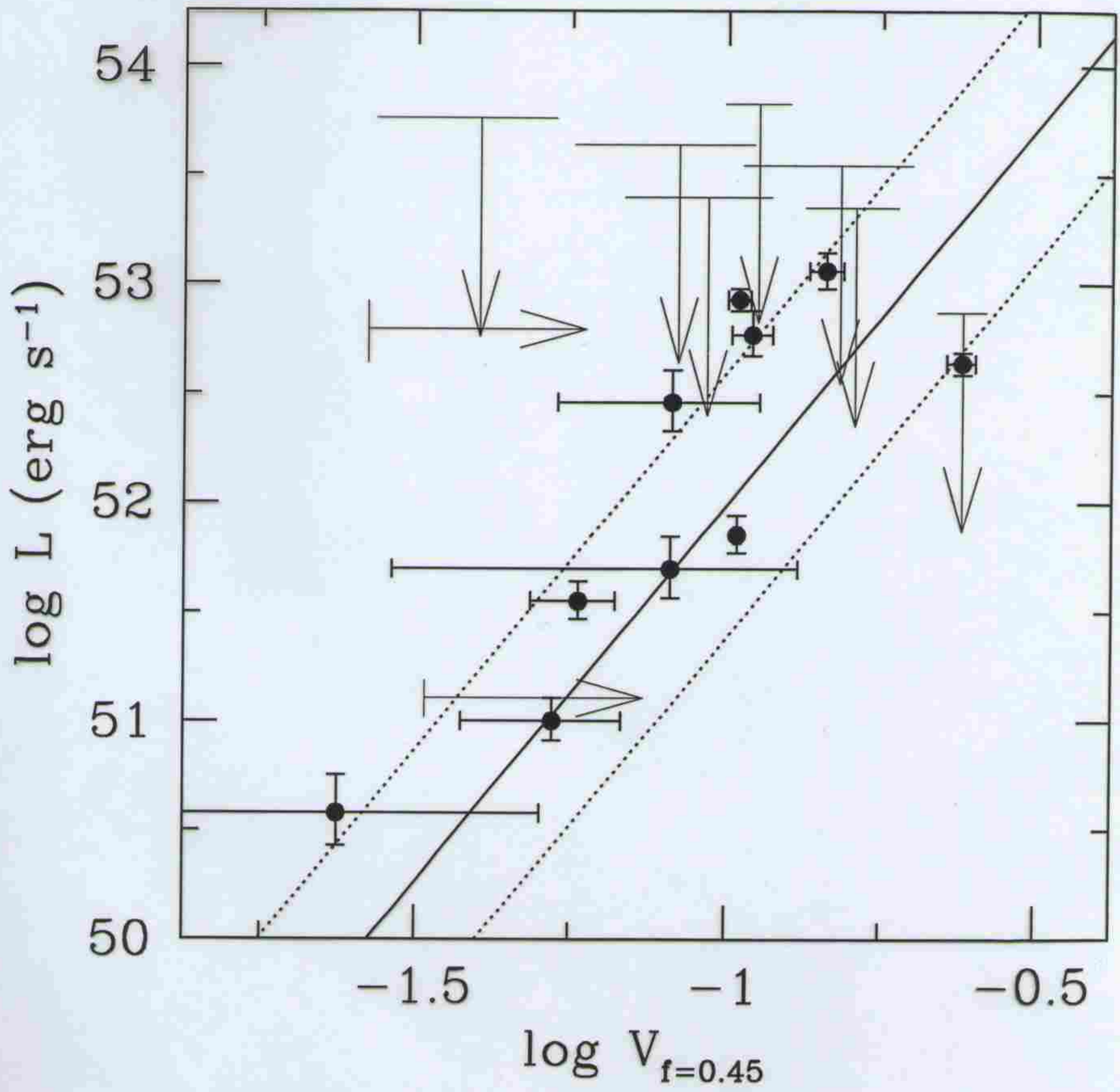
or
(HAVING A GLAST
IN HUNTSVILLE!)

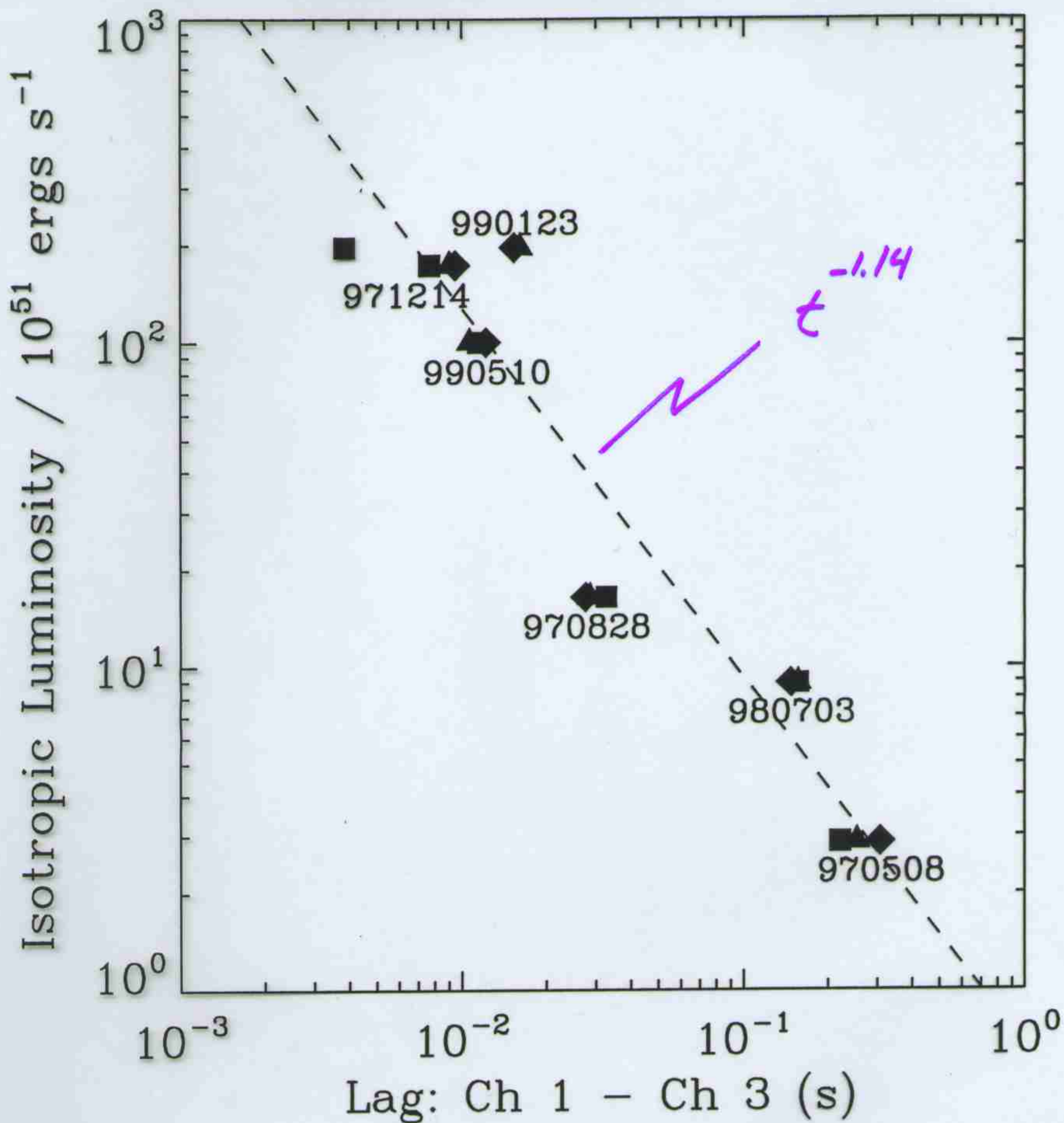
by
Jay Salmonson

Lawrence Livermore National Lab

@ the GLAST SWG 9/12/02

Reichart et al. 2001





Norris, Bonnell & Marani
2000

Lag-Luminosity

Relationship

(Norris et al.)

$$L_{pk} \sim \frac{1}{\Delta t}$$

Doppler

Factor

$$D \equiv \frac{1}{\gamma(1-\beta \cos \theta)(1+z)} \sim \frac{\gamma}{(1+(\gamma \theta)^2)(1+z)}$$

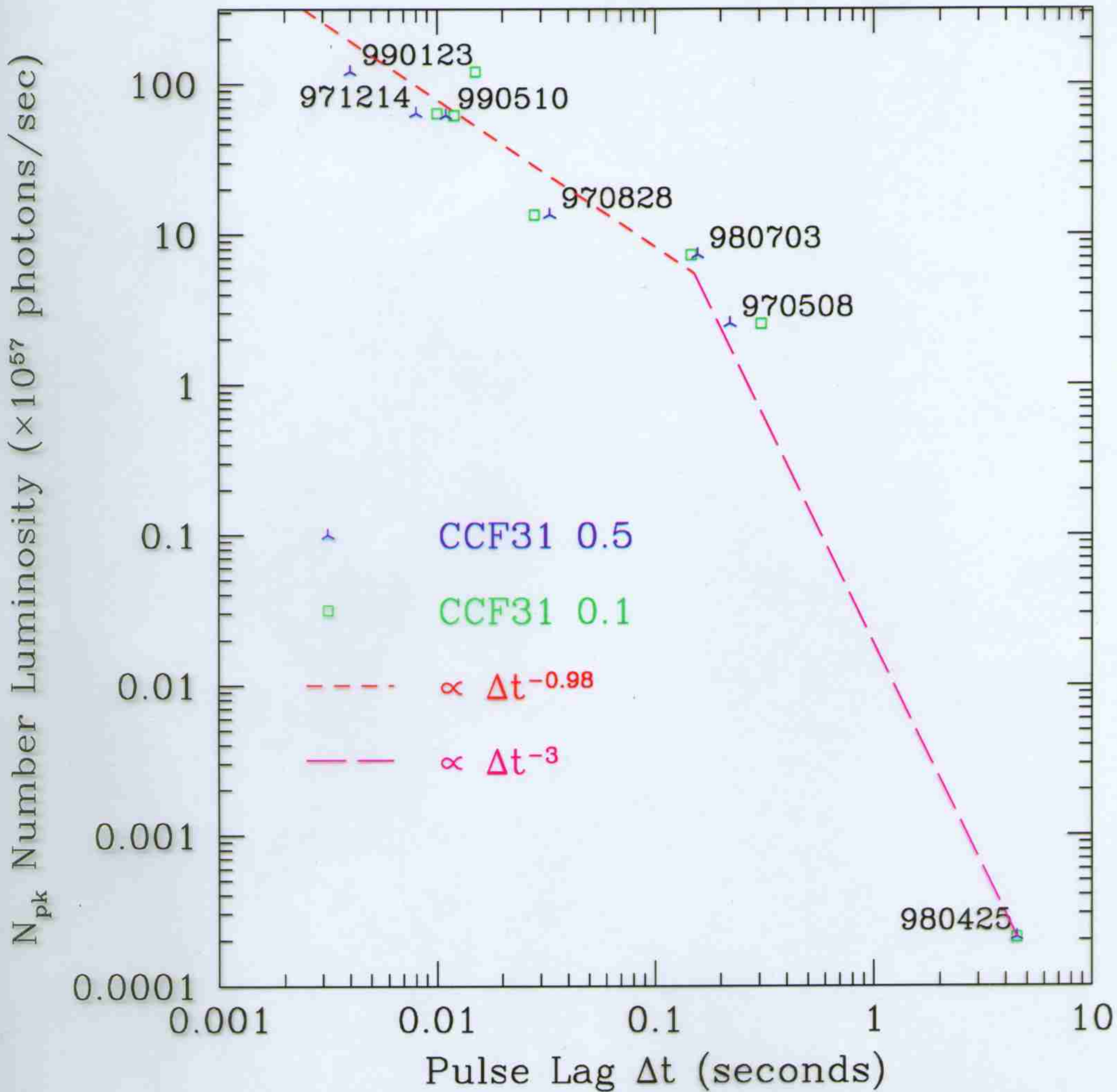
Spectrum $\phi(E) \sim \begin{cases} E^{-\alpha}, & E < E_0 \\ E^{-\beta}, & E > E_0 \end{cases} \Rightarrow E^{-\alpha}$

$$L_{pk} \propto \int_{E_1}^{E_2} E \phi\left(\frac{E}{\delta}\right) dE \propto D^\alpha (\alpha \approx 1)$$

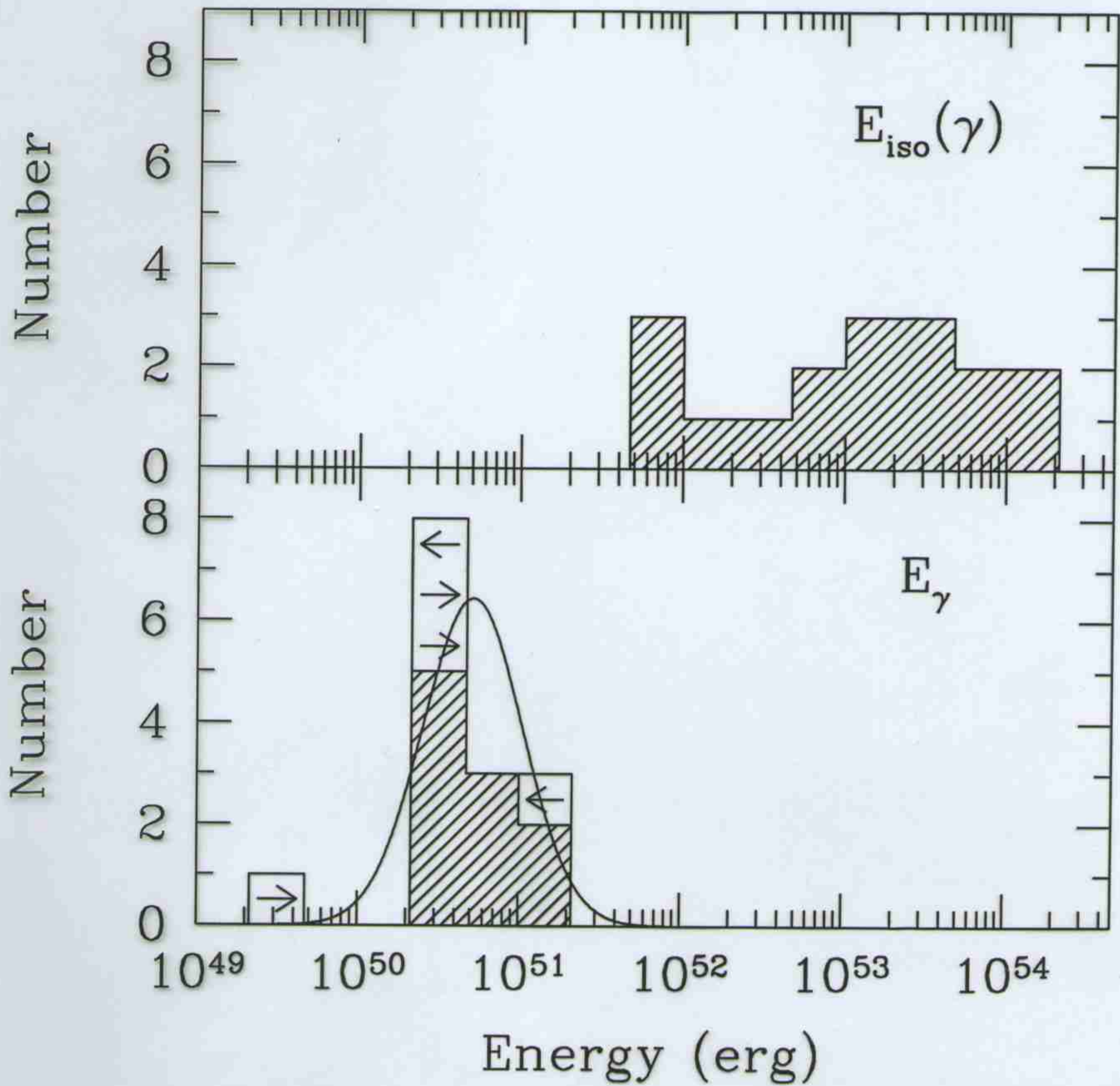
$$\Delta t \propto \frac{\Delta t'}{D}$$

so $L_{pk} \propto \left(\frac{1}{\Delta t}\right)^\alpha (\alpha \approx 1)$

Salmonson 2001

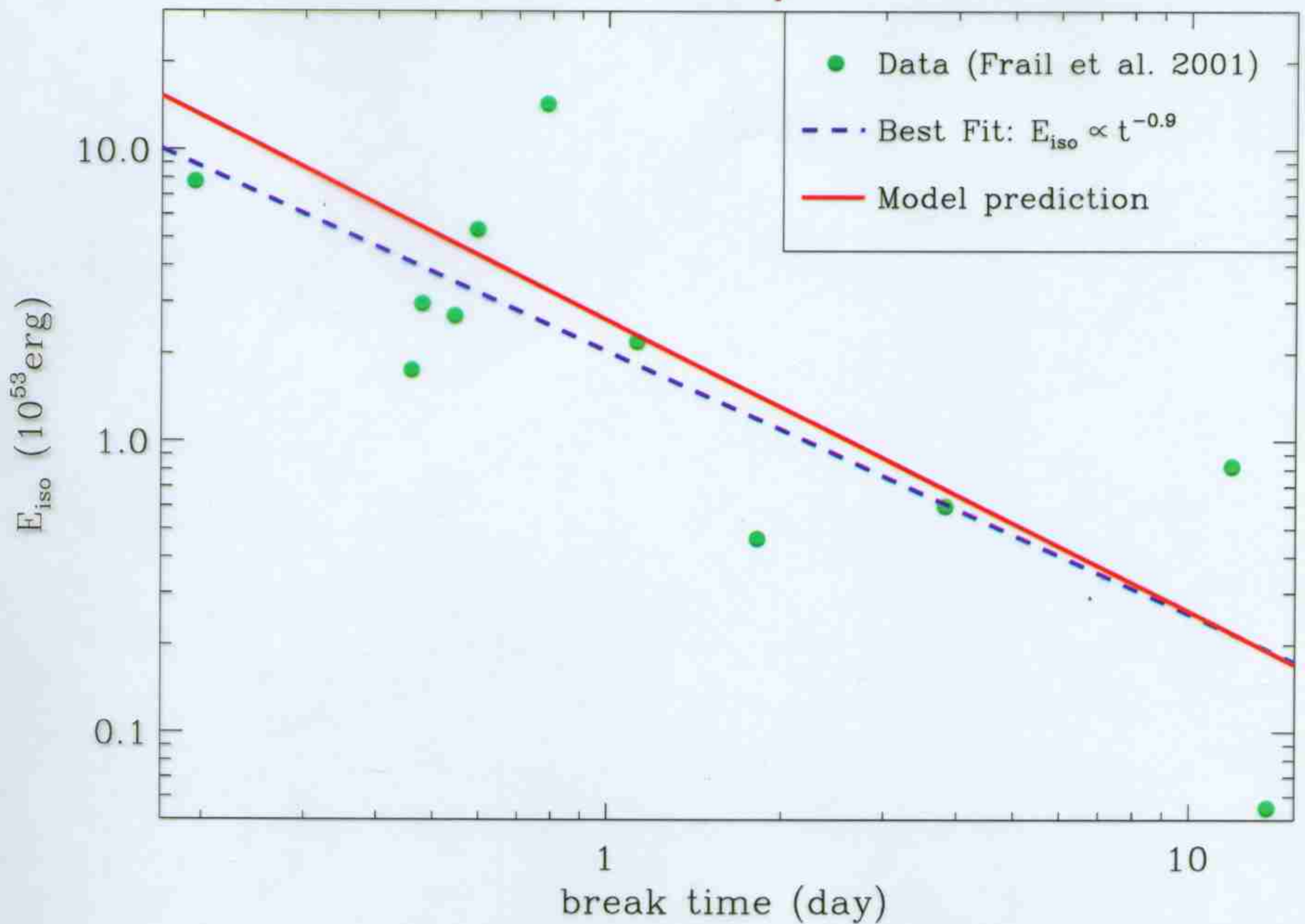


Frail et al. 2001



$$E_\gamma \propto \Theta^2 E_{iso} \approx \text{constant}$$

Rossi, Lazzati & Rees



Frail et al. noticed: $E_{\text{iso}} \propto t_j^{-1}$

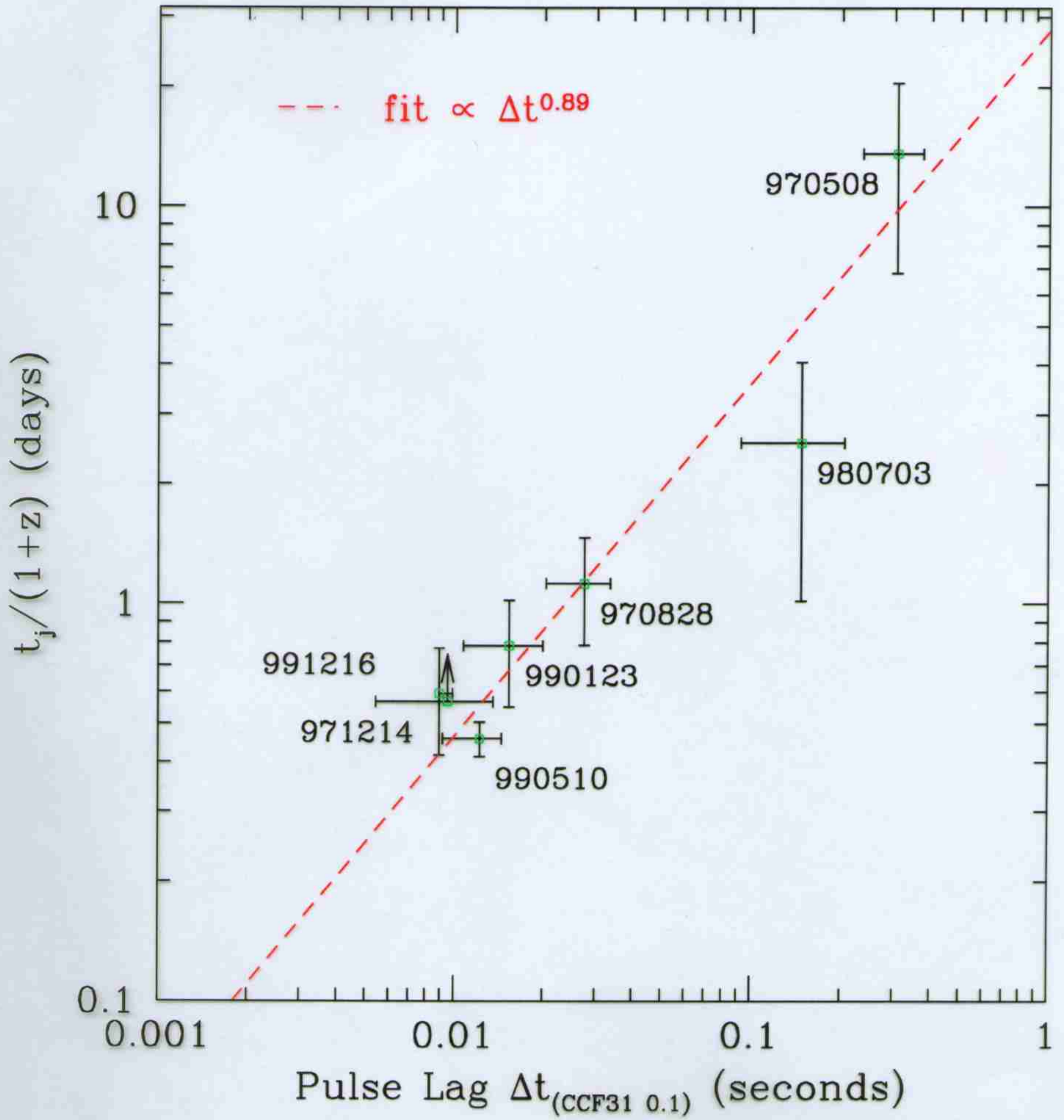
so since $t_j \propto E^{1/3} \Theta^{8/3}$

then $E \sim E^{-1/3} \Theta^{8/3}$

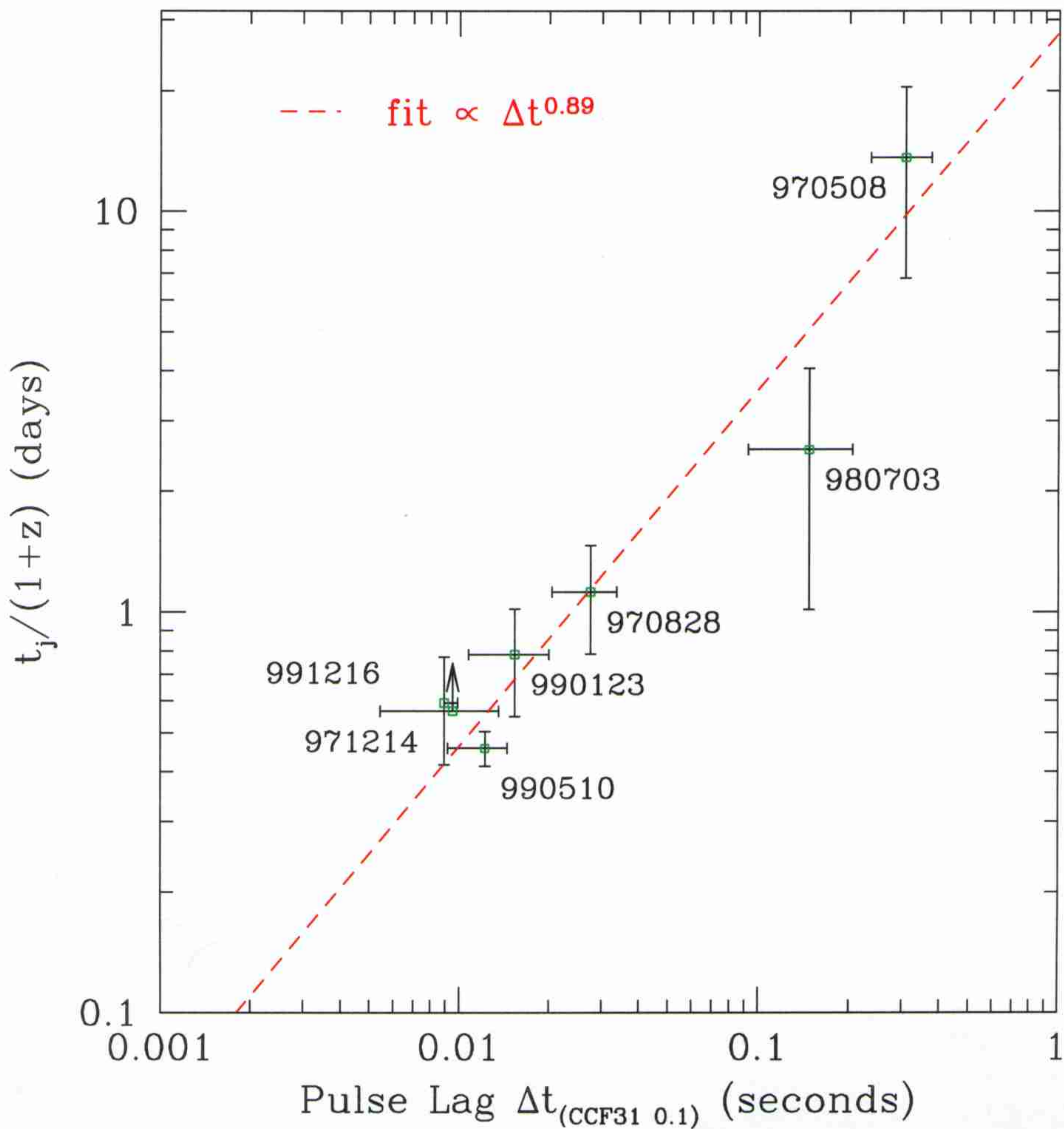
or

$$E_{\text{iso}} \propto \Theta^{-2}$$

Salmonson & Galama 2002



Salmonson & Galama 2002



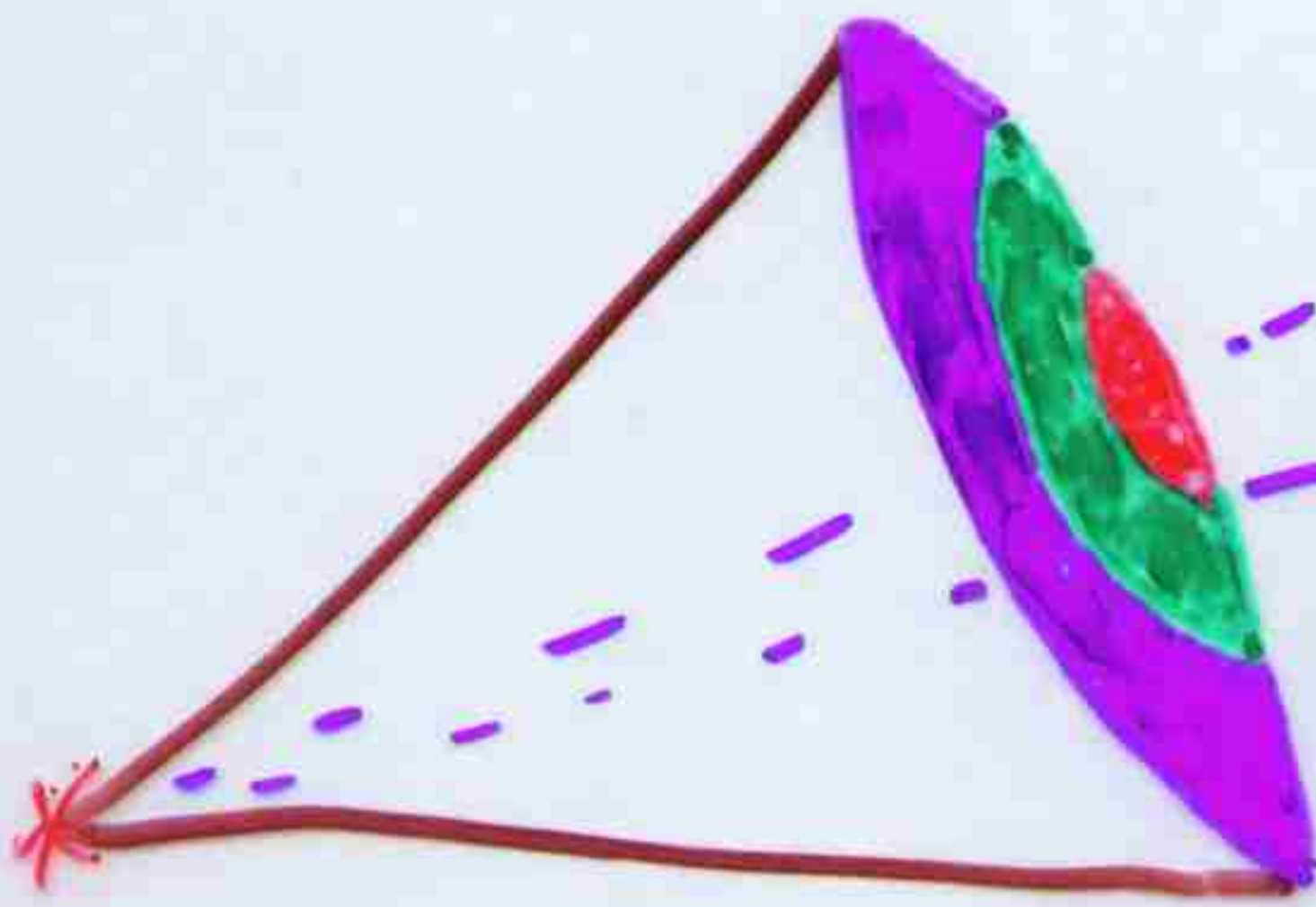
Two Paradigms

Many bursts:



$$\underline{\epsilon \propto \theta_0^{-2}}$$

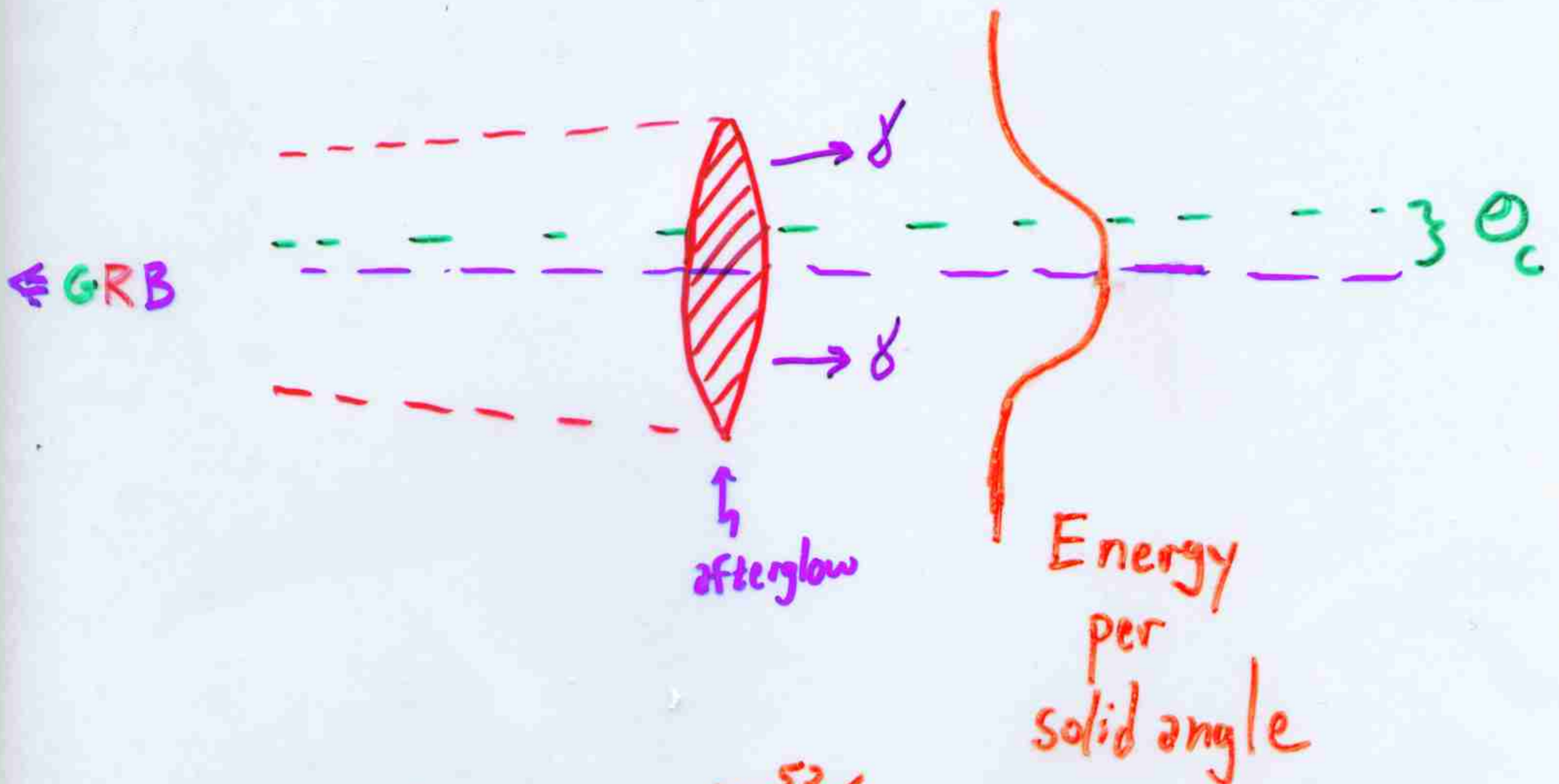
Universal burst:



$$\epsilon \propto \theta_0^{-2}$$

$$\underline{\epsilon \propto \theta_0^{-2}}$$

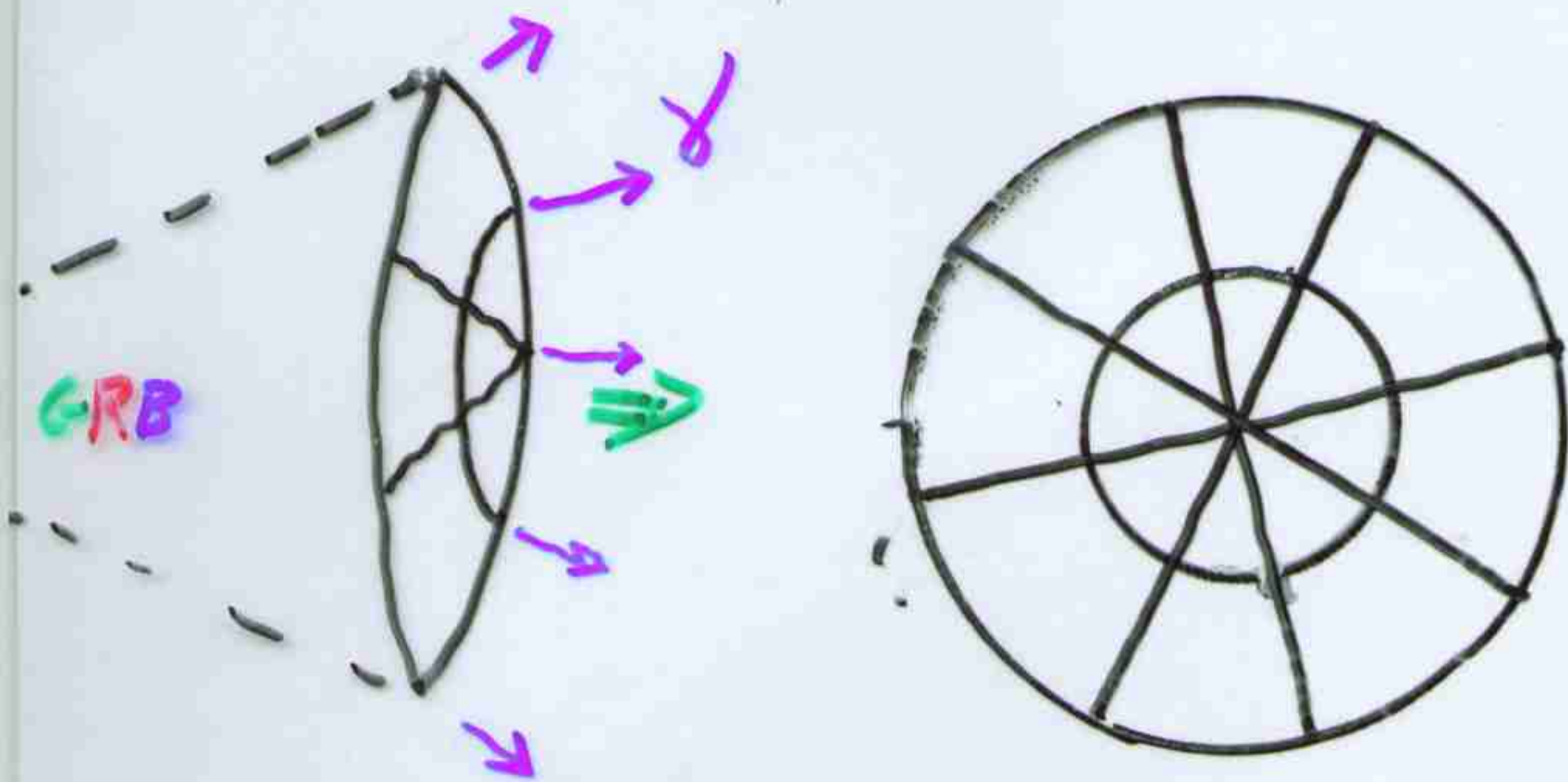
Structured Jet



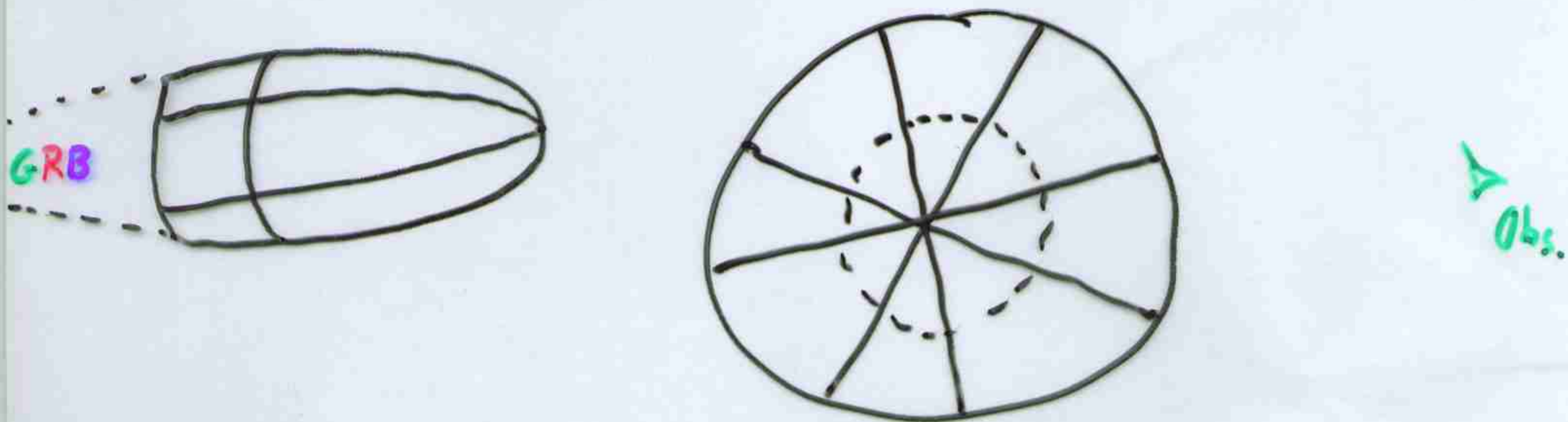
$$\xi(\theta) = \frac{10^{52}/4\pi}{1 + (\theta/\theta_0)^2} \propto \theta^{-2}$$

$$\chi(\theta) = \Gamma_0 = 100$$

Discretize jet surface in θ & ϕ :



Stuff moving toward the observer evolves faster than stuff moving at an angle:



Fireball Dynamics

An afterglow of initial mass, M_0 , and Lorentz factor, Γ_0 , sweeps into the ISM, it decelerates according to conservation of energy and radial momentum which give the following evolution equations (Paczynski & Rhoads '93 ApJ, 418,L6).

$$f = \frac{M(R)}{M_0} = \frac{1}{M_0} \int_0^R r^2 \Omega_m \rho_{ISM} dr$$

swept up mass fraction

$$\Omega_m = \pi(\theta_0 + v_{\perp} t' / R)^2$$

jet opening solid angle

$$\beta = \frac{v}{c} = \sqrt{1 - 1/\Gamma^2}$$

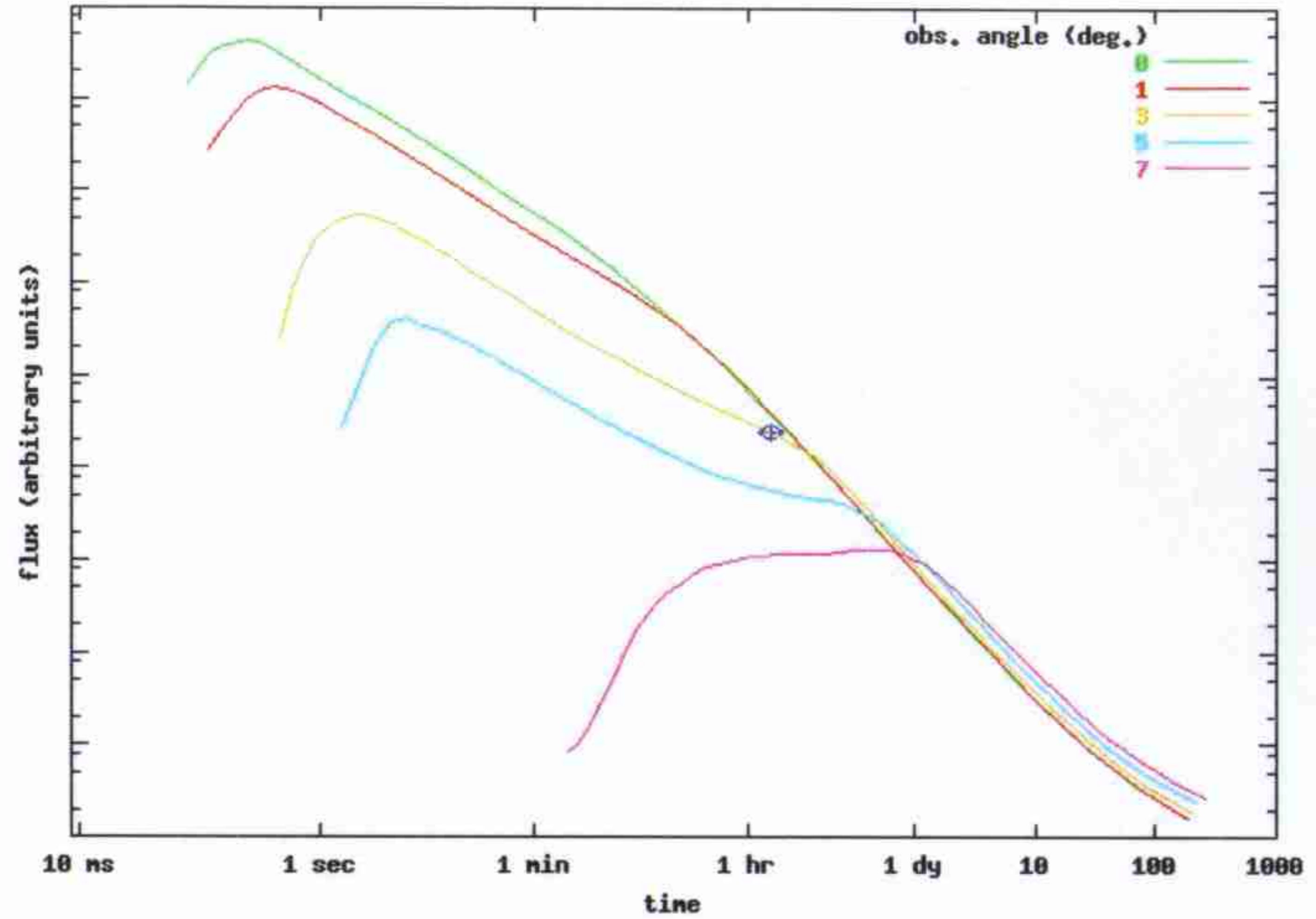
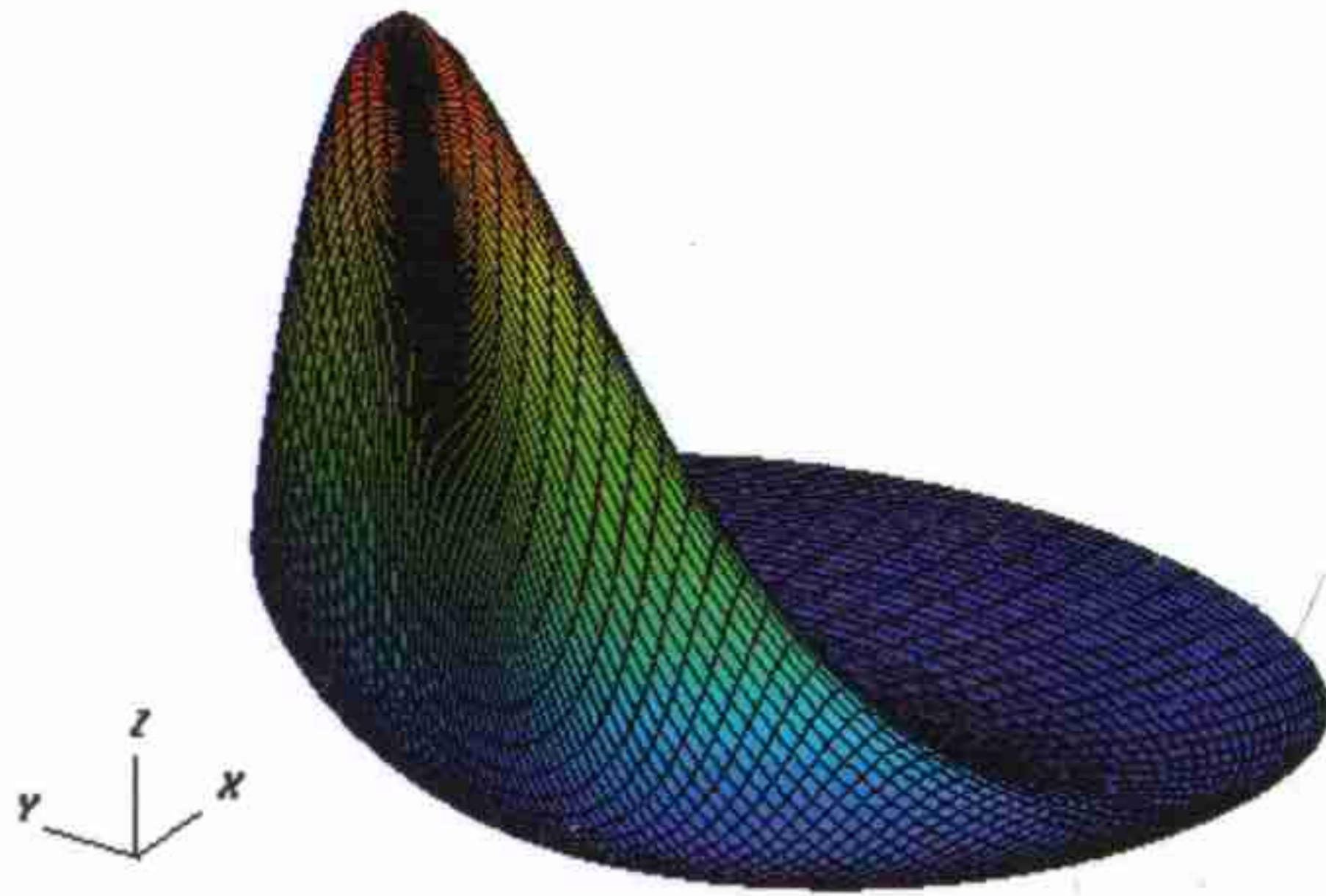
velocity

$$\Gamma = \frac{\Gamma_0 + f}{\sqrt{1 + 2\Gamma_0 f + f^2}}$$

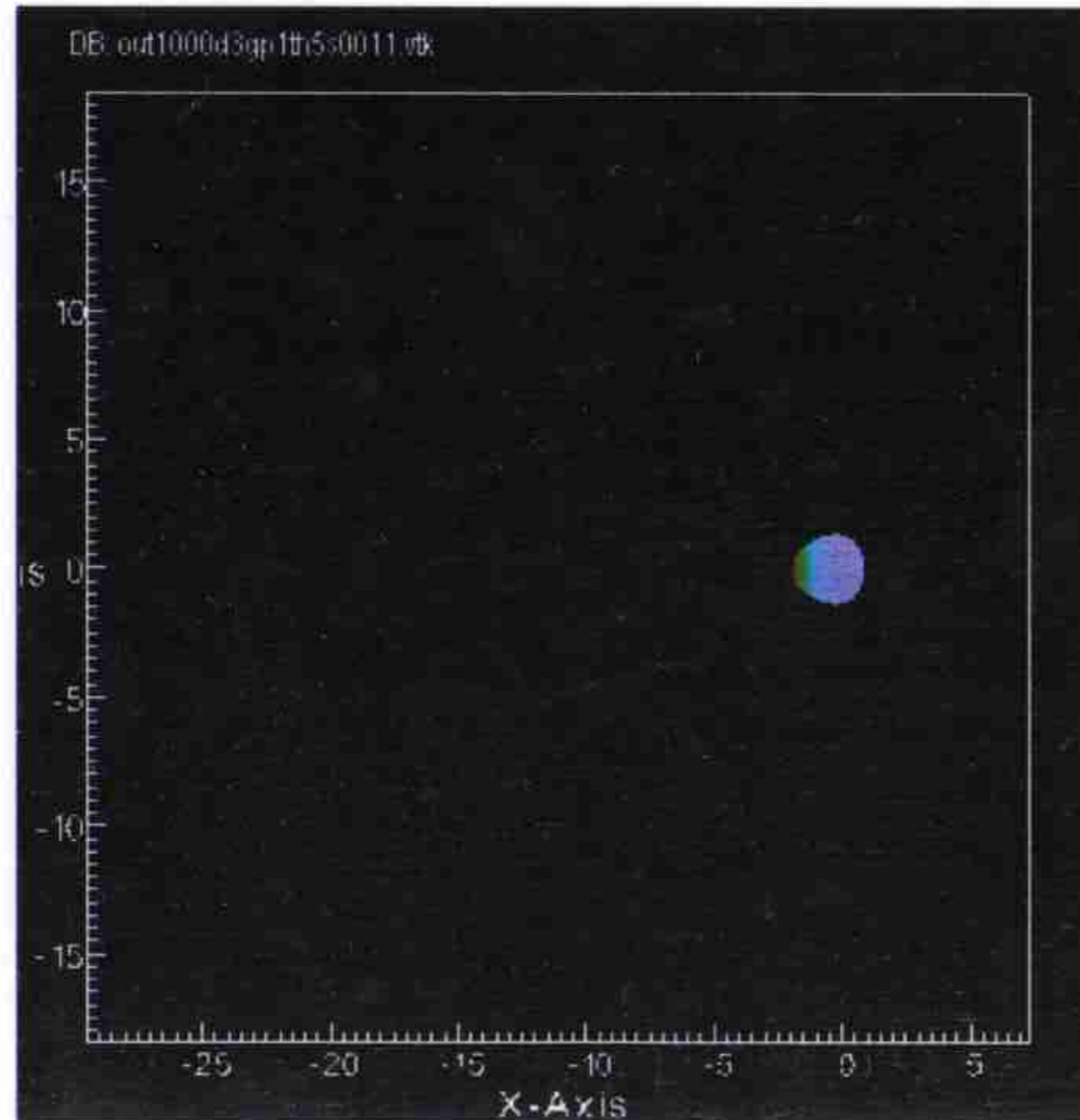
Lorentz factor

$$t' = \int_0^R \frac{dr}{\Gamma \beta c} = \int_0^R \frac{dr}{U}$$

proper time

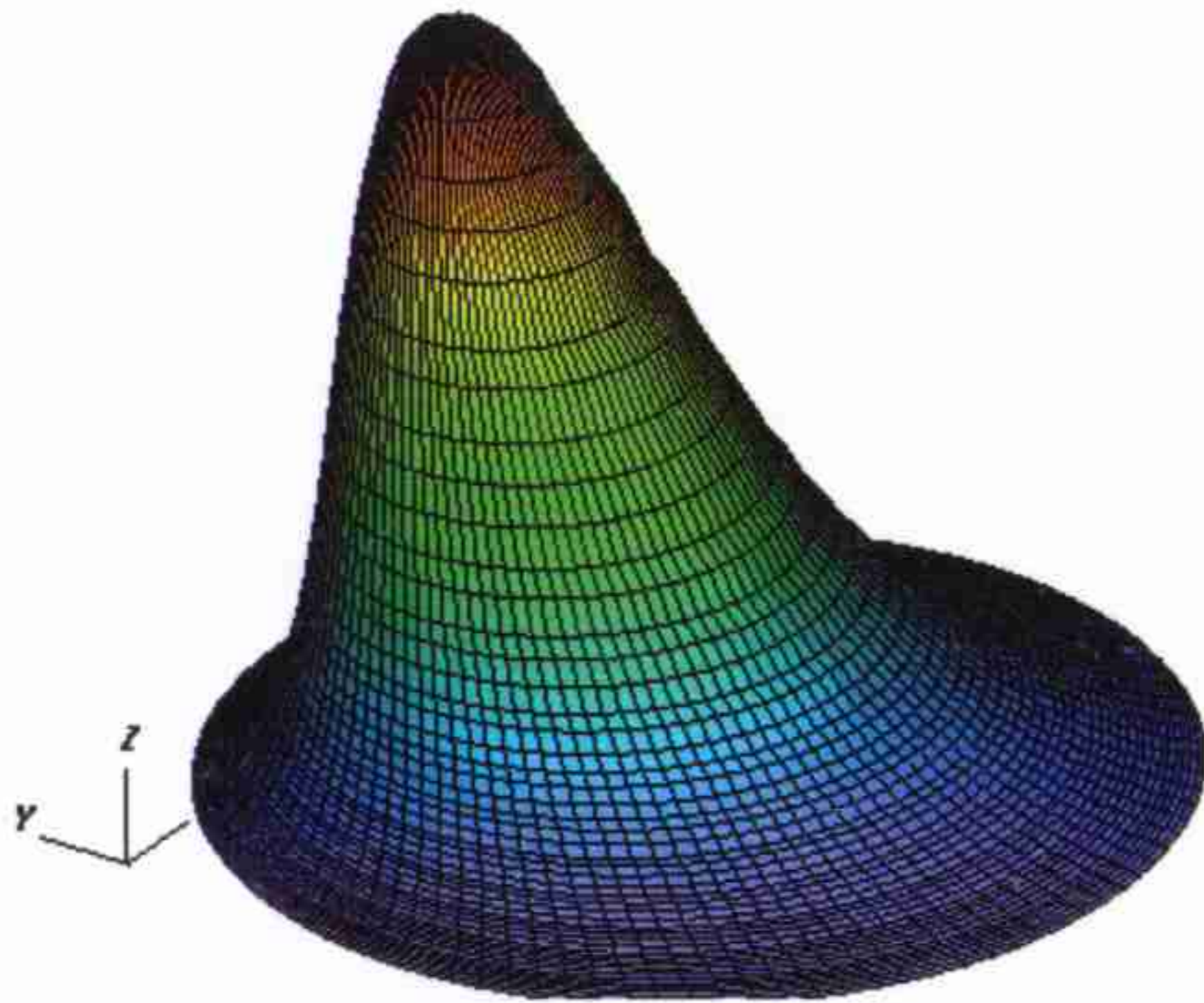


Above: the pole of the jet is almost aligned with the flux peak. Once the peak passes by the pole, the jet will break (upper right).

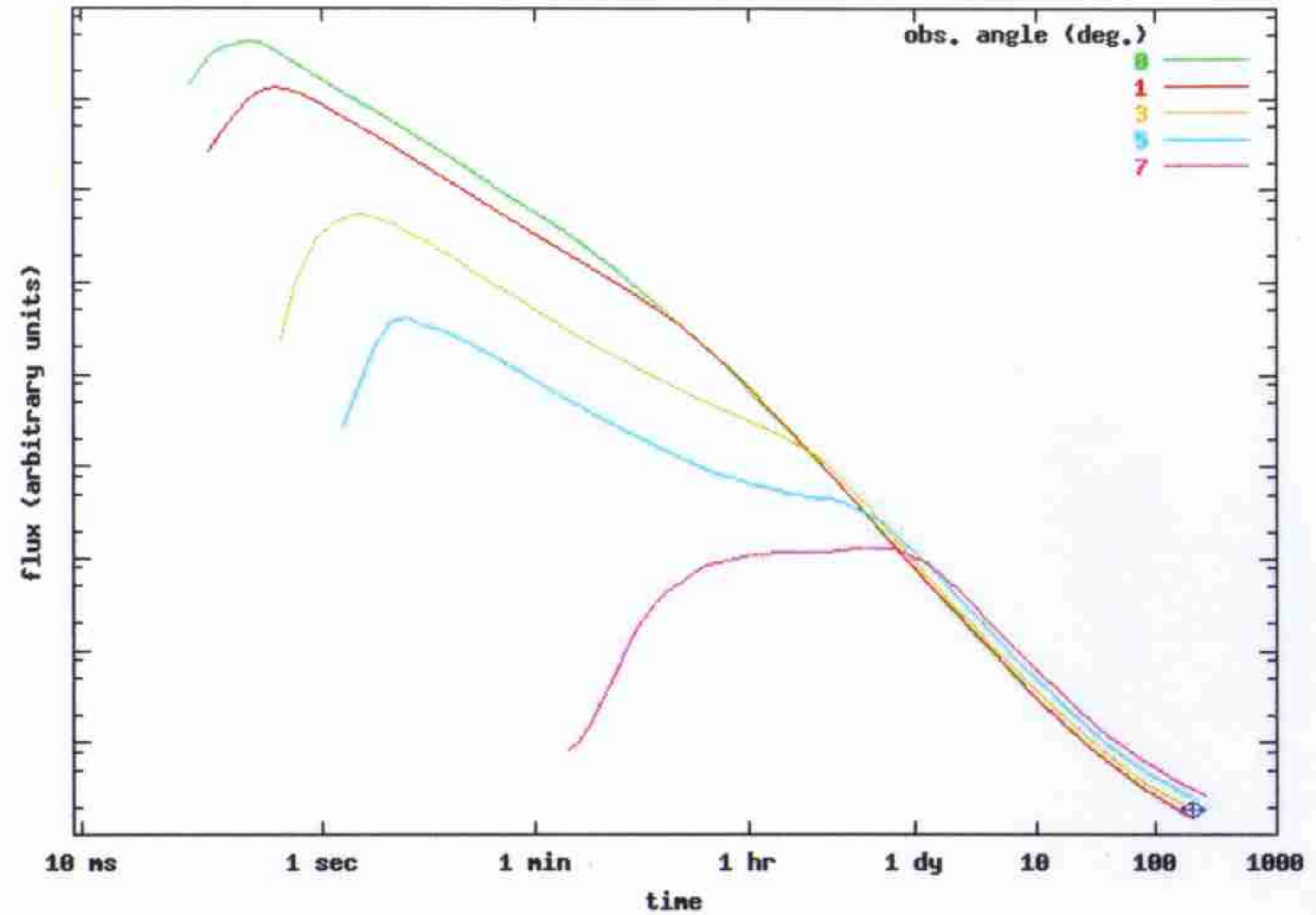


Left: a "view" of what the afterglow would look like at this time. X-Y scale is in light-days.

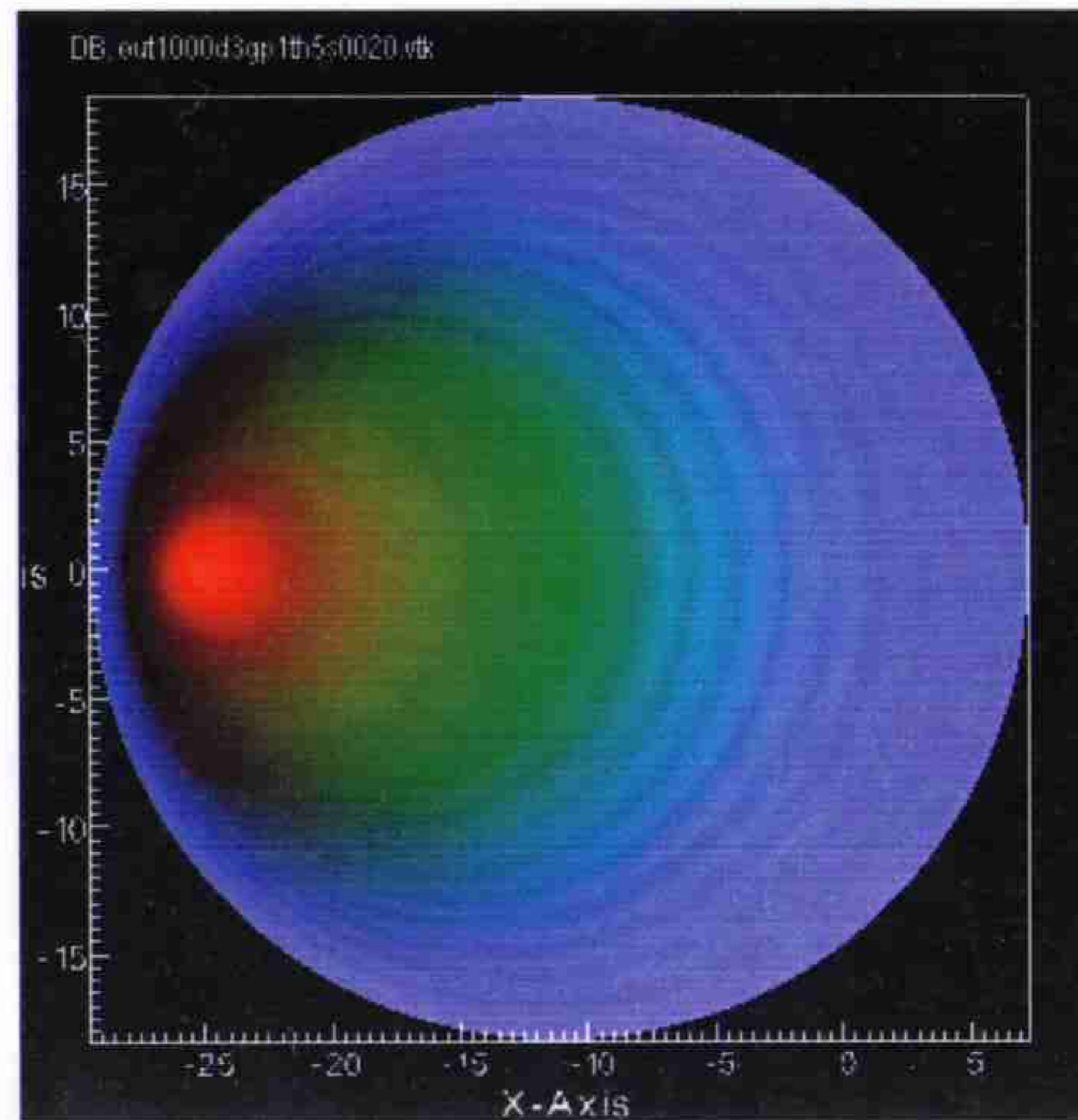
DB: out1000d3gp 1th5s0020.vtk
Cycle: 20



Observer angle = 3, gamma = 1000, gamma_perp = 1 theta0 = 5 J.D.Salmonson 2002

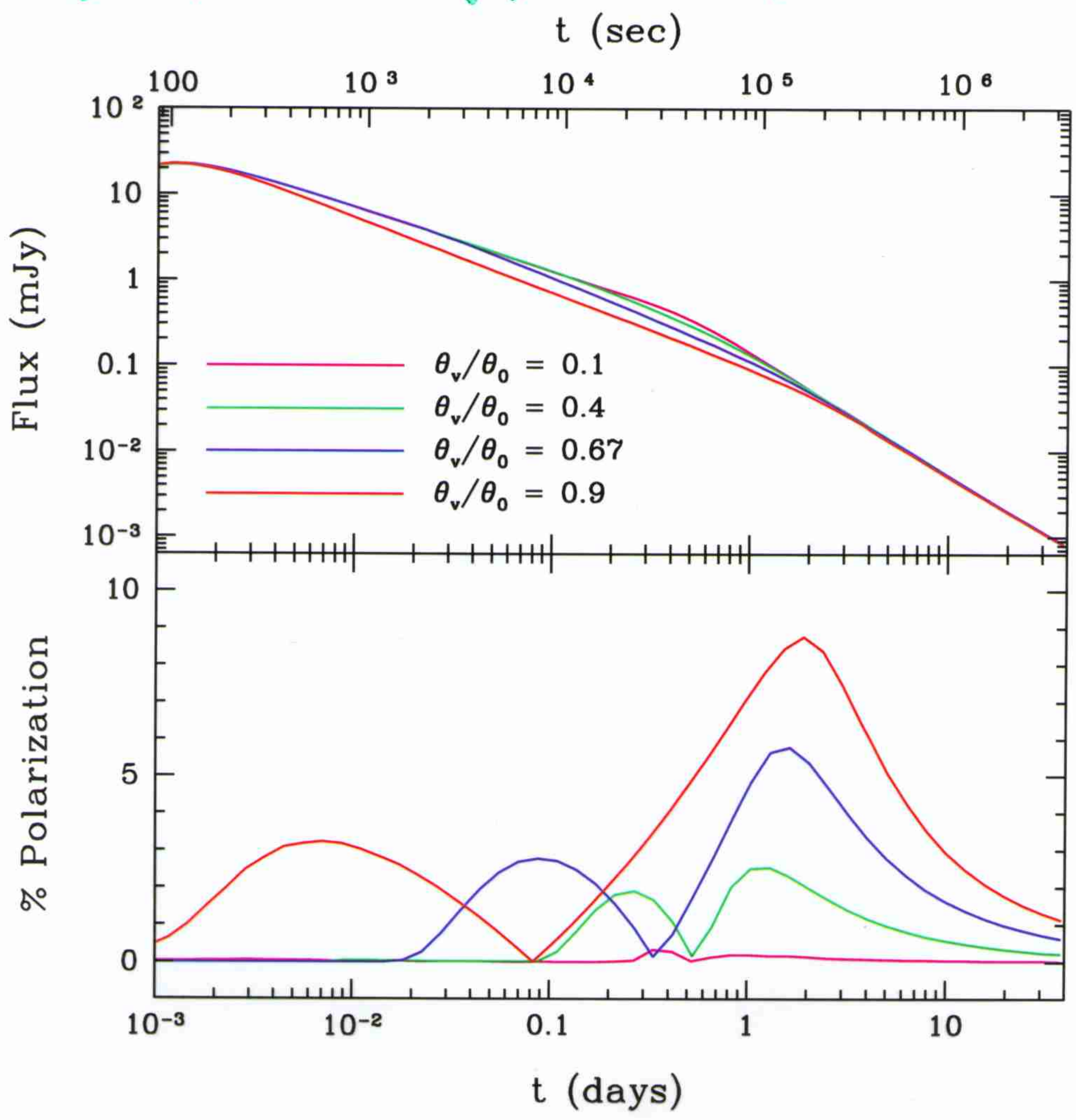


Above: the "sombbrero" shape of the structured jet is apparent as the Lorentz factor has decreased to almost unity.



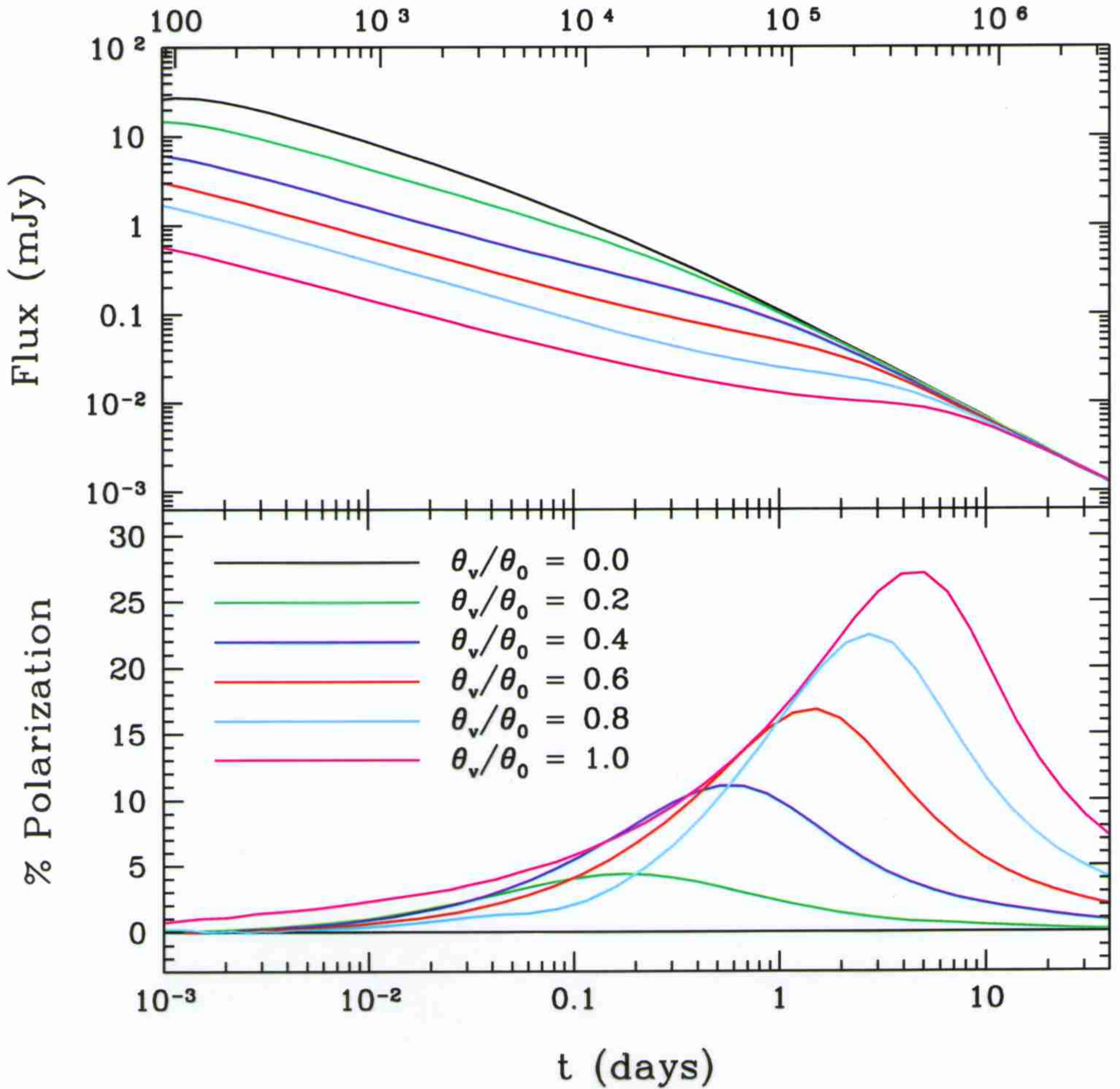
Left: a "view" of what the afterglow would look like at this time. X-Y scale is in light-days.

Homogeneous, $\dot{E} = 10^{52}$ ergs, $\theta_0 = 5^\circ$, $\Gamma_0 = 100$



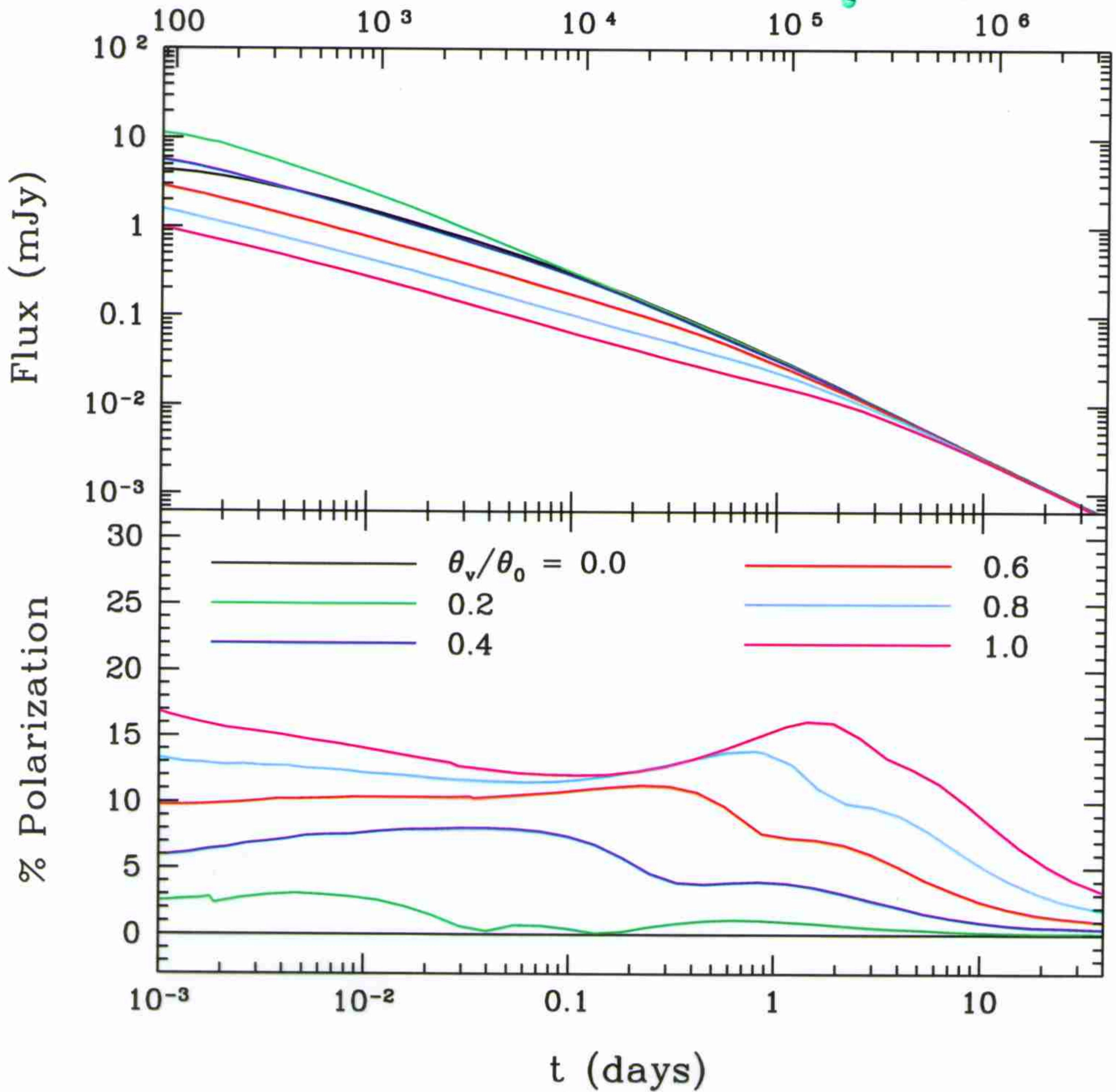
$R_K = 0$

Structured, $\dot{E} = 10^{52}$ ergs, $\theta_c = 3^\circ$, $\theta_o = 15^\circ$
 $\Gamma_o = 100$



$$R_K = 0$$

Structured, $E = 10^{52}$ ergs, $\theta_e = 3^\circ$, $\theta_0 = 15^\circ$
 $\Gamma_0 = 100$



$R_k = 0.1$

THE GLAST TIMES

High Energies still
a mystery

▶ $E_{pk} \sim \text{constant}$ \therefore
doesn't follow other trends

▶ More data will help diagnose
lags

