

Winter Sampling of Shallow Firn Air at the South Pole to Understand Processes Affecting Firn Atmospheric Histories and Ice Core Gas Records

JEFFREY P. SEVERINGHAUS

Scripps Institution of Oceanography, University of California, San Diego, La Jolla, CA 92093-0244

In January and July 1998 a cooperative project was carried out at the South Pole CMDL Clean Air Laboratory to sample air from the top 15 m of the snow (firn). The goal was to better understand gas fractionation processes that may affect the atmospheric gas histories obtained from the air in the deep firn column. Previous studies in summer have shown a marked isotopic anomaly in the top 10 m of the firn, driven by thermal fractionation in response to the strong temperature gradient in the top few meters. No samples had been taken in winter, however, so Scripps Institution of Oceanography (SIO) was curious to know (1) if the opposite sign of the temperature gradient produced an opposite isotope anomaly, as expected, (2) if the air in the firn convects spontaneously during winter in response to colder air overlying warmer air, which could bias the ice core record of gas isotopes, and (3) if adsorption and desorption on snow grains could explain anomalies in some halogenated species. For comparison a summer sampling was done first. Both data sets show the expected isotope fractionation because of temperature gradients. The winter data reveal a robust isotope anomaly of opposite sign, as expected, that exceeds a diffusion-only model prediction driven by Automatic Weather Station surface temperature data. No evidence for vigorous gas mixing because of convection was found, although the misfit of the data to the model may suggest slow downward motion or a sampling artifact. It appears unlikely that rectification of seasonality biases the deep firn or ice core record at this site, although this study is just a "snapshot" and may have missed episodic events. No evidence for adsorption or desorption was found.

INTRODUCTION

The porous and permeable layer of snow on top of polar ice sheets is typically 50-100 m thick and is known as the firn layer. At the base of this layer the firn is continuously transformed into impermeable ice. The air in firn mixes slowly with the atmosphere, primarily by molecular diffusion [Schwander *et al.*, 1988]. Consequently, the mean age of the air (defined as the time elapsed since the air crossed the surface) increases downward in the firn and reaches several decades at sites with thick firn [Schwander *et al.*, 1993].

A number of recent studies have exploited this fact to reconstruct atmospheric concentration histories of various gases, such as halocarbons [Butler *et al.*, 1999], O₂/N₂ ratios and N₂O [Battle *et al.*, 1996], and ¹³C of CO₂ [Francey *et al.*, 1999]. Other recent studies have taken advantage of the fact that temperature gradients in the firn arising from rapid climate variation leave an isotope record of abrupt climate change in

trapped air in ice cores [Severinghaus *et al.*, 1998; Leuenberger *et al.*, 1999; Severinghaus and Brook, 1999; Lang *et al.*, 1999]. At many polar sites, surface temperature varies by >30°C seasonally. This creates strong temperature gradients in the top few meters of the firn due to the fact that the firn at ~10 m depth closely maintains the mean annual temperature at all times [Paterson, 1969]. Theory predicts that these gradients will lead to a characteristic pattern of thermally fractionated gases in the upper ~15 m of the firn column. A model of heat and gas diffusion in the firn enables a precise numerical prediction of the magnitude and shape of this pattern, when forced with the known surface temperature history. The model has no adjustable parameters and thus serves as a stringent test of our understanding of the physical processes governing gas-isotope ratios in shallow firn:

$$\frac{\partial C}{\partial t} = \frac{\partial}{\partial z} \left(D(z, T) \left[\frac{\partial C}{\partial z} - \frac{\Delta m g}{RT} + \Omega \frac{dT}{dz} \right] \right) \quad (1)$$

where

- C isotope delta value (e.g., $\delta^{15}\text{N}$)
- t time
- z depth
- D effective molecular diffusivity of gas in porous snow
- T temperature, K
- Δm mass difference between isotopes
- g gravitational acceleration
- R gas constant
- Ω thermal diffusion sensitivity, ‰ K⁻¹

A separate model of heat delivers temperatures to the gas model [Alley and Koci, 1990]. Comparison of the modeled temperature profiles with several measurements (taken from thermistors buried in the borehole) shows good agreement (Figure 1a).

SAMPLING AND ANALYSIS

Because of the difficulty of handling samples outdoors in winter, gas sampling was arranged to be done from indoors. Using a hand coring tool, on January 18, 1998, a 15-m deep borehole located 42 m due grid E of the NE corner of the Clean Air Building was drilled. Tubes with stainless steel intake screens on the ends were placed in the hole at desired depths, and

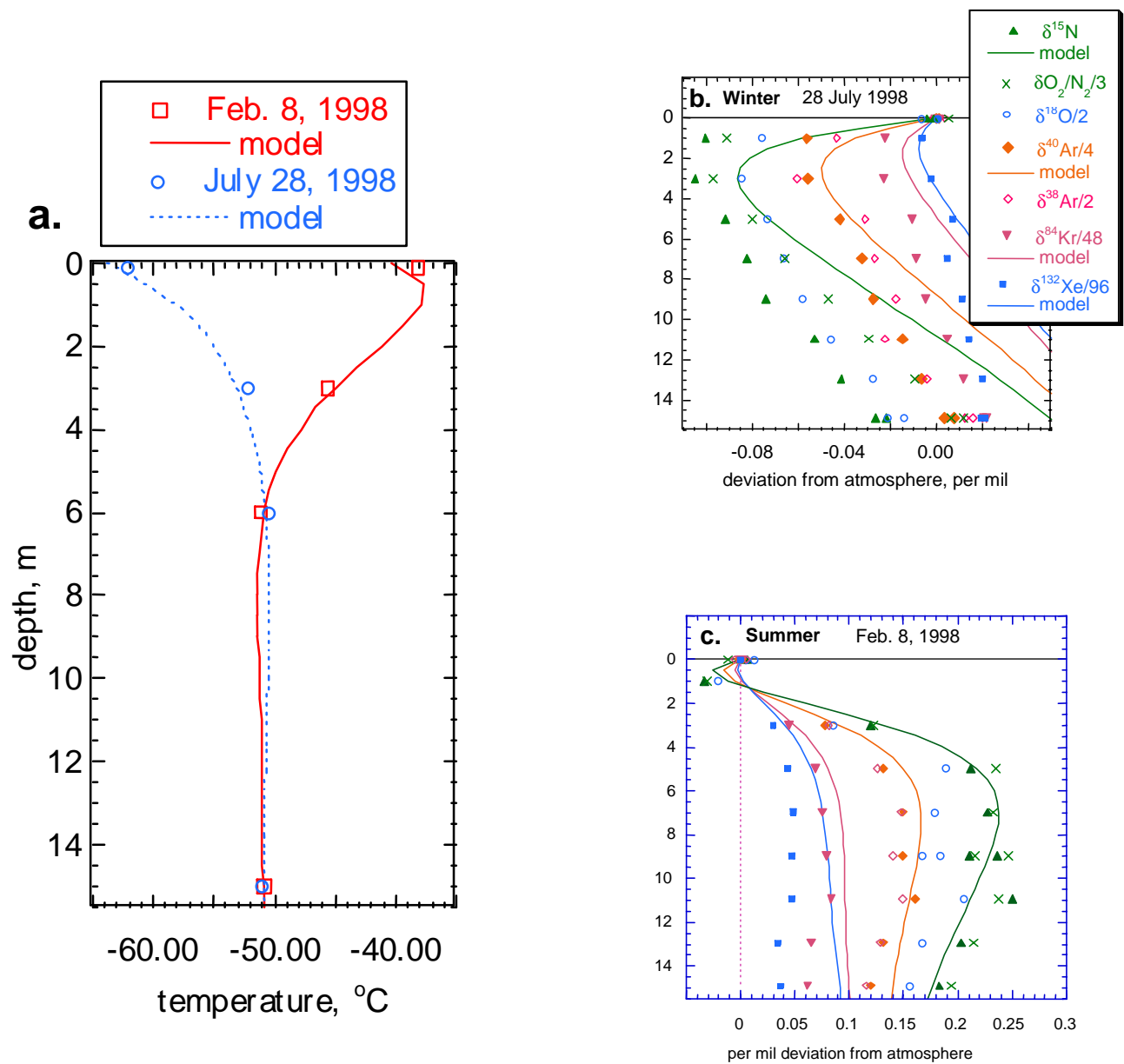


Fig. 1. Depth profiles of measured temperature (a) and gas composition in the shallow firn at the south pole. Winter sampling (b) reveals a large negative isotope excursion due to thermal isotope fractionation caused by the cold temperatures at the surface relative to the temperature at 10 m depth. Summer sampling (c) reveals the opposite pattern, as expected. Model calculations (shown as lines) are based on the measured surface temperature history for this site and the known diffusivities of these gases. Gravitational settling causes a steady linear downward increase that is independent of season and that is additive to the thermal fractionation signal. The strong winter anomaly, even stronger than predicted from the model, argues against any vigorous convective mixing that would have erased the signal by homogenizing it with the atmosphere.

the tubes were extended into the building for samples to be taken indoors. The hole was backfilled with snow and layers of slush to impede downward movement of surface air. After backfilling, the tubes were pumped on to remove as much contamination as possible. Sampling by pumping on the tubes was done on February 8 and again on July 28. This technique is similar to

that used to sample air in sand dunes [Severinghaus *et al.*, 1997]. Samples were pumped at 4 L min^{-1} for 16 minutes to flush the lines and the sample flasks. The air was dried by passing it through a column packed with granular phosphorous pentoxide. The air was stored in 2-L flow-through glass flasks as described by Battle *et al.* [1996]. Flasks were analyzed mass

spectrometrically at the University of Rhode Island and SIO for $^{15}\text{N}/^{14}\text{N}$ and $^{40}\text{Ar}/^{36}\text{Ar}$ ratios and other gas ratios including $^{40}\text{Ar}/^{38}\text{Ar}$, $^{18}\text{O}/^{16}\text{O}$ of O_2 , O_2/N_2 , $^{84}\text{Kr}/^{36}\text{Ar}$, and $^{132}\text{Xe}/^{36}\text{Ar}$. Results of two aliquots, measured for $\delta^{15}\text{N}$, $\delta^{18}\text{O}$, and $\delta\text{O}_2/\text{N}_2$ from 3 m depth at south pole February 8, were rejected on the basis of high CO_2 concentrations (22 and 31% above normal) that suggested gross contamination with laboratory air.

RESULTS

Gravitational settling enriches the heavier molecules at the bottom of a stagnant air column in proportion to the mass difference Δm , so gravity should affect $\delta^{40}\text{Ar}$ four times as much as $\delta^{15}\text{N}$ [Lindemann and Aston, 1919; Craig *et al.*, 1988; Sowers *et al.*, 1989; Schwander, 1989]. In other words, if gravity were the only process acting on a gas mixture, $\delta^{15}\text{N}$ would be equal to $\delta^{40}\text{Ar}/4$. To facilitate comparison all results are presented here with the δ value divided by Δm .

SIO wanted to know what the shallow firn gas profile did in winter when the temperature gradient was reversed from the more commonly studied summer profile. In winter cold dense air overlies warmer air in the firn and might become unstable and undergo Bernard convection [Powers *et al.*, 1985]. If the wintertime thermal diffusion anomaly were not expressed because of convection, the mean annual thermal fractionation signal would be nonzero even when annual mean temperature change was zero. In other words, there would be a rectifier of seasonal temperature change, and this signal would be transmitted to the deep firn and create a bias in the bubble air-based thermal diffusion paleoindicator.

The tendency to undergo Bernard convection in a porous media can be estimated from the dimensionless Rayleigh number [Powers *et al.*, 1985]. Taking the coefficient of thermal expansion of air as 0.0045 K^{-1} , the depth of the layer with strong temperature gradient as 5 m, the temperature difference across this layer as 10 K, the permeability of the firn as $4 \times 10^{-9}\text{ m}^2$, the dynamic viscosity of air as $1.4 \times 10^{-5}\text{ kg m}^{-1}\text{ s}^{-1}$, and the thermal diffusivity as $0.22\text{ m}^2\text{ s}^{-1}$, a Rayleigh number of 4 is obtained. This is sufficiently close to the critical value for the onset of convection (10-30) that natural convection in snow cannot be ruled out by this calculation. Further, this assumes a one-dimensional firn without irregularities (snow dunes or sastrugi) that might get convection started much more easily. If vigorous convective stirring during winter were important, one would expect to see a greatly attenuated thermal fractionation signal. In fact a fully developed negative isotope anomaly, with an amplitude slightly larger than that predicted by a pure diffusion model driven by observed surface air temperature (Figure 1b) was observed. The summer data are shown for comparison (Figure 1c).

The marked deviation of the winter data from the deeper part of the model is not understood. It seems possible that gradual downward flow is advecting the values obtained at shallower depths to deeper levels, as hypothesized at Siple Dome (J. Severinghaus, Thermal fractionation of air in polar firn by seasonal temperature gradients, in preparation, 2000). Alternatively, intermittent turbulent convective mixing may be

weak enough that it does not homogenize and erase the isotope fractionation, yet strong enough that it acts to transport the isotope signal downward into the firn faster than would molecular diffusion alone. In other words, mixing by eddy diffusion would double or triple the effective diffusivity over the molecular diffusivity. Adding an arbitrary intermittent eddy diffusion term to our model produces this effect, demonstrating plausibility but not proof.

CONCLUSIONS

Seasonal effects penetrate no deeper than 40 m in the firn at current measurement precision. Therefore, the firn-ice transition at 114 m is effectively isolated from seasonality and records a time-averaged value of the fluctuations in the shallow firn. Trapped air in ice is, therefore, susceptible to rectification in principle. The gas isotope profile observed in the shallow firn during winter is nearly symmetrical to that observed during summer, suggesting that they should cancel in the annual mean. However, this conclusion assumes that the sampling dates were representative of average conditions and did not miss strong episodes of convection by chance. Deep firn isotopes do fit a pure gravitational model well, though, indicating that seasonal anomalies probably do cancel [Battle *et al.*, 1996]. No sign of thermally driven adsorption or desorption was seen. These conclusions should be regarded as valid for one site at one time, and generalized to other sites with caution.

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