High Temperature Vaporization Behavior of Oxides. I. Alkali Metal Binary Oxides

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In order to assess the high temperature vaporization behavior and equilibrium gas phase compositions of binary alkali metal oxides, the relevant thermodynamic and molecular constant data have been compiled and critically evaluated. Selected values of the Gibbs energy and enthalpy functions of condensed and vapor phases are given in the form of equations valid over wide temperature ranges, along with the standard entropies and enthalpies of formation. These data were used to generate plots of the equilibrium partial-ressures of vapor species as functions of temperature for representative conditions ranging from reducing to oxidizing. Maximum vaporization rates have been calculated using the Hertz–Knudsen equation. Literature references are given.

Key words: alkali metal; critically reviewed data; enthalpy increment; enthalpy of formation; Gibbs energy function; high temperature; oxide; partial pressure; vaporization; vaporization rate.

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1. Introduction 1.1. Background

Oxide materials are used or encountered in a wide variety of high temperature applications where vaporization rates and thermodynamic stabilities are often limiting factors. The efficient design and operation of high temperature devices and processes requires reliable information about the stability and volatility of these oxides so that vaporization losses and component lifetimes can be predicted.

Despite continuing research efforts and the use of increasingly sophisticated techniques, there are still many gaps in our knowledge and understanding of the detailed vaporization thermodynamics of metal oxide systems. This puts a premium on critical review of the literature and selection of the necessary thermodynamic data. Even in cases where most of the requisite data are compiled, however, the user still must resort to a significant amount of additional calculations, sometimes unfamiliar, in order to evaluate vapor composition and vaporization behavior for specific environmental conditions. Thus we perceive a definite need for critically evaluated data, presented in a format that gives users ready access to the detailed vaporization chemistry and equilibrium partial pressures of the various species.

Reviews and compilations of data on oxide thermodynamics and vaporization behavior have been done by Brewer, 1 Coughlin, 2 Kelley, 3 Schick, 4 Ackermann and Thorn, 5 Brewer and Rosenblatt, 6,7 Kelley and King,8 Olette and Ancey-Moret, Wicks and Block, the Bureau of Standards, 11-16 the JANAF group, 17,21 Barin and Knacke, 22 Samsonov et al., 23,24 the IVTAN (High Temperature Institute, State Institute of Applied Chemistry, National Academy of Sciences of the U.S.S.R.) group, 25,26 and Pankratz.27 Pedley and Marshall²⁸ have submitted a review of thermodynamic data of gaseous monoxides to this journal. The later reviews are more useful, since much of the relevant data has appeared in the last few years. However, the recent critical reviews by JANAF and IVTAN cover only certain portions of the periodic table and do not include many oxides of interest. This work is the first part of a program sponsored by the Office of Standard Reference Data of the National Bureau of Standards to provide critically reviewed thermodynamic data of oxides, along with overviews of their high temperature data of oxides, along with overviews of their high temperature vaporization phenomena. Participants in the project are Leo Brewer of the University of California at Berkeley, Gerd M. Rosenblatt of the Los Alamos Scientific Laboratory, and the present authors.

1.2. Scope

The present study is intended to serve as a source of critically reviewed thermodynamic data, and to provide an overview of high temperature vaporization phenomena for condensed alkali oxides. Substances considered are condensed phases stable over a portion of the temperature range between 298 and 3000 K and vapor phase species that have been, or under the proper conditions might be, observed in equilibrium with the condensed phases. In the absence of accurate phase boundary data, the solid phases are treated as line compounds. Data are tabulated for liquid phases with compositions identical to the solid phases. The limited amount of thermodynamic data for homogeneous oxide phases of variable composition are cited in the text but were not evaluated. The properties of these phases, and of solutions of oxygen in liquid alkali metals, are beyond the scope of this work, and data necessary for their treatment are for the most part unavailable. The information presented consists of thermodynamic data for the various chemical species involved in the vaporization processes, equilibrium partial pressures of significant vapor species found over condensed oxide phases as functions of temperature and oxygen partial pressure, and maximum vaporization rates calculated from the Hertz-Knudsen equation of classical kinetic theory. Graphs of partial pressures of significant vapor species in equilibrium with condensed phases are provided as a concise reference to vaporization behavior under various conditions of temperature and oxygen potential.

1.3. Literature Reviewed

A thorough search was made in Chemical Abstracts and in the Bulletin of Chemical Thermodynamics and Thermochemistry for entries related to the substances under consideration. The literature review encompassed publications up to about December 1982. When an adequate review of earlier work on a subject is available, reference is given to the review only, except when specific references are desirable.

2. Thermodynamic Properties

2.1. Discussion

a. Evaluation of Data

The methods of data evaluation used in the present work are similar to those used for the JANAF tables¹⁷ and need not be repeated in detail. In addition to consistency of data within a given source and general assessments of experimental methods, criteria for data evaluation included the degree of agreement of second- and third-law reaction enthalpies evaluated from equilibrium data. For gaseous alkali metal monoxides, the Rittner ionic model^{29,30} was used as an aid in data evaluation. Data from original sources were recalculated, when necessary, using selected values of auxiliary data. CODATA³¹ values were used when available. Selected thermodynamic data derived from experimental information are consistent with the original observations within experimental uncertainty.

The mathematical procedures used to process thermodynamic data have been discussed by JANAF, ¹⁷ among others. Data for substances undergoing transitions between solid phase forms or between solid and liquid phases were treated by using data for the lower temperature form below the normal transition temperature, taking the enthalpy and entropy of transition into account at the transition temperature, and using data for the higher temperature form above the transition temperature. The Gibbs energy functions of condensed phases calculated by this procedure are consistent with the enthalpies of formation of the condensed phase stable at 298.15 K and 1.013 25 bar.

b. Symbols

The thermodynamic symbols used in the present work are $\Delta_{\rm f} H_{298}^{\circ}/R$ The standard molar enthalpy of formation at 298.15 K divided by the molar gas constant $(H_{298}^{\circ} - H_{0}^{\circ})/R$ The molar enthalpy increment between 0 and 298.15 K divided by the molar gas constant $-(G_T^{\circ}-H_{298}^{\circ})/RT$ The molar Gibbs energy function divided by the molar gas constant The molar enthalpy increment $(H_T^{\circ} - H_{298}^{\circ})/R$ between 298.15 and T K divided by the molar gas constant S_T°/R The standard molar entropy at TK divided by the molar gas constant

stant

Solid-phase transition tempera-

ture and molar enthalpy of transi-

tion divided by the molar gas con-

 $T_{trs}, \Delta_{trs}H^{\circ}/R$

 $T_{
m fus}$, $\Delta_{
m fus}H^{\circ}/R$

Melting temperature and molar enthalpy of fusion divided by the molar gas constant.

c. Units

Following the suggestion of Pitzer and Brewer, 32 thermodynamic quantities are expressed in dimensionless units, e.g., S°/R , G°/RT , or in Kelvin units, e.g., $\Delta H^{\circ}/R$, $\Delta G^{\circ}/R$. The symbols K and kK, where 1 kK = 10^3 K, are used to represent Kelvin units. Values of R used in the present work are 1.987 19 cal/mol K and 8.3144 J/mol K. The temperature scale is nominally the 1968 International Practical Temperature Scale (IPTS) for measured quantities and the thermodynamic temperature scale for estimated quantities and calculated gaseous quantities. If conversion to the 1948 IPTS is desired, subtract 1 K between 1015 and 1380 K and 2 K between 1381 and 1926 K, for an accuracy within 1 K. The atomic weights are from the 1975 Report of the Commission on Atomic Weights, Inorganic Division of IUPAC. 33

Molecular constants cited in the text are internuclear distances r in cm, vibrational fundamentals ω in cm⁻¹, the molecular symmetry number σ , electronic state quantum weights g_i , and excited electronic state term values T_e in cm⁻¹. A more thorough discussion of these quantities and their use in calculating thermodynamic functions has been given in JANAF.¹⁷

Table 1. Alkali metal oxides melting above 300 K, their melting temperatures, and enthalpies of fusion.

	T _m (K)	$\Delta_{fus}^{H/R}$ (kK)
Li ₂ 0	1711 <u>+</u> 5	6.5 <u>+</u> 1.0
Li ₂ 0 ₂	1250 <u>+</u> 250 ^b	5.0 ± 2.0 ^b
ao ₂	825 <u>+</u> 10	3.3 ± 1.0
Na ₂ 0	1405 <u>+</u> 100°	5.7 <u>+</u> 0.5°
Na ₂ 0 ₂	948 <u>+</u> 150 ^c	3.9 ± 0.7
ко ₂	865 <u>+</u> 80°	1.6 <u>+</u> 0.7°
ς ₂ 0 ·	1190 <u>+</u> 50	4.8 <u>+</u> 1.0
202	763 <u>+</u> 100°	3.1 <u>+</u> 1.0
1b0 ₂	685 <u>+</u> 100°	2.7 + 0.6
tb ₂ 0	900 <u>+</u> 50	3.7 <u>+</u> 1.0
h ₂ n ₂	843 <u>+</u> 100 ^c	3.4 ± 0.7
^{2s0} 2	723 ± 50°	3.0 <u>+</u> 0.5
3 ⁰⁸ 3	343 <u>+</u> 70 ^d	
Cs ₂ 0	768 <u>+</u> 5	3.1. <u>+</u> 0.5
.s ₂ 0 ₂	863 <u>+</u> 100°	3.5 <u>+</u> 0.5
Cs ₃ 0	437 <u>+</u> 4 ^e	

a See table 5 for references.

b Estimated

C Uncertainty estimated (this work).

d Peritectic decomposition to CsO₂ and liquid.

e Peritectic decomposition to Cs20 and liquid.

Table 2. Gibbs energy functions of substances encountered in the high temperature vaporization of alkali metal oxides.

Table 2. Gibbs energy functions of substances encountered in the high temperature vaporization of alkali metal oxides. -- Continued

		-(G ^o T-H	° ₂₉₈)/rt	- A + BT + C	T ² + DT ³			-(G°T-H°	298 ^{)/RT} =	A + BT + CT	2 + DL3
	A	10 ³ B	10 ⁶ 0	10 ⁹ D	T Range (K)		A	10 ³ B	10 ⁶ c	10 ⁹ D	T Range (K
(g)	19.92 18.23 18.54	-4.808 3.200 2.639	11.711 -0.891 -0.551	-6.412 0.122 0.053	298 - 800 800 -1600 1600 -3000	Na ₂ 0(в)	12.07 5.85 5.74	-24.530 8.562 8.638	58•798 -0•633 -0•764	-36.029 -0.103 0.000	298 - 600 600 -1300 1300 -1405
₂ (g)	25.46 23.11	-6.723 4.240	16.189 -0.889	-8.707 0.088	298 - 800 800 -1600	Na ₂ 0(1)	7.67.	6.193	0.000	0.000	1405 -3000
. (-)	23.35	3.886 26.852	-0.726 -28.145	0.066	1600 -3000 298 - 600	Na ₂ 0(g)	33.81 28.82	-19.835 6.843	47.718 -0.428 -1.673	-29.514 -0.304 0.168	298 - 600 600 -1200 1200 -3000
₃ (ε)	18.68 21.20 23.86	14.143	-6.413 -1.892	1.381 0.214	600 -1400 -1400 -= 3000	Na ₂ 0 ₂ (s)	28.65 14.11 7.25	7.801 -22.460 10.565	53.337 0.000	-28.544 0.000	298 - 785 785 - 948
i(s)	3,98	-3.544	6.775	0.000	298 -453	Na ₂ O ₂ (1)	9.33	8.368	0.000	0.000	948 -3000
i(1)	1.42 3.33	6.096 2.796	-2.229 -0.300	0.386 0.000	453 -1500 1500 -3000	Na ₂ 0 ₂ (g)	38.25 31.80	-25.730 8.719	61.721 -0.327	-38.044 -0.461	298 - 600 600 -1200
1(8)	17.23 15.64	-4.681 2.942	11.351 -0.713	-6.194 0.080	298 - 800 800 - 3000		31.47	10.223	-2.170	0.216	1200 -3000
10(g)	26.79	-11.266	27.053	-16.660	298 - 600	K(s)	8.55	-3.293	-3.967	21.487	298 - 336
10(8)	23.96 23.86	3.847 4.403	-0.170 -0.857	-0.163 0.079	600 -1400 1400 -3000	K(1)	8.07 5.88 6.74	-7.660 6.301 4.407	27.588 -2.422 -1.018	-21.265 0.463 0.113	336 - 500 500 -1300 1300 -3000
10 ₂ (g)	31.29 27.24 27.08	-15.850 5.515 6.366	37.839 -0.227 -1.227	-23.079 -0.225 0.115	298 - 600 600 -1400 1400 -3000	K(g)	20.16	-7.220	17.477	-10.913	298 - 600
i ₂ (g)	25.25 22.02	-12.740 4.591	50.728 -0.578	-19.081 -0.071	298 - 600 600 -1400	ro()	18.32	2.665 2.624	-0.428 -0.553	-0.016 0.054	600 -1400 1400 -3000
	22.14	4.671	-0.910	0.081	1400 -3000	KO(g)	30.30 26.77 27.11	-13.720 5.249 4.846	33.313 -1.084 -1.027	-20.900 0.057 0.100	298 - 600 600 -1400 400 -3000
i ₂ 0(s)	7.22 2.00 6.98	-21.201 6.391 -2.441	50.163 0.899 5.633	-30.238 -0.631 -1.311	298 - 600 600 -1300 1300 -1711	K02(8)	18.26 11.55	-28.382 8.027	67.924 1.568	-41.676 -1.154	298 - 600 600 - 865
1 ₂ 0(1)	-8.27	18.200	-4.168	0.410	1711 -3000	KO ₂ (1)	7.69	16.211	-4.146	0.478	865 -3000
₂ 0(g)	29.78 25.18 25.11	-17.904 6.391 7.070	42.895 -0.501 -1.399	-26.299 -0.189 0.131	298 - 600 600 -1400 1400 -3000	KO ₂ (g)	33.98 29.83 29.40	-17.218 5.207 6.937	41.356 0.578 -1.453	-25-457 -0-598 0-144	298 - 600 600 -1200 1200 -3000
1202(a)	10.18 3.56	-26.989 7.970	63.832 1.514	-38.277 -0.867	298 - 600 600 -1250						
i ₂ 0 ₂ (1)	-6.65	23.805	-6.842	0.910	1250 -2000	K ₂ (g)	31.67 28.28 28.29	-13.248 4.858 5.247	32.007 -0.654 -1.247	-19.914 -0.095 0.127	298 - 600 600 -1400 1400 -3000
i ₂ 0 ₂ (g)	33.79 28.15 27.76	-22.112 7.556 9.099	52.593 -0.101 -1.769	-31.942 -0.380 0.164	298 - 600 600 -1400 1400 -3000	K ₂ 0(s)	14.68 8.73	-27.005 5.772	64.596 4.132	-39.642 -2.363	298 - 600 600 - 890
i ₃ (g)	35.23	-19.938	48.156	-30.024	298 - 600		2.42	26.192	-17.967	5.636	890 -1190
	30.14 30.43	7.247	-1.212 -1.506	-0.022 0.145	600 -1400 1400 -3000	K ₂ 0(1)	0.80	18.889	-4.532	0.467	1190 -3000
i ₃ 0(g)	35.73 29.44 29.42	-24.797 8.625 9.505	59.413 -0.554 -1.852	-36.522 -0.314 0.173	298 - 600 600 -1400 1400 -3000	K ₂ 0(g)	36.52 31.63 31.83	-19.066 7.039 7.112	45.948 -1.027 -1.465	-28.502 -0.062 0.141	298 - 600 600 -1400 1400 -3000
ı(s)	7.27	-7.591	12.955	0.000	298 - 371	K ₂ 0 ₂ (s)	16.08	-24.847	59-193	-31.877	298 - 763
·(1)	4.78	3.249	2.713	-2.360	371 - 700	K ₂ 0 ₂ (1)	1.90	25.223	-7.055	0.859	763 -3000
	4.10 5.29	6.542 4.045	-2.625 -0.853	0.509 0.083	700 -1400 1400 -3000	K ₂ 0 ₂ (g)	41.01 34.06 34.32	-27.084 9.978 10.122	65.212 -1.428 -2.079	-40.422 -0.095 0.199	298 - 600 600 -1400 1400 -3000
a(g)	19-03 17-44	-4.681 2.935	11.351 -0.708	-6.194 0.079	298 - 800 800 -3000	Rb(s)	9.15	0.278	0.000	0.000	298 - 313
0(g)	29.00 25.95 25.75	-12.251 4.104	29.466 -0.075 -1.060	-18.185 -0.252 0.106	298 - 600 600 -1200 1200 -3000	Kp(1)	8.93 7.28	-2.497	14.303	-10.022	313 - 600
0 ₂ (s)		4.932 -16.637	39.727	-20.900	298 - 825		8.49	6.671 4.095	-2.722 -0.872	0.536 0.087	600 -1400 1400 -3000
02(1)	4.82	19.292	-5.267	0.623	825 -3000	Rb(g)	21.33 19.48 19.61	-7.220 2.721 2.617	17.477 -0.492 -0.551	-10.913 0.008 0.054	298 - 600 600 -1400 1400 -3000
0 ₂ (g)	32.89 28.65 28.33	-17.038 5.644 6.905	40.869 0.013 -1.442	-25.116 -0.377 0.142	298 - 600 600 -1200 1200 -3000	RbO(g)	31.31 27.44	-15.061 5.774	36.458 -1.306	-22.913 0.092	298 - 600 600 -1400
₂ (g)	25.96	-13.120 4.770	31.694 -0.653	-19.719 -0.056	298 - 600 600 -1400	RbO ₂ (s)	27.88 18.55	5.169 -23.671	-1.127 56.381	0.112 -32.530	1400 -3 000 298 - 685
	25.93	5.111	-1.101	0.103	1400 -3000	RbO ₂ (1)	6.64	20.954	-6.185	0.777	685 -3000

Table 2. Gibbs energy functions of substances encountered in the high temperature vaporization of alkali metal oxides.--Continued

 $-(G^{O}_{T}-H^{O}_{298})/RT = A + BT + CT^{2} + DT^{3}$ 10⁶C 10⁹D T Range (K) 10³B A 41.958 -0.689 -1.336 -25.875 -0.122 0.126 298 - 600 600 -1400 1400 -3000 35.49 31.01 31.05 -17.470 RbO2(g) 6.308 6.699 298 - 600 31.551 -0.870 -0.984 -19.709 0.008 0.096 $Rb_2(g)$ 34.09 13.042 600 -1400 1400 -3000 30.75 30.98 4.882 298 - 675 675 - 750 750 - 810 810 - 900 18.49 11.84 13.01 60.654 -35.906 Rb₂O(s) 9.461 -0.805 1.174 0.000 22.56 -27.259 40.596 -15.312 Rb20(1) 0.744 900 -3000 4.60 21.797 -6.188 298 - 600 600 -1400 1400 -3000 38.69 33.76 33.98 46.426 -1.065 -1.466 Rb20(g) -19.261 -28.837 7.098 0.052 298 - 843 Rb₂O₂(s) 18.54 -22.449 53.321 -27.871 843 -3000 4.69 24.083 -6.542 0.782 Rb₂0₂(1) 43.19 36.17 36.49 -27.341 10.088 10.138 65.845 -1.495 -2.081 298 - 600 -40-844 Rb202(g) -0.079 0.199 600 -1400 1400 -3000 cs(s) 10.20 0.143 0.000 0.000 290 - 302 301 - 600 600 -1400 1400 -3000 -9.509 0.371 10.00 -2.498 6.048 13.761 Ca(1) 8.47 9.34 -0.938 4.284 0.095 22.00 20.14 20.26 -7.220 2.721 2.644 298 - 600 600 -1400 1400 -3000 Cs(g) -0.492 0.008 -0.566 0.057 298 - 600 600 -1400 1400 -3000 31.75 27.99 -21.217 CsO(g) -14.342 34.392 5.474 5.389 -0.899 -1.175 -0.034 0.116 298 - 723 Cs02(s) 19.16 -19.641 47.344 -26.434 -5.268 0.615 723 -3000 19.410 Cs0₂(1) 8.20 36.29 31.81 31.82 -26.044 298 - 600 Cs0₂(g) -17-557 42.187 6.285 6.735 -0.658 -1.353 1400 -3000 35.14 32.38 32.50 298 - 600 -12-719 Cs₂(g) -8.705 21,922 -0.822 -0.905 600 -1400 1400 -3000 298 - 768 -18.388 44.069 -23.923 Cs20(s) 19.86 Cs₂O(1) 768 -3000 11.10 -1.953 0.000 14.255 40.61 35.68 35.89 46.327 -1.063 -1.427 298 - 600 Cs₂O(g) -19.222 -28.754 7.100 -0.053 0.142 600 -1400 1400 -3000 Cs₂O₂(s) 19.88 -21.301 50.979 -26.536 298 - 863 0.748 863 -3000 6.32 23.533 -6.310 00202(1) -40.513 -0.089 0.200 44.36 37.38 37.67 -27.134 10.047 10.151 65.374 -1.461 -2.088 298 - 600 Cs₂O₂(g) 600 **-1**400 1400 **-3**000

Table 3. Enthalpy increment above 298.15 K of substances encountered in the high temperature vaporization of alkali metal oxides. (Units = kK)

		(H° _T -H° ₂	₉₈)/R = a	+ bT +cT ²	+ aT ³
	a	10 ³ b	10 ⁶ c	10 ⁹ d	T Range (K)
0(5)	-0.81	2.849	-0.460	0.220	298 - 800
	-0.74	2.550	- 0.022	0.004	800 -3 000
0 ₂ (g)	-1.00	3.201	0.517	0.027	298 - 800
	-1.21	3.665	0.301	-0.026	800 - 3000
0 ₃ (g)	-1.11	2.540	4.423	-1.655	298 - 800
	-1.79	5.136	1.061	-0.178	800 -1400
	-3.41	8.872	-1.787	0.540	1400 -3000
Li(s)	-0.71	1.825	1.880	0.000	298 - 454
Li(1)	-0.83	3.893	-0.382	0.109	454 -1500
	-0.66	3.452	0.000	0.000	1500 -3000
Li(g)	-0.75	2.505	-0.011	0.007	298 -1 000
	-0.79	2.598	-0.069	0.016	1000 -3 000
LiO(g)	-1.03	3.102	1.356	-0.421	298 - 1000
	-1.49	4.346	0.176	-0.029	1000 - 3000
Li0 ₂ (g)	-1.32	3.523	3.448	-1.277	298 - 1000
	-2.03	6.054	0.382	-0.036	1000 - 3000
Li ₂ (g)	-1.26	4.064	0.573	-0.173	298 -1000
	-1.05	3.626	0.811	-0.180	1000 -2000
	-2.84	6.172	-0.395	0.009	2000 -3000
Li ₂ 0(a)_	-1.50 -2.20 43.02 -68.98	3.113 6.271 -63.456 122.314	2.475 28.986 -64.562	-2.706 -0.285 0.000 12.854	298 = 800 800 -1300 1300 -1400 1400 -1711
Li ₂ 0(1)	1.56	12.077	0.000	0.000	1711 -3000
Li ₂ 0(g)	-1.59	4.555	2.933	-1.090	298 -1000
	-2.32	6.959	0.195	-0.025	1000 -3000
Li ₂ 0 ₂ (s)	-1.85	3.431	10.570	-4.095	298 - 600
	-2.59	7.212	4.044	-0.299	600 - 1250
Li ₂ 0 ₂ (1)	-2.97	16.103	0.000	0.000	1250 -2000
Li ₂ 0 ₂ (g)	-1.72	4.185	6.036	-2.307	298 -1000
	-3.17	9.068	0.336	-0.043	1000 -3000
Li ₃ (g)	-2.00	6.576	0.497	-0.205	298 - 1000
	-2.10	6.951	0.020	-0.003	1000 - 3000
Li ₃ 0(g)	-2.20	6.372	3.833	-1.467	298 - 1000
	-2.73	8.650	0.712	-0.092	1000 - 3000
Na(s)	-0.73	1.534	3.025	0.000	298 - 371
Na(1)	-1.00	4.704	-1.519	0.648	371 - 700
	-0.87	4.189	-0.796	0.295	700 -1400
	-0.80	3.417	0.167	-0.024	1400 -3000
Na(g)	-0.75	2.505	-0.011	0.007	298 -1 000
	-0.77	2.548	-0.033	0.007	1000 -3 000
NaO(g)	-1.15	3.502	1.334	-0.501	298 -1000
	-1.57	4.806	-0.050	0.002	1000 -3000
NaO ₂ (s)	-2.37	7.200	2.470	-0.005	298 - 825
Na0 ₂ (1)	-1 -86	12.581	0.000	0.000	825 -3000
NaO ₂ (g)	-1.52	4.341	2.849	-1.184	298 - 800
	-2.07	6.368	0.255	-0.036	800 -3000
Na ₂ (g)	-1.33	4.373	0.260	-0.021	298 -1000
	-1.49	4.535	0.440	-0.196	1000 -2000
	-3.20	7.463	-1.237	0.124	2000 -3000

Table 3. Enthalpy increment above 298.15 K of substances encountered in the high temperature vaporization of alkali metal oxides. (Units = kK) --Continued

Table 3. Enthalpy increment above 298.15 K of substances encountered in the high temperature vaporization of alkali metal oxides. (Units = kK) --Continued

		(H ^o T-H ^o 2	₂₉₈)/R = 6	+ bT +cT ² +	ar ³			(HoT-HoS	98)/R = a	+ bT +cT ² +	dT ³
	a	10 ³ ъ	10 ⁶ c	10 ⁹ a	T Range (K)		a	10 ³ b	10 ⁶ c	10 ⁹ d	T Range (K)
Na ₂ 0(s)	-2.28 -16.15	6.973 50.726	2.245	-0.008 16.019	298 -1000 1000 -1300	RbO ₂ (s)	-2.38	6.225	6.930	-3.450	298 - 685
	-9.12	15.618	-43.680 0.000	0.000	1300 -1405	Rb0 ₂ (1)	-1.87	12.581	0.000	0.000	685 -3000
Na ₂ 0(1)	0.89	12.581	0.000	0.000	1405 -3000	вю ₂ (g)	-1.62 -2.13	4.900 6.542	1.984 0.165	-0.713 -0.021	298 - 1000 1000 - 3000
1a ₂ 0(x)	-1.87 -2.30	5.774 7.298	1.934 0.072	-0.776 -0.009	298 -1 000 1000 -3 000	Rb ₂ (g)	-1.34	4.495	0.002	0.000	298 -1000
Na ₂ 0 ₂ (s)	-2.70 -3.91	7.093 13.661	7.379	-2.661 0.000	298 - 785 785 - 948	Rb ₂ 0(s)	-1.38 -2.29	4.571 5.844	-0.050 7.205	0.011 -3.662	1000 -3 000 298 - 675
Na ₂ 0 ₂ (1)	-2.90	16.774	0.000	0.000	946 -3000	20(0)	-1.70 42.92	5.629 -111.056	3•757 80•01 <i>3</i>	0.000	675 - 750 750 - 810
Na ₂ O ₂ (g)	-2.27	6.470	4.420	-2.026	298 - 800		-21.15	47.127	-17.613	0.000	810 - 900
	-3.02	9.508	0.202	-0.029	800 -3000	Rb ₂ 0(1)	-0.65	12.581	0.000	0.000	900 -3000
K(s)	-5.31	44.899	-134.115	145.167	298 - 336	Rb ₂ 0(g)	-1.87 -2.12	5.963 6.885	1.190 0.042	-0.484 -0.006	298 - 1000 1000 - 3000
K(1)	-0.97 -0.44	4.470 2.917	-1.145 0.366	0.491	336 -1000 1000 -3000	Rb ₂ O ₂ (s)	-2.84	7.387	8.193	-3.258	298 - 843
K(ε)	-0.75	2.505	-0.011	0-007	298 =1000	Rb ₂ 0 ₂ (1)	-3.44	16.774	-0.001	0.000	843 -3000
	-0.83	2.694	-0.139	0.032	1000 -3000	Rb ₂ O ₂ (g)	-2.63 -3.04	8.322 9.827	1.958 0.037	-0.805 -0.008	298 -1000 1000 -3000
KO(g)	-1.49 -1.32	5•177 4•586	-0.737 -0.033	0.287 0.004	298 - 1000 1000 -3 000	- ()	4 45	~ ~~			000 #00
102(a)	-2.45	6.796	5.390	-2.202	298 - 865	Cs(s) Cs(1)	-1.15 -0.95	3.860 4.187	0.000 -0.652	0.000	298 - 302 302 -1000
KO ₂ (1)	-3.22	12.581	0.000	0.000	865 -3000	(3(1)	-0.83	3.719	-0.026	0.019	1000 -3000
KO ₂ (g)	-1.59 -2.14	4.760 6.537	2.150 0.166	-0.783 -0.021	298 - 1000 1000 - 3000	Cs(g)	-0.75 -0.78 -1.58	2.505 2.597 3.809	-0.011 -0.087 -0.708	0.007 0.026 0.133	298 -1000 1000 -2000 2000 -3000
K ₂ (g)	-1.31 -2.27	4.189 6.708	0.807 -1.357	-0.430 0.173	298 -1000 1000 -3000	CsO(g)	-1.13 -1.78	2.627 5.832	4.790 -0.588	-2.958 0.094	298 - 600 600 -1400
K ₂ 0(s)	-2.29 -145.88 -25.75	6.249 460.411 45.078	5.508 -473.284 -12.511	-2.362 165.926 0.000	298 - 890 890 -1070 1070 -1190	Cs0 ₂ (s)	-1.45 -2.78	5.236 9.206	-0.235 0.187	0.027 0.687	1400 -3 000 298 - 723
K ₂ 0(1)	0.02	12.581	0.000	0.000	1190 -3000	CsO ₂ (1)	-1.89	12.581	0.000	0.000	723 -3000
к ₂ о(е)	-1.85 -2.13	5.865 6.876	1.296 0.045	-0.525 -0.006	298 -1000 1000 -3 000	Cs0 ₂ (g)	-1.60 -1.82 -2.21	4•739 5•758 6•668	2.430 0.820 0.106	-1.066 -0.202 -0.012	298 - 600 600 -1400 1400 -3000
K ₂ 0 ₂ (s)	-3.35	10.417	2.473	1.023	298 - 763	Cs ₂ (g)	-1.35	4.447	0.257	-0.071	298 - 600
x ₂ 0 ₂ (1)	-3.19	16.772	0.000	0.000	763 -3000	2(6)	-1.35 -1.50	4.497 4.797	0.149 -0.067	0.001 0.054	600 -1400 1400 -3000
K ₂ 0 ₂ (g)	-2.50 -2.88	7•434 9•360	3.773 0.444	-2.062 -0.113	298 - 600 600 -1400	Cs ₂ 0(s)	-2.54	7.925	2.036	~0.010	298 - 768
	-3.09	9.859	0.046	-0.006	1400 -3000	Cs ₂ 0(1)	-1.80	12.581	0.000	0.000	768 -3000
Rb(s)	-1.14	3.820	0.000	0.000	298 - 313	Cs ₂ 0(g)	-1.83 -2.02	5.673 6.611	1.889	-1.008 -0.073	298 - 600 600 - 1400
Rb(1)	-1.11	5.039	-1.745	0.648	313 - 800		-2.14	6.907	0.032	-0.004	1400 -3000
	-1.02 -0.72	4•744 3•387	-1.403 0.132	0.514 0.001	800 -1400 1400 -3000	Cs ₂ O ₂ (s)	-3.19	9.820	3.178	-0.354	298 - 863
Rb(g)	-0.75 -0.84	2.505 2.713	-0.011 -0.154	0.007 0.036	298 -1000 1000 -3000	Cs ₂ O ₂ (1)	-3.60	16.771	0.001	-0.000	863 -3000
Rb0(g)	-1.56	5.330	-0.318	-0.026	298 -1000	Cs ₂ O ₂ (g)	-2.55 -2.88	7.747 9.403	3.261 0.422	-1.755 -0.108	298 - 600 600 -1400
	-1.33	4.869	-0.136	0.018	1000 -3000		-3.07	9.863	0.047	~0.006	1400 -3000

Table 4. Values of $^{\Delta}_{1}\text{H}^{o}_{298}/\text{R}$, $\text{S}^{o}_{298}/\text{R}$, and $(\text{H}^{o}_{298}\text{-H}^{o}_{0})/\text{R}$ of substances encountered in the high temperature vaporization of alkali metal oxides.

Table 4. Values of $\Delta_f H^0_{298}/R$, S^0_{298}/R , and $(H^0_{298}-H^0_0)/R$ of substem concurred in the high temperature vaporization of alkali metal on the second substantial second seco

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	. 10 /5	a0 /p	(110 140)/D			· · · · · · · · · · · · · · · · · · ·	Cont
	^1 ^{H°} 298/R (kK)	s ^o ₂₉₈ /R	(H ^O 298 ⁻ H ^O O)/R (kK)		4fH ⁰ 298/R	s°298/R	(H°298-H°0)/R
g)	29.97 ± 0.01	19.36 ± 0.00	0.81		(kK)		(kK)
(g)	0.00 <u>+</u> 0.00	24.66 <u>+</u> 0.01	1.04	K2(8)	15.30 <u>+</u> 0.50	30.03 <u>+</u> 0.01	1.29
(g)	17.06 <u>+</u> 0.24	28.73 + 0.01	1.25	K ₂ 0(s)	-43.58 <u>+</u> 0.25	11.32 + 0.75	
(s)	0.00 + 0.00	3.50 <u>+</u> 0.02	0.56	K ² 0(8)	-7.05 <u>+</u> 0.26	34.17 <u>+</u> 0.75	1.64
(g)	19.17 ± 0.05	16.68 ± 0.00	0.75	K ₂ O ₂ (s)	-59.43 <u>+</u> 0.50	13.09 <u>+</u> 0.75	
.0(g)	8.30 <u>+</u> 0.50	25.39 + 0.30	1.06	K ₂ 0 ₂ (g)	-20.00 <u>+</u> 1.00	37.66 + 1.00	2.10
0 ₂ (g)	9.30 <u>+</u> 3.50 ^a	29.32 <u>+</u> 0.75	1.31	()			
₂ (g)	25.85 ± 0.15	23.68 <u>+</u> 0.01	1.16	Rb(s)	0.00 <u>+</u> 0.00	9.23 <u>+</u> 0.04	0.90
₂ 0(s)	-71.91 ± 0.04	4.56 <u>+</u> 0.02	0.87	Rb(g)	9.73 <u>+</u> 0.03	20.44 <u>+</u> 0.00	0.75
. ₂ 0(g)	-20.78 ± 0.45	27.56 ± 0.75	1.50	RbO(g)	7.10 ± 0.50	29.45 <u>+</u> 0.30	1.25
. ₂ 0 ₂ (s)	-76.30 <u>+</u> 1.00	6.79 ± 0.50		RbO ₂ (s)	-33.57 <u>+</u> 0.30	15.65 <u>+</u> 0.08	2.02
1 ₂ 0 ₂ (g)	-35.00 ± 3.00	31.03 ± 1.00	1.63	RbO ₂ (g)	4.50 <u>+</u> 2.50 ^a	33.32 <u>+</u> 0.75	1.49
3(g)	38.25 <u>+</u> 4.00	32.77 ± 1.00	1.75	Rb ₂ 0(s)	-40.76 <u>+</u> 1.56	15.32 <u>+</u> 0.75	0.00
-3(6) -30(g)	-28.14 <u>+</u> 5.00	32.65 ± 1.00	1.88	Rb ₂ (g)	13.89 ± 0.30	32.48 <u>+</u> 0.30	1.30
30(6)	-20114 _ 5100)210) <u>7</u> 1100		Rb ₂ O(g)	-8.05 ± 2.50	36.31 <u>+</u> 0.75	1.66
ı(s)	0.00 <u>+</u> 0.00	6.17 + 0.02	0.78	Rb ₂ O ₂ (s)	-57.57 <u>+</u> 2.40	15.85 <u>+</u> 0.75	
a(g)	12.88 <u>+</u> 0.05	18.47 <u>+</u> 0.00	0.75	Rb ₂ O ₂ (g)	-19.10 <u>+</u> 4.00	39.81 <u>+</u> 1.00	2.14
10(g)	12.58 <u>+</u> 0.50	27.48 <u>+</u> 0.30	1.12	202(g/	-19:10 - 4:00)3.07 <u>-</u> 1.00	2.14
10 ₂ (s)	-31.35 <u>+</u> 0.35	13.94 <u>+</u> 0.15	2.21	Cs(s)	0.00 <u>+</u> 0.00	10.25 <u>+</u> 0.05	0.93
10 ₂ (g)	-1.01 ± 0.25ª	30.78 ± 0.75	1.42-	-Cs(-g)-	9.22 ± 0.25	-21-11 <u>+</u> 0-00	0.75
a ₂ (g)	17.17 ± 0.15	27.68 ± 0.01	1.25	Cs0(g)	4.90 <u>+</u> 0.50	29.96 <u>+</u> 0.30	1.19
1 ₂ 0(s)	-50.17 ± 0.50	9.03 <u>+</u> 0.02	1.49	CsO ₂ (s)	-34.42 <u>+</u> 0.25	16.81 <u>+</u> 0.75	
2 ₂ 0(g)	-3.82 <u>+</u> 1.00	31.35 ± 0.75	1.64	Cs0 ₂ (g)	1.00 ± 2.50ª	34.11 <u>+</u> 0.75	1.51
1 ₂ 0 ₂ (s)	-61.65 <u>+</u> 0.60	11.40 ± 0.15	1.89	Cs ₂ (g)	13.74 ± 0.51	34.15 <u>+</u> 0.10	1.33
1 ₂ 0 ₂ (g)	-10.00 <u>+</u> 4.00	35.05 <u>+</u> 1.00	1.91	02 ₂ 0(2)	-41.61 <u>+</u> 0.14	17.66 <u>+</u> 0.05	2.13
	_			Cs ₂ 0(g)	-19.10 <u>+</u> 3.00	38.23 <u>+</u> 1.00	1.66
s)	0.00 ± 0.00	7.78 ± 0.02	0.85	Cs ₂ O ₂ (s)	-59.88 <u>+</u> 1.26	17.36 <u>+</u> 0.75	
g)	10.73 + 0.03	19.27 <u>+</u> 0.00	0.75	00202(6)	-30.20 <u>·</u> 3.00	41.01 <u>+</u> 1.00	2.14
(g)	7.20 ± 0.50	28.61 <u>+</u> 0.30	1.28				
) ₂ (s)	-34.22 <u>+</u> 0.25	14.73 <u>+</u> 0.50	2.00				
(g)	0.60 ± 2.50ª	31.84 ± 0.75	1.46	a Estimate	d lower limit (see to	ext).	

Table 5. Sources of data in tables 1 through 4.

Table 5. Sources of data in tables 1 through 4 .-- Continued

	-H ^o ₂₉₈)/T, -H ^o ₂₉₈)	△ f ^{H°} 298	s° ₂₉₈	(H° ₂₉₈ -H° ₀)	Tfus, Afus ^H		(g° _T -H° ₂₉₈)/T, (h° _T -H° ₂₉₈)	≜ f ^H ⁰ 298	s° ₂₉₈	(H°298-H°O)	^T fus, [≜] fus ^H
)(g)	17	17 ⁸ ,31	17 ^a ,31	17ª,31	······································	к ₂ 0(в)	c	đ	c	c	
) ₂ (g)	17	17 ⁸ ,31	17 ⁸ ,31	17 ⁸ ,31		K ₂ 0 ₂ (s,1)	17 ^b	đ	d,e		17 ^h
₃ (g)	25	25	25	25		K202(8)	c	đ	e	c	
i(s,1)	48 ^b	48,31	48,31	48		Rb(s,1)	48 ^b		31,48	31,48	48
i(g)_	17_	48	.17.	17		Rb(- g)	108	-48-	-108-	-106	108
i0(g)	c	105	e	c		RbO(g)	c	d,e	c	c	
LiO ₂ (g)	c	d,e	c	c		RbO ₂ (s,1)	66 _p	36 ¹	66	66	67 ^h
Li ₂ (g)	77 ^d	78	77 ^d	77 ^đ		RbO ₂ (g)	c	d,e	c	c	
Li ₂ 0(s,1)	17 ^b	106	17	17	46 ^h	Rb ₂ (g)	c,d	74 [£]	c,d	c,d	
ii ₂ 0(g)	c	đ	c	c		Rb ₂ 0(s,1)	đ,e	· 16	đ,e		d,e,h
1 ₂ 0 ₂ (s,1)	17 ^b	17	d,e		•	Rb ₂ 0(g)	c	82 8	c	c	
i ₂ 0 ₂ (g)	c	å .	c	c .		Rb ₂ O ₂ (s,1)	d,e	68 8	d,e	d,e	67 ^h
₫ ₃ (g)	c	95 [£]	c .	c		Rb202(8)	c	d,e	c	¢	
1 ₃ 0(g)	С	d	с	c		Cs(s,1)	17 ^b		17,31	17,31	109
a(s,1)	b,đ		48,31	48,31	48	Ca(6)	17	17	17	17	10)
-(0)-)					•-			à,e	c	c	
a(g)	17	48	17	17		CsO(g)	¢			·	
a0(g)	c	76	c	C		CsO ₂ (s,1)	e	36 ¹	e		70 ^h
a0 ₂ (s,1)	17 ^b	17	17	17	17 ^h	CsO ₂ (g)	C	d,e	c	c	
10 ₂ (g)	c	d,e	c	C		Cs ₂ (g)	17 ^d	74 [£]	17ª	17 ^d	
a ₂ (g)	7 7	74 ^f	77	77		Cs ₂ O(s,1)	72 ^b	110	72	72	111 ^h
a ₂ 0(s,1)	50 ^b , d	ď	17	17	53	Ca ₂ O(g)	C	d,e	c .	G	•
a ₂ 0(g)	c	đ	c	c		Cs ₂ O ₂ (s,1)		73 ^d .g	d,e		70 ^h
1 ₂ 0 ₂ (s,1)	17 ^b	17 ¹	17 ,	17	17h	Cs ₂ O ₂ (g)	c	d,e	c	c	
a ₂ 0 ₂ (g)	c	d,e	c	c		<u> </u>					
(s,1)	48 ^b		48,31	48,31	48		1, 1977 supple				
(g)	17	48	17	17			lues were estinated from molec			xt.	
)(g)	c	107 ^đ	c	c		d See tex					
0 ₂ (s,1)	18 ^b	36 ¹	18	18	57	e Estimat f The ent	ed. halpy of forms:	tion was cal	culated fr	om the dissocia	ition
0 ₂ (g)	c	d,e	c	c		energy.			*		
2(g)	77	74 [£]	77	77			om the cited so			culate value.	
20(s,1)	đ,e	17 ^d .	d,e		d,e,h	2110 0111	charpy of ruston			iary data.	

2.2. Selected Thermodynamic Data

Selected values of the thermochemical properties of substances encountered in the high temperature vaporization of alkali metal oxides are shown in Tables 1–4. Where possible, sources of the data in Tables 1–4 are given in Table 5. Table 1 lists the solid high temperature alkali metal oxide phases and their melting or dissociation temperatures. Table 2 presents Gibbs energy functions of oxygen and alkali oxide species fit to the equation

$$-(G_T^{\circ} - H_{298}^{\circ})/RT = A + BT + CT^2 + DT^3.$$
 (1)

Table 3 presents values of the enthalpy increment above 298 K fit to the equation

$$(H_T^{\circ} - H_{298}^{\circ})/R = a + bT + cT^2 + dT^3.$$
 (2)

Values of $\Delta_{\rm f}H_{298}^{\circ}/R$, S_{298}°/R , and $(H_{298}^{\circ}-H_{0}^{\circ})/R$ are presented in Table 4.

The quantity S_T°/R for a substance may be found from the relationship

$$S_T^{\circ}/R = -(G_T^{\circ} - H_{298}^{\circ})/RT + (H_T^{\circ} - H_{298}^{\circ})/RT.$$
(3)

The equations given for $-(G_{\ r}^{\circ}-H_{\ 298}^{\circ})/RT$ and $(H_{\ r}^{\circ}-H_{\ 298}^{\circ})/R$ represent the tabulated values within typical limits of \pm 0.01 or \pm 0.01 kK over their applicable temperature ranges. The equations were derived by fitting tabular data for the particular functions cited; they are not direct integrations of heat capacity equations. Because of the nature of the polynomial fits, values derived by differentiating these equations do not accurately represent thermodynamic properties, e.g., heat capacities.

2.3. Condensed Phase Thermochemical Data

This section gives background information on alkali oxide condensed phases and additional information on the selected thermodynamic data presented in Tables 1–4 and referenced in Table 5.

a. Li-O Phases

The Li–O phases important above room temperature are ${\rm Li_2O}$ and ${\rm Li_2O_2}$. ³⁴ Lithium superoxide ${\rm LiO_2}$ and ozonide ${\rm LiO_3}$ are reported to be unstable, even at relatively high oxygen pressures and low temperatures. ^{34–38}

Li(s,l)—thermodynamic functions for the liquid above 1300 K were calculated for an estimated heat capacity $C_p^{\circ}/R = 3.45$. The enthalpy of fusion of Malaspina *et al.*³⁹ is in good agreement with the selected value.

Li₂O(s,l)—values of the enthalpy and Gibbs energy function above 1125 K were calculated by integrating estimated heat capacities including the effect of a diffuse disordering transition⁴⁰ in the antifluorite-type crystal lattice.¹ The recent heat capacity measurements of Tanifuji and Nasu⁴¹ agree with the selected values within experimental error. JANAF¹⁷ selected 1843 K as the melting point of Li₂O on the basis of the work of Van Arkel *et al.*⁴² The value shown in Table 1 was taken from the work of Brewer and Margrave⁴³ and the subsequent determinations by Papin *et al.*,⁴⁴ Akiyama *et al.*,⁴⁵ and Ortman and Larsen.⁴⁶

 $\text{Li}_2\text{O}_2(s,l)$ —the melting temperature, enthalpy of fusion, and liquid heat capacity were estimated.

b. Na-O Phases

The Na-O phases important above room temperature are Na₂O, Na₂O₂, and NaO₂. ^{34,35} The data on the thermodynamics of liquid sodium oxide solutions by Leffler and Wiederhorn ⁴⁷ were not reviewed in the present study. The selected thermodynamic properties of the individual substances were taken from the sources cited below.

Na(s,l)—the heat capacities of Hultgren et al.⁴⁸ were accepted below 1100 K; above this temperature they were smoothly merged with the values of Fredrickson and Chasanov,⁴⁹ which extended to 1500 K. Above this temperature, the heat capacity of Na(l) was estimated to be $C_p^{\circ}/R = 3.80$. The heat capacities were integrated to give values of the Gibbs energy function and enthalpy increments above 298.15 K.

NaO₂(s,l)—the thermodynamic properties of the liquid were calculated using estimated values of the enthalpy of fusion and liquid heat capacity.

Na₂O(s,l)—values of the Gibbs energy function and enthalpy increments between 298.15 and 1300 K were calculated by integrating the heat capacities of Fredrickson and Chasanov. 50 Na₂O(s) is expected to undergo a diffuse disordering transition of its antifluorite crystal lattice^{1,40} above the temperatures of the heat capacity measurements of Fredrickson and Chasanov and of Grimley and Margrave. 51 Estimated heat capacity values that included contributions of the disordering transition estimated by analogy with $K_2S(s)$ were used to calculate values of the Gibbs energy function and enthalpy increments above 1300 K. Neither Fredrickson and Chasanov⁵⁰ nor Henry et al.⁵² observed the solidphase transitions reported by Bouaziz et al.53 at 1023.2 and 1243.2 K, which were not included in the present work. The selected enthalpy of formation was derived from the measurements of O'Hare⁵⁴ and Gross and Wilson.⁵⁵

Na₂O₂(s,l)—estimated values of the enthalpy of fusion and liquid heat capacity were used to derive Gibbs energy function and enthalpy increment values for the liquid.

c. K-O Phases

The solid phases of the K–O system were reviewed by Elliott³⁴ and Shunk,⁵⁶ and more recently by Byker et al.⁵⁷ The phases reported are K_2O , K_2O_2 , KO_2 , and KO_3 . KO_3 reportedly decomposes at 333 ± 2 K⁵⁶; thus, it is not a factor in high temperature vaporization processes. Data on the potassium oxides is limited and of uncertain quality. Different observers have reported significantly different vaporization behavior for solid K_2O and K_2O_2 ; in particular, Leffler and Wiederhorn⁴⁷ and Petrocelli⁵⁸ observed greater O_2 dissociation pressures over K_2O_2 than those calculated from thermodynamic data. The selected enthalpies of formation shown in Table 4 are in good agreement with the NBS¹⁶ and JANAF¹⁷ values and are reasonably consistent with trends in the other alkali oxides. Nevertheless, new studies are needed

K(s,l)—thermodynamic functions were calculated using the estimated heat capacity values $C_p/R = 2.92 + 7.24 \times 10^{-4} T$ above 1300 K.

 $KO_2(s,l)$ —the melting temperature is uncertain because of large impurity effects; experimental values of T_{fus} range

from 653 to 782 K. Byker et al.⁵⁷ estimated $T_{\rm fus} = 865$ K and $\Delta_{\rm fus}H/R = 1.624$ kK from an analysis of phase diagram data. The heat capacity of the liquid was estimated as $C_p/R = 12.6$ for calculation of Gibbs energy function and enthalpy increment values.

K₂O(s,l)—no experimental heat capacity or entropy data were found. The entropy at 298.15 K was estimated by the Latimer relation $S^{\circ}(298 \text{ K})/R = 1.5 \ln M + 3B$, where M is the species gram formula weight and B a constant whose value was interpolated for $K_2O(s)$ on a plot of $\ln M vs B$ for solid alkali metal M_2O oxides. The resulting value, $S^{\circ}(298)$ K)/R = 11.34, is in good agreement with the estimated value given in JANAF.¹⁷ Both JANAF and Eliezer et al.⁵⁹ estimated heat capacities for solid K2O. The JANAF values are too large, and the values of Eliezer et al. are for temperatures over 1000 K. The melting temperature is uncertain; the experimental value of 919 ± 5 K of Natola and Touzain⁶⁰ based on DTA of a 30 mg specimen was rejected because the likely interactions with the Pyrex container and the noncongruent vaporization of K₂O would both lower the observed melting temperature. Simmons et al.61 reported a melting temperature of 1370 K based on analogy to the melting temperature of KCl, NaCl, and Na2O. The thermodynamic functions for condensed K2O of the present work were calculated using heat capacities estimated by analogy with the other alkali M₂O oxides, including the effects of a diffuse disordering transition in the antifluorite crystal lattice. 1,40 The melting temperature given in the present work was estimated by analogy to K₂S and Na₂S (Ref. 17, March 31, 1978 supplement) and Na₂O. The enthalpy of formation and K₂O(s) reported by JANAF¹⁷ and others was taken from the solution calorimetric work of Rengade⁶²; a modern determination is needed.

 $\rm K_2O_2(s,l)$ —heat capacity values given in JANAF¹⁷ were used, along with the value of $S^{\circ}(298~\rm K)/R$ estimated using Latimer's rule, to calculate Gibbs energy function and enthalpy increment values for the solid. The calculated values of the Gibbs energy function were used along with the pressure data of Riley⁶³ and Kazarnovskii and Raikhshtein⁶⁴ for the reaction

$$2KO_2(s) = K_2O_2(s) + O_2(g)$$

to derive the enthalpy of formation. Byker⁵⁷ had done this calculation with the JANAF Gibbs energy values. The JANAF value of the enthalpy of formation, derived from earlier sources, is in good agreement with the selected value.

d. Rb-O Phases

The Rb-O phases important above room temperature are Rb₂O, Rb₂O₂, and RbO₂.^{34,35,56} The ozonide RbO₃ has been reported,³⁴ as well as the suboxide phases Rb₆O and Rb₉O₂⁶⁵; these phases are not stable at high temperatures. Thermodynamic data for individual substances was taken from the sources given below.

Rb(s,l)—thermodynamic values above 1300 K were calculated from estimated heat capacities.

RbO₂(s,l)—Paukov et al.⁶⁶ measured the heat capacity of RbO₂(s) from 13 to 298 K. Their values of heat capacity, entropy, and $(H_{298}^{\circ} - H_{0}^{\circ})/R$ were adopted. The melting

temperature was taken from Centnersziver and Blumenthal.⁶⁷ Gibbs energy function and enthalpy increment values were calculated using estimated values of the heat capacity above 298.15 K and an estimated enthalpy of fusion.

Rb₂O(s,l)—no experimental heat capacity or entropy data were found. The entropy at 298.15 K was estimated by the method used for K_2O . Heat capacity values for the solid were estimated by analogy to the other alkali metal M_2O oxides and included the effects of a diffuse disordering transition in the antifluorite crystal lattice. ^{1,40} The melting temperature was estimated from a plot of T_{fus} vs $\ln M$, where M is the gram formula weight, for the alkali metal M_2O oxides. The heat capacity and enthalpy of transition data were used to calculate Gibbs energy function and enthalpy increment values.

 ${
m Rb_2O_2(s,l)}$ —the heat capacity above 298.15 K was estimated by analogy to that of ${
m Rb_2O}$ and the difference between C_p° values for ${
m Na_2O_2}$ and ${
m Na_2O}$. The entropy at 298.15 K was estimated as for ${
m K_2O_2}$. The liquid heat capacity was estimated. The estimated entropy and heat capacities were used to calculate Gibbs energy function and enthalpy increment values. The enthalpy of formation was calculated using the pressure measurements of Kraus and Petrocelli⁶⁸ for the reaction

$$2RbO_2(s) = Rb_2O_2(s) + O_2(g).$$

This enthalpy of formation was calculated by both secondand third-law methods; the second-law method gave a value of $\Delta_f H^\circ_{298}/R$ 1.6 kK less negative than the selected value. Earlier values of the enthalpy of formation included a value estimated by de Forcrand⁶⁹ and a value derived from the doubtful pressure measurements of Centnersziver and Blumenthal.⁶⁷ The selected value is 0.5 kK more negative than the NBS¹⁶ value and 2.9 kK more negative than the IVTAN²⁶ value.

e. Cs-O Phases

Knights and Phillips⁷⁰ reviewed the literature on condensed phases of the Cs-O system and reported the solid phases Cs₂O, Cs₂O₂, CsO₂, Cs₇O, Cs₄O, Cs₇O₂, Cs₃O, and CsO₃. Simon⁶⁵ reported the phase Cs₁₁O₃ rather than Cs₇O₂. Except for Cs₂O, Cs₂O₂, CsO₂, and Cs₃O, which decomposes peritectically at 437 K, the solid phases undergo decomposition at relatively low temperatures.⁷⁰ Knights and Phillips reported thermodynamic data, which were not evaluated in the present work, for liquid cesium oxide solutions. Sources of data selected for individual substances are given below.

Cs(s,1)—Behrens et al.⁷¹ found that the Gibbs energy functions derived from the heat capacities of Hultgren et al.⁴⁸ are inconsistent with Cs vaporization data. On this basis, the JANAF¹⁷ values of the thermodynamic functions for condensed Cs were adopted. The data were estimated above 1300 K.

CsO₂(s,l)—The entropy at 298 K was estimated using Latimer's rule. Values of the Gibbs energy function and enthalpy increments were calculated from heat capacities estimated by adding the heat capacity of Cs₂O to the heat capacity difference between NaO₂ and Na₂O.

 $Cs_2O(s,l)$ —Flotow and Osborne⁷² measured the heat capacity of Cs_2O from 5 to 350 K, and extrapolated their results to the melting temperature. Their data for the Gibbs energy function, $(H_{298}^{\circ} - H_{0}^{\circ})/R$, and $(H_{T}^{\circ} - H_{298}^{\circ})$ were accepted. Cs_2O crystallizes in a $CdCl_2$ -type rhombohedral structure, ¹ and thus does not undergo the high temperature antifluorite disordering transition expected for the other dialkali monoxides. Thermodynamic functions of the liquid were calculated from estimated heat capacities.

Cs₂O₂(s,l)—the entropy at 298.15 K was estimated by Latimer's rule. Values of the Gibbs energy function and enthalpy increments were calculated from heat capacities estimated by analogy to Cs₂O and the difference between Na₂O and Na₂O₂. The enthalpy of formation was calculated from the dissociation pressure measurements of Berardinelli and Kraus⁷³ for the reaction

$$2CsO_2(s) = Cs_2O_2(s) + O_2(g).$$

Berardinelli and Kraus⁷³ began their measurements with specimen compositions near CsO₂ and monitored composition change by mass loss, which was assumed to be due to O₂ only. This assumption, valid for the decomposition of CsO₂ to Cs₂O₂, is demonstrably wrong for vaporization of Cs₂O₂; thus, the pressure data reported for the reaction

$$Cs_2O_2(s) = Cs_2O(s) + 1/2O_2(g)$$

probably do not correspond to this equilibrium. IVTAN²⁶ reported the value $\Delta_{\rm f} H_{298}^{\circ}/R = -54 \pm 3$ kK, probably based on the estimation of de Forcrand.⁶⁹

2.4. Vapor Species Thermochemical Data

a. Discussion

This section gives background information on alkali oxide gas species and additional information on the selected thermodynamic data presented in Tables 1–4 and referenced in Table 5. The alkali metal—oxygen vapor species reviewed below are grouped by generic molecular formula rather than by their alkali metal constituent because of the generally related molecular properties of the homologous species.

b. M(g) Species

M(g) species—Hultgren's enthalpies of formation,⁴⁸ which incorporated data not reviewed by JANAF¹⁷ were used.

c. MO(g) Species

The spectroscopic constants of the gaseous alkali monoxides are based almost entirely on *ab initio* calculations, as summarized for the most part by Huber and Herzberg. Additionally, the vibrational frequency of RbO and the internuclear distance in CsO were estimated by analogy with trends in the diatomic alkali fluorides, chlorides, and oxides, while the energy of the first excited electronic state of CsO was estimated by extrapolating the values for the lighter alkalis. The selected constants used in calculating thermodynamic functions are summarized in Table 6.

Because of the scarcity of thermochemical data for all of the gaseous monoxides but LiO, the dissociation energies of the MO species were calculated by means of the Rittner polarizable ion model^{29,30} as an aid in evaluating the experi-

Table 6. Molecular constants of the alkali metal MO vapor species. a

	r _e x10 ⁸ (cm)	ω (cm ⁻¹)	Electronic ground state degeneracy	Electronic excited state degeneracy	T _e (cm ⁻¹)
LiO	1.695	852	2	. 2	2330
NaO	2.05	500 ^b	2	2 ·	1500
ко	2.22	384	2	2	347
RЪО	2.28	340°	2	2	606
Cs0	2.39°	314 ^d	2	2	1000°

- a Unless otherwise indicated, values were taken from reference 74.
- b From the calculations of O'Hare and Wahl112.
- c Estimated by the method given in the text.
- d From matrix isolation studies of Spiker and Andrews¹¹³.

mental results and fixing trends in the series. Input parameters for the Rittner calculations are given in Table 7, while the calculated and experimental dissociation energies are compared in Table 8. In view of the reasonably good agreement between the two values for LiO, NaO, and KO, the Rittner calculated values for RbO and CsO were selected for evaluating the standard-enthalpies-of-formation.

d. MO₂(g) Species

The gaseous alkali dioxide species LiO₂, NaO₂, KO₂, RbO₂, and CsO₂ have all been identified by low temperature matrix isolation spectroscopy, but these species have not been directly observed by mass spectrometry in any of the investigations reported to date. However, it seemed prudent

Table 7. Values of atomic and molecular constants used in
Rittner electrostatic model calculations of the
dissociation energies of alkali metal MO vapor
species. A

	r x 10 ⁸ b	α(M ⁺)x10 ^{24 c}	k x 10 ^{-5 d} (dyne/cm)	IP ^e (eV)
Li0	1.695	0.03	2.08	5.392
Na0	2.05	0.20	1.54	5.139
ко	2.217	1.00	1.00	4.341
RbO	2.278	1.70	1.00	4.197
Cs0	2.39	2,50	0.90	3.894

- a r-internuclear distance; α -dipole polarizability; k--force constant; IP--ionization potential; the electron affinity of 0 was 1.465 eV¹⁷; the polarizability of 0 was 1.0 x 10⁻²⁴ cm³.105
- b References in table 5.
- c Reference 114.
- d Calculated from constants given in table 6.
- e Reference 115.

Table 8. Enthalpies of formation and dissociation energies of alkali metal MO vapor species.

	△ _f H° ₂₉₈ /R (kK)	D ^O O/R (kK)	Method	Reference
LiO	9.1 <u>+</u> 1.0	39.6 <u>+</u> 1.0	Experimental	94
	8.3 <u>+</u> 0.8	40.4 <u>+</u> 0.8	Experimental	105
	8.6 <u>+</u> 2.5	40.1 <u>+</u> 2.5	Experimental	80
	8.7 <u>+</u> 0.5	40.0 <u>+</u> 0.5	Rittner model	This work
	8.3 <u>+</u> 0.8	40.4 <u>+</u> 0.5	(Selected value)	
NaO	12.6 <u>+</u> 0.3	29.8 <u>+</u> 0.3	Experimental	76
	13.3 <u>+</u> 0.5	29.1 <u>+</u> 0.5	Rittner model	This work
	12.6 <u>+</u> 0.3	29.8 <u>+</u> 0.3	(Selected value)	
ко	7.4 <u>+</u> 2.5	33.0 <u>+</u> 2.5	Experimental	107
	7.2 <u>+</u> 0.5	33.2 <u>+</u> 0.5	Rittner model	This work
	7.2 -0.5	33.2 10.5	(Selected Value)	
RbO	7.1 <u>+</u> 0.5	32.8 <u>+</u> 0.5	Rittner model	This work
			(Selected	Value)
Cs0	4.9 <u>+</u> 0.5	33.9 <u>+</u> 0.5	Rittner model	This work
			(Selected	Value)

to include data for these species, since they may be important under oxidizing conditions. The molecular constants used to calculate thermodynamic functions for the alkali MO_2 vapor species are shown in Table 9.

The only experimental enthalpy data found for these

Table 9. Molecular constants of the alkali metal mo_2 vapor species.^a

, 	r(M-0)x10 ⁸		ω ₁ (cm ⁻¹)	ω ₂ (cm ⁻¹)	ω ₃ (cm ⁻¹)
LiO2 ^{b,c,d}	1.77	44	1097	699	492
NaO2 ^{c,e,f}	2.07	37.5	1080	391	333
KO2c.g	2.28	32.6	1108	308	300
RbO2c,g,h	2.57	28.8	1110	255	283
CsO2c,g,h	2.67	27.7	1114	237	269

a The molecules have $C_{2\gamma}$ geometry, a symmetry number of 2, and were assumed to have an electronic ground state degeneracy of 2.

species are those for NaO₂ deduced from the flame studies of McEwan and Phillips⁷⁵ who reported data for the reaction

$$NaO_2(g) = Na(g) + O_2(g)$$
.

However, these data predict NaO₂(g) to be a major species over Na₂O(s) under neutral Knudsen cell conditions, contrary to observations.⁷⁶ We believe that the flame data on NaO₂(g), which are indirect, are in error.

Since none of the MO₂ vapor species has been observed under neutral Knudsen cell conditions, lower limits for the enthalpies of formation were estimated by considering the reaction

$$2MO_2(g) = M_2O(g) + 3/2O_2(g),$$

for which the Gibbs energy change is given by the equation

$$\Delta G/RT = \Delta \left[-(G_T^{\circ} - H_{298}^{\circ})/RT + \Delta_f H_{298}^{\circ}/R \right]$$

= \ln[p(M_2O)p(O_2)^{3/2}/p(MO_2)^2], (4)

and taking 10^{-10} bar as the limit of detection in the mass spectrometric experiments. Third-law calculations combined with the established M_2O data yield the lower limit values for the MO_2 enthalpies of formation. The values are uncertain, since detection limits in the mass spectrometric investigations are somewhat variable.

e. M₂(g) Species

M₂(g) species—the low dissociation energies of the diatomic alkali metal species require that the statistical methods used to evaluate thermodynamic functions from molecular data utilize accurate rotation-vibration energy levels and correct cutoff procedures for dissociating states. Several groups have performed these calculations for Li2, Na2, and K_2 ; the calculations of Feber and Herrick⁷⁷ were used in the present work. In the absence of more accurate calculations, data for Cs₂(g) were taken from JANAF. ¹⁷ Spectroscopic data from Hultgren et al.48 were used to calculate the Gibbs energy functions for Rb₂ given in the present work. Only the nondissociating ground state and one excited state were used in these calculations. Since the diatomic elemental species are minor components of the vapor, these approximations should not introduce significant error into the vapor pressure calculations. The enthalpy of formation of Li₂ was derived from the accurate ab initio dissociation energy of Konowalow and Olson. 78 The enthalpies of formation of the remaining M2 species were calculated from the dissociation energies compiled by Huber and Herzberg.74 The data of Piacente et al. 79 were also considered for Rb₂.

f. M₂O(g) Species

On the whole, the gaseous dialkali monoxide species have been studied in considerably more detail than the other alkali oxide species. Nevertheless, there are considerable uncertainties in some of the results, and the molecular constants needed for calculation of thermodynamic functions are incomplete. The equilibrium data for the Li and Na species appear to be reliable; those for the K, Rb, and Cs species are less certain, and further experimental investigations are needed.

Table 10 lists the selected molecular constants used in calculating thermodynamic functions. Experimentally un-

b Reference 86.

c Reference 89.

d Reference 116.

g Reference 88.

e Reference 87.

h Reference 118.

f Reference 117.

Table 10. Molecular constants of the alkali metal M_20 vapor species.^a

	r(M-0)x10 ⁸ <m-0-m< th=""><th>ω_1</th><th>ω_2</th><th colspan="2">ω_3</th></m-0-m<>		ω_1	ω_2	ω_3	
	(cm)	(deg)	(cm ⁻¹)	(em ⁻¹)	(cm ⁻¹)	
Li ₂ 0	1.60 ^b	180°	(760)°	(140) ^{c,d}	987°	
Na ₂ 0	(1.95)	(180)	(450)	(100) ^d	(600)	
к ₂ 0	(2.12)	160°	(301)	(80)	500 ¢	
Rb ₂ 0	(2.17)	160 ^e	(284)	(70)	473 ^e	
Cs ₂ 0	(2.28)	140 ^e	(305)	(60)	455 ^e	

a r--internuclear distance; <--angle; ω_1 , ω_2 , ω_3 --wave numbers of fundamental vibrations. The symmetry number was assumed to be 2. The electronic ground state was assumed to be nondegenerate, and no excited electronic states were included in the calculations. Values in parentheses are estimated, and those not referenced are from the present study.

known M-O distances were estimated using the ratio of Li-O distances in LiO and Li₂O and the known or estimated M-O distance in the monoxides. Except for Na₂O, the unsymmetrical stretching frequency has been observed by matrix infrared spectroscopy. For Na₂O, this frequency was estimated by interpolation. For the Li and Na species, estimated values of the other vibrational frequencies were taken from White et al. ⁸⁰ and from Hildenbrand and Murad. ⁷⁶ Values for the K, Rb, and Cs species were estimated by analogy to the values for the Li and Na species.

The selected enthalpies of formation of the M_2O species shown in Table 4 are all based on equilibrium measurements. For Li₂O(g), the value was taken from the concordant results of Hildenbrand *et al.*, ⁸¹ White *et al.*, ⁸⁰ Norman and Winchell, ⁸² and Ikeda *et al.* ⁸³ The value for Na₂O was taken from Hildenbrand and Murad⁷³ and Norman and Winchell. ⁸² The results of Gorokhov *et al.* ^{84,85} and of Norman and Winchell were used to derive the selected enthalpies of formation for Rb₂O(g) and Cs₂O(g).

g. M₂O₂(g) Species

Data for the alkali metal M_2O_2 vapor species are less complete and less certain than for the corresponding M_2O or MO species. All have been observed by low temperature matrix isolation spectroscopy, but the fundamental vibrational frequencies have been only partially assigned, and information about molecular geometries is incomplete. Mass spectrometric studies of reaction equilibria involving Li_2O_2 , K_2O_2 , and Cs_2O_2 have been made, but no data is available for Na_2O_2 or Rb_2O_2 .

Table 11. Molecular constants for the alkali metal M₂O₂
vapor species.^a

	r(0-0)x10 ⁸	r(M-M)x10 ⁸	<0-M-0	I _x I _y I _z x10 ¹¹⁷	
	(cm)	(cm)	(deg)	(g ³ cm ⁶)	
Li ₂ 0 ₂	1.50 ^b	2.91°	54	1.16x10 ²	
Na ₂ 0 ₂	1.50	3.67	44	2.21x10 ³	
K ₂ 02	1.50	4.02	41	8.70x10 ³	
Rb ₂ 0 ₂	1.50	4.15	40	4.56x10 ⁴	
Cs ₂ O ₂	1.50	4.37	38	1.35x10 ⁵	

a Unless otherwise indicated, values were estimated as described

The selected molecular constants are shown in Table 11. In accordance with the matrix spectra of Andrews and co-workers, $^{86-89}$ the M_2O_2 species were regarded as planar D_{2h} molecules with symmetry numbers of 4. The O-O distances in the molecules were taken as 1.5×10^{-8} cm, in accord with data on the solid alkali peroxides 35 and the *ab initio* results of Yates and Pitzer 90 on Li_2O_2 . The M-O distances were estimated by multiplying the value of r(M-O) of the MO(g) species by the ratio $r(\text{M-O},\text{Li}_2O_2,g)/r(\text{M-O},\text{Li}_O,g) = 1.64/1.695. ^{90,91}$

The selected vibrational frequencies of the M_2O_2 molecules are given in Table 12. Unobserved vibrational frequencies were estimated by multiplying the frequencies of $\text{Li}_2O_2^{86,90}$ by the ratio of comparable frequencies in the species M_2O to those in Li_2O . For this correlation, the A_g , A'_g , and B_{1g} M_2O_2 modes were assumed to scale as the M_2O ω_1 mode; the B_{1u} mode of M_2O_2 as the ω_2 mode of M_2O ; and the B_{3u} mode of M_2O_2 as the ω_3 mode of M_2O . Experimental values are available for the B_{2u} modes for each of the M_2O_2 species, so that predicted and observed values can be compared; except for Na_2O_2 , where the predicted frequency is high by 43 cm⁻¹, the agreement is within 21 cm⁻¹.

Enthalpies of formation derived from mass spectrometric data were selected from the concordant experimental values for Li₂O₂, ^{80,92-94} K₂O₂, ⁸⁵ and Cs₂O₂. ^{82,85} For Na₂O₂ and Rb₂O₂, enthalpies of formation were estimated by interpolating the energy of the reaction.

$$M_2O_2(g) = 2MO(g),$$

via a plot $[D_0(M_2O_2 - 2D_0(MO)]$ vs $D_0(MO)$ for the alkali metals.

b Reference 119.

c Reference 17.

d Doubly degenerate vibrational mode.

e Reference 113.

in the text. The symmetry was assumed to be $\mathrm{D}_{2\mathrm{h}}$, with a symmetry number of 4. The electronic ground state was assumed to be nondegenerate, and no excited electronic states were included in the calculations of thermodynemic functions.

b Reference 90.

C Reference 86.

Table 12. Vibrational fundamentals of alkali metal ${\rm M_2O_2} \ \, {\rm vapor \ species.} \ \, {\rm (Units: \ cm^{-1})}$

	Vibrational Mode							
	^B 1u	B _{2u}	B _{3u}	Ag	A'g	B _{1g}		
Li ₂ 0 ₂	94 ⁸	796 ^b	446 ^b	954 ^a	741 ⁸	639ª		
Na ₂ 02	70°	525 ^d	254 ^đ	565°	440°	380°		
K ₂ 0 ₂	55 °	433 ^e	226°	-380°	290 ^c	250°		
Rb ₂ 0 ₂	47 ^c	389 ^e	214 ^C	356°	280°	240 ^C		
Cs ₂ 0 ₂	40°	357 [£]	206°	380°	300°	260°		

- a Reference 90. ab initio calculations.
- b Reference 86, experimental.
- C This work, estimated.
- d Reference 87, experimental.
- e Reference 88, experimental.
- f Reference 89, experimental.

h. Other Species

The species Li₃ was reported in a mass spectrometric study of the vaporization of pure lithium. Skimoto et al. Skimoto et a

The species Li₃O has been reported in mass spectrometric investigations of the vaporization of Li₂O by Wu et al., ⁹² Ikeda et al., ⁸³ and Kimura et al. ⁹⁴ The pressures reported for this species were near instrumental sensitivity limits and differ significantly among the different investigators. Approximate thermodynamic functions for Li₃O(g) were estimated from the molecular constants for Li₂O₂(g) and rhombic Li₄ clusters ^{99,100} and were used in third-law calculations to estimate the enthalpy of formation from the vaporization data of Wu et al. and of Kimura et al.

3. High Temperature Equilibria and Vaporization Rates

3.1. Equilibrium Properties of Alkali Oxides

The solid phases relevant to high temperature vaporizaiton processes are the M_2O dialkali monoxide, the M_2O_2 peroxide, and the MO_2 superoxide. Phases of the composition M_2O_3 have been reported for alkali metals except lithium, 35 but their existence as pure equilibrium phases has not been substantiated. 1,56,57,70 Several general trends exist in the properties of alkali oxides. The higher oxide condensed phases become more stable relative to the dialkali monoxide

as the size of the metal ion increases. This trend is illustrated by the synthetic methods used for oxide production. 35,57,72,101 Direct oxidation of the metal at moderate temperatures in dry air produces Li₂O, Na₂O₂, KO₂, RbO₂, and CsO₂. NaO₂ is produced by high temperature reaction of Na₂O₂ with high pressure oxygen. Thermal decomposition of the corresponding superoxides yields K_2O_2 , Rb₂O₂, and Cs₂O₂, and decomposition of Na₂O₂ gives Na₂O. The dial-kali monoxides of K, Rb, and Cs are produced by reacting the peroxides with excess metal, then removing the remaining metal by distillation.

For lithium, the only important high temperature solid oxide is the dialkali monoxide. For sodium, the dialkali monoxide is stable under neutral or reducing conditions above about 400 K, but in dry air the peroxide is stable up to its melting point. Either the peroxide or superoxide of K, Rb, or Cs is stable under neutral or oxidizing conditions at high temperatures. The existence of several possible oxide phases complicates the study and description of vaporization processes for K, Rb, and Cs oxides, especially when low-melting liquids are considered.

The information needed to characterize the vaporization behavior of oxides was discussed by Brewer in his 1953 review. The most important facts to be established are the identities of the important vapor species and the effects of vaporization on condensed phase compositions. When these two factors are known, the vaporization equilibria can be described and quantitatively calculated if the necessary thermodynamic data are available. The important vapor species are reasonably well established for alkali metal—oxygen systems; unfortunately, changes in condensed phase composition during vaporization are only partially known.

According to the phase rule, the intensive variables of an alkali metal-oxygen system depend only on temperature if two condensed phases and one gas phase are present. If only one condensed phase and a gas phase are present, one additional variable must be specified to define the state of the system. This requirement can be met by specifying the oxygen potential or the elemental composition of the vapor phase. Congruent vaporization, in which the overall elemental composition of condensed and vapor phases are the same, is a particular example of a one-solid-phase vaporization process in which the state of the chemical system is specified by temperature and gas phase composition. Experimentally, solid Li₂O and Na₂O vaporize congruently under neutral conditions. ^{76,81,102} Thermodynamic data indicate that solid K₂O₂, Rb₂O₂, and Cs₂O₂ also vaporize congruently.

3.2. Mathematics of Vaporization Calculations

The calculations that follow are applicable to the vaporization of pure solid and liquid phases under conditions where condensed phase composition changes are negligible; vaporization from solution phases is beyond the scope of the present work, and the necessary data are for the most part unavailable. The properties considered are the vapor pressures of principal vapor species and the rate of mass loss per unit surface area, both as functions of temperature.

For many vaporization conditions one or two vapor species are predominant, and calculation of the chemical equilibrium from thermodynamic data is straightforward. Calculations can be more complicated when several species are present in significant quantities. The following procedure is a general method for calculating the equilibrium pressures of vapor species in equilibrium with a condensed phase of known composition.

The chemical species present at equilibrium in the vaporization of a condensed binary oxide can be described mathematically as being formed by linear combination of the vaporizing oxide $M_b O_c$ and oxygen O_2 according to the general reaction for the production of species i,

$$d(i)M_b O_c(s,l) + e(i)O_2(g) = M_x(i)O_v(i)(s,l,or g),$$

where e(i) may be zero or a negative number. The possible species include the vaporizing phase M_bO_c (s or 1), considered to be species 1, and $O_2(g)$ as well as the species produced by their linear combination. From mass balance,

$$d(i) = x(i)/b, (5)$$

and

$$e(i) = [by(i) - cx(i)]/2b.$$
 (6)

We define the Gibbs energy of species i at temperature T as

$$G_T^{\circ}(i)/RT = (G_T^{\circ} - H_{298}^{\circ})/RT - \Delta_f H_{298}^{\circ}(i)/RT.$$
 (7)

The Gibbs energy change for the reaction producing gaseous species i from $M_b O_c$ (s or l) and O_2 (g), $\Delta_r G^{\circ}_T$ (i), is given by

$$\Delta_{r}G_{T}^{\circ}(\mathbf{i})/RT = G_{T}^{\circ}(\mathbf{i})/RT - d(\mathbf{i})G_{T}^{\circ}[\mathbf{M}_{b}O_{c}(\mathbf{s} \text{ or } \mathbf{l})]/KI$$
$$= e(\mathbf{i})G_{T}^{\circ}[O_{2}(\mathbf{g})]/RT, \tag{8}$$

and the equilibrium constant for producing gas species i is

$$K(i) = \exp\left[-\Delta_r G_T^{\circ}(i)/RT\right] = p(i)/\{a(1)^{d(i)}p[O_2(g)]^{e(i)}\},$$
(9)

where the pressures p(i) are considered low enough that they equal fugacities, and a(1) is the activity of $M_b O_c$. If $M_b O_c$ is a pure condensed phase, its activity is unity, and

$$p(\mathbf{i}) = K(\mathbf{i})p(\mathbf{O}_2)^{e(\mathbf{i})},\tag{10}$$

for the gas species. If the oxygen pressure is known, the pressure of each species i can be calculated from this equation.

For congruent vaporization, the vapor composition is the same as the condensed phase, and

$$\sum [p(\mathbf{i})x(\mathbf{i})] / \sum [p(\mathbf{i})y(\mathbf{i})] = b/c.$$
 (11)

This leads to the expression

$$\sum [K(i)p(O_2)^{e(i)}\{cx(i) - by(i)\}] = 0,$$
 (12)

where summation is over the gaseous species present. In practice, this equation is solved for $p(O_2)$ by graphical interpolation or by numerical methods. The value found for $p(O_2)$ is then used in Eq. (10) to find the other partial pressures.

For the above calculations to correspond to equilibrium vaporization, they must be used for vaporization of the condensed oxide phase or phases stable under the particular vaporization conditions. Where more than one oxide phase is a possibility, the stable phase or phases must be determined.

Table 13. Oxygen pressures of two-solid-phase regions.

	K	Q	10 ³ R	10 ⁻³ s	10 ⁻⁶ ¥	T Range
Li ₂ 0-Li ₂ 0 ₂	879.52	-311.10	117.36	-54.94	3.21	298-450
Na ₂ 0-Na ₂ 0 ₂	-114.32	41.37	-11.03	0.22	-0.94	298-948
K20-K202	288.04	-95•49	25.88	-34.13	1.53	298-763
к ₂ 0 ₂ -ко ₂	25.25	-8.44	3.62	-4.06	-0.10	298-763
Rb ₂ 0_Rb ² O ²	585.01	-194.92	50.91	-59-31	3.75	350-843
Rb202-Rb02	176.06	-60.68	19.69	-14.65	0.65	298-685
Cs ₂ 0-Cs ₂ 0 ₂	-94.23	36.87	-12.67	-8.96	-0.54	350-768
Cs ₂ 0 ₂ -Cs0 ₂	143.81	-49.06	15.08	-12.83	0.59	298-723

Two condensed oxide phases of fixed composition at a given temperature are related by the equilibrium

Oxide (1) = Oxide (2) +
$$fO_2(g)$$
.

Oxide (1) is assumed to have the greater oxygen activity. The equilibrium constant for this reaction is

$$K_{eq} = a(2)p(O_2)^f/a(1).$$
 (13)

If both condensed phases are present at equilibrium, their activities are unity, and the pressure of O_2 is given by the equation

$$p_{\rm eq}(O_2) = K_{\rm eq}^{(1/f)}$$
. (14)

Oxide (1) is the stable phase for O_2 pressure greater than $p_{eq}(O_2)$, and oxide (2) is stable for O_2 pressures less than $p_{eq}(O_2)$ at the given temperature. The calculated oxygen pressures for two-solid-phase regions are given as functions of temperature in Table 13.

3.3. Vaporization Rate Calculations

Maximum vaporization rates were calculated using the Hertz-Knudsen equation of classical kinetic theory, which applies to free vaporization from an uncontaminated surface at low pressures. For gas species i evolving from the surface, the maximum vaporization rate dn(i)/dt in mol cm⁻² s⁻¹ is given by the equation

$$dn(i)/dt = p(i)/[2\pi M(i)RT]^{1/2},$$
 (15)

where M (i) and p(i) are the gram molecular weight and equilibrium partial pressure of vapor species i, and R and T are the molar gas constant and absolute temperature. Since 1 bar is 10^6 dyne cm⁻² and R is 8.3144×10^6 erg mol⁻¹ deg⁻¹, for pressures in bar, this is equivalent to

$$dn(i)/dt = 43.8\{p(i)/[M(i)T]^{1/2}\} \text{ mol cm}^{-2} \text{ s}^{-1}.$$
 (16)

Calculation of maximum vaporization rates is more difficult for noncongruent than for congruent processes. In general, solid phase activities during the vaporization process are not known, and calculated maximum vaporization rates of oxides that deviate from their original stoichiometry should be considered indicative rather than exact. ¹⁰³ In the present work, maximum vaporization rates have been calcu-

lated for noncongruent, as well as congruent, vaporization. Only initial vaporization rates, for which no condensed phase composition changes have taken place, are considered. Vaporization rates are expressed in terms of the mass of originally present oxide $M_b\,O_c$ lost as measured by the volatilization of metal atoms. The total mass loss thus calculated is

$$dm/dt = 44.3 [M(M_b O_c)/b] \sum \{p(i)x(i)/[M(i)T]^{1/2}\},$$
(17)

where dm/dt is the mass loss rate in $g \text{ cm}^{-2} \text{ s}^{-1}$, $M_b O_c$ is the formula of the condensed oxide, and x(i) and M(i) are the number of gram atoms of metal per mole and gram molecular weight of gas species i.

Although actual vaporization may be slower, the calculated maximum vaporization rate is valuable as a rough guide to vaporization kinetics. Recent references on maximum vaporization rate calculations include works by Turkdogan 104 and Beruto $et\ al.$

3.4. Vaporization Behavior under Representative Conditions

The calculated partial pressures of major species in equilibrium with stable alkali metal oxides are shown in the figures of this section as plots of $\log P$ vs 1/T for vaporization under conditions representative of those likely to be encountered experimentally. For each alkali metal-oxygen system, equilibrium partial pressures were calculated for vaporization in 0.2 bar of O₂, representing oxidizing conditions, for vaporization at a fixed oxygen potential of 10⁻¹⁵ bar, representing reducing conditions, and for vaporization of solid under neutral conditions. The calculations do not cover certain situations of practical interest. One of these is the vaporization of alkali metal-oxygen liquid solution phases, for which adequate thermodynamic data do not exist. Another is vaporization under conditions of decreased condensed phase activity, as in solutions of alkali oxides in silicate melts, which is beyond the scope of this work.

Species representing less than about 10^{-6} of the total pressure were omitted from the graphs, as were the gaseous dioxides, for which only upper limits of partial pressures could be calculated. The limiting stabilities indicate that the dioxides may be major species under oxidizing conditions, and it is important that their properties be determined.

a. Li-O System

Lithium oxide Li₂O is the stable condensed oxide phase at temperatures above about 380 K for vaporization at O₂ pressures up to 0.2 bar. The partial pressures of the major species in equilibrium with condensed Li₂O are shown in Figs. 1–3 for representative vaporization conditions. Liquid Li₂O was assumed to vaporize congruently under neutral conditions. Calculated maximum mass loss rates are shown in Fig. 4.

b. Na-O System

The equilibrium oxygen pressure for equilibrium of $Na_2O_2(s)$ with NaO_2 is greater than 100 bar above 298 K, and disodium monoxide Na_2O is stable above about 400 K for

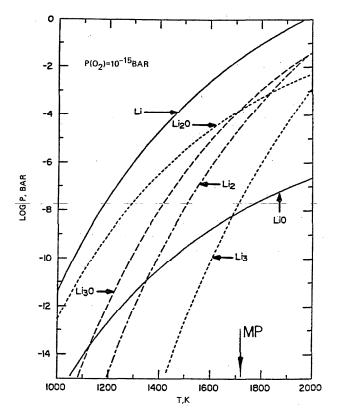


Fig. 1. Li₂O vaporization in 10⁻¹⁵ bar O₂.

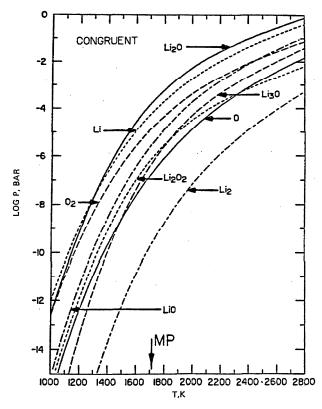


Fig. 2. Li₂O congruent vaporization.

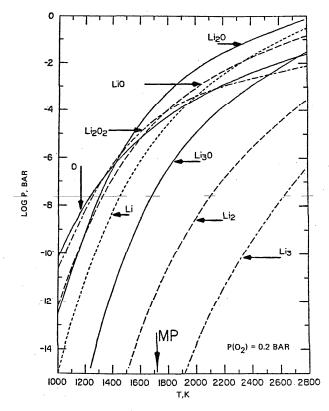


Fig. 3. Li₂O vaporization in 0.2 bar O₂.

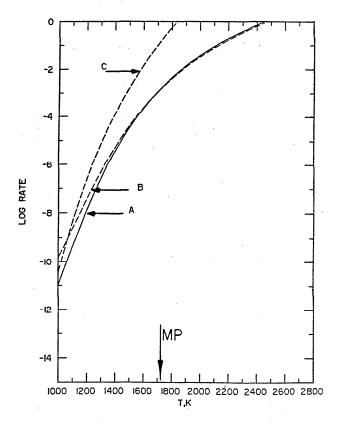


FIG. 4. Li₂O maximum vaporization rates. A—0.2 bar O_2 ; B—congruent vaporization; C— 10^{-15} bar O_2 . Rate units—g cm⁻² s⁻¹.

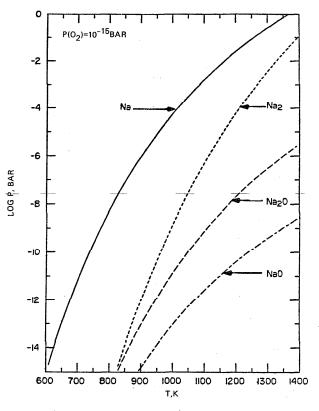


Fig. 5. Na₂O vaporization in 10⁻¹⁵ bar O₂.

congruent vaporization conditions. Liquid Na_2O was assumed to vaporize congruently under neutral conditions. The partial pressures of the major species in equilibrium with condensed Na_2O are shown in Figs. 5 and 6 for vaporization at an O_2 pressure of 10^{-15} bar and for congruent vaporization, respectively. The partial pressure curves in Fig. 7 arc for vaporization of Na_2O_2 solid and liquid below 1150 K, and for vaporization of condensed Na_2O above this temperature, in 0.2 bar of O_2 . Calculated mass loss rates for these vaporization conditions are shown in Fig. 8.

c. K-O System

Very little data are available on the vaporization of condensed potassium, rubidium, and cesium oxides, and the corresponding thermodynamic properties and composition changes during vaporization are known with less certainty than for lithium and sodium oxides. New studies of condensed phase vaporization equilibria and gas species thermodynamic properties are needed.

Partial pressures of species in equilibrium with K_2O at an O_2 pressure of 10^{-15} bar, for which K_2O is the stable solid phase above about 600 K, are shown in Fig. 9. Calculated partial pressures of the major species in equilibrium with condensed K_2O_2 are shown in Fig. 10 for congruent vaporization below the melting temperature; thermodynamic data indicate that K_2O_2 is the stable solid phase under these conditions. Byker *et al.*⁵⁷ concluded from the available evidence that under neutral conditions solid K_2O vaporizes to potassium-rich vapor with a resulting shift in overall condensed phase composition toward K_2O_2 , and that the congruently

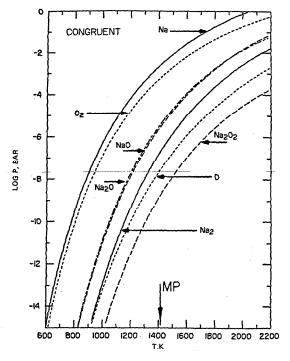


Fig. 6. Na₂O congruent vaporization.

Fig. 8. Sodium oxide maximum vaporization rates. A—Na₂O in 10^{−15} bar O₂; B—Na₂O congruent vaporization; C—Na₂O in 0.2 bar O₂ above 1150 K; D—Na₂O₂ in 0.2 bar O₂ below 1150 K. Rate units—g cm^{−2} s^{−1}.

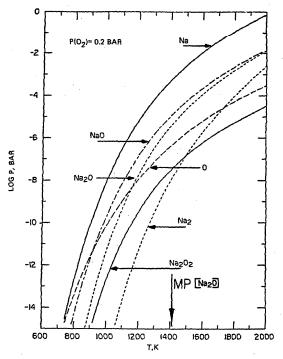


Fig. 7. Vaporization of sodium oxides in 0.2 bar $\rm O_2$. The condensed phase is Na $_2\rm O_2$ below and Na $_2\rm O$ above 1150 K.

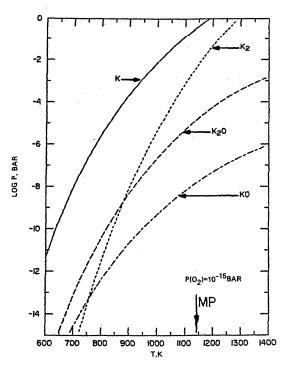


Fig. 9. K_2O vaporization in 10^{-15} bar O_2 .

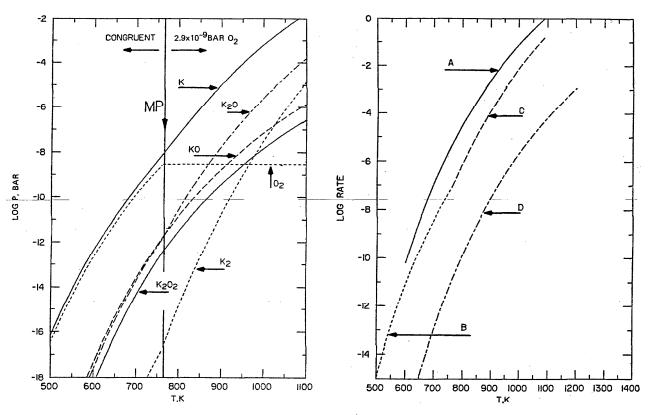


Fig. 10. Congruent vaporization of solid K_2O_2 and vaporization of liquid K_2O_2 in 2.9×10^{-9} bar O_2 .

Fig. 12. Potassium oxide maximum vaporization rates. A— K_2O in 10^{-15} bar O_2 ; B—Congruent vaporization of solid K_2O_2 ; C—Liquid K_2O_2 in 2.9×10^{-9} bar O_2 ; D— KO_2 in 0.2 bar O_2 . Rate units—g cm⁻² s⁻¹.

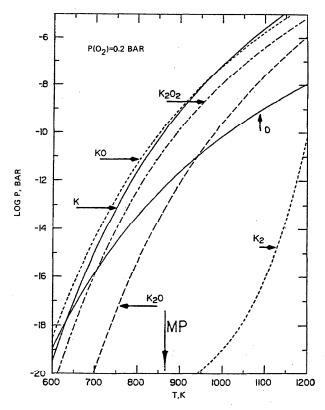


Fig. 11. KO₂ vaporization in 0.2 bar O₂.

vaporizing liquid composition is slightly more potassium-rich than K_2O_2 . Since congruent vaporization of K_2O_2 liquid is thus physically unrealistic, the partial pressures of species in equilibrium with liquid K_2O_2 shown in Fig. 10 were calculated using the O_2 pressure found for the congruently vaporizing solid at its melting temperature. Since the liquid composition changes with time under these conditions, the calculated partial pressures apply only for initial vaporization conditions. These comments also apply to the vaporizations of liquid Rb_2O_2 and Cs_2O_2 under neutral conditions discussed below. The partial pressures of species in equilibrium with condensed KO_2 in 0.2 bar of O_2 are shown in Fig. 11, and maximum mass loss rates for potassium oxides under various conditions are shown in Fig. 12

d. Rb-O System

Partial pressures of species in equilibrium with Rb₂O at an O₂ pressure of 10⁻¹⁵ bar, for which Rb₂O is the stable solid phase above 500 K, are shown in Fig. 13. The partial pressures of species calculated for the congruent vaporization of solid Rb₂O₂ are shown in Fig. 14. Since the congruently vaporizing liquid composition is not known, the partial pressures over liquid Rb₂O₂ shown in this figure were calculated for initial vaporization conditions at the oxygen potential of the congruently vaporizing solid at the melting temperature. Figure 15 shows the calculated partial pressure of species in equilibrium with RbO₂ at an oxygen partial pres-

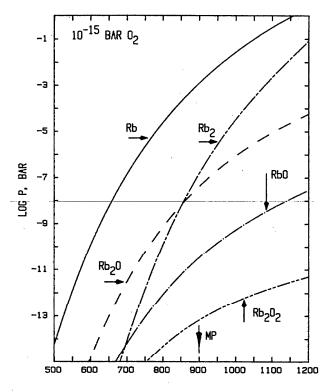


Fig. 13. Rb₂O vaporization in 10⁻¹⁵ bar O₂.

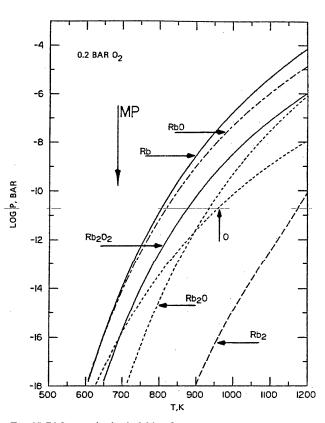


FIG. 15. RbO₂ vaporization in 0.2 bar O₂.

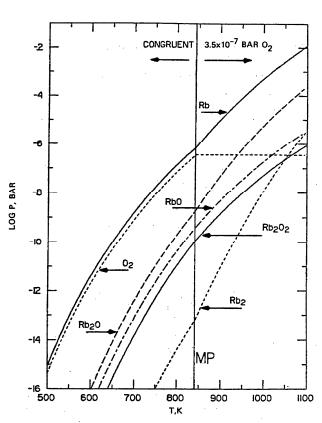


FIG. 14. Congruent vaporization of solid Rb_2O_2 and vaporization of liquid Rb_2O_2 in 3.5×10^{-7} bar O_2 .

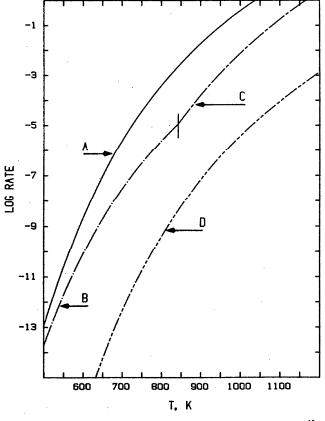
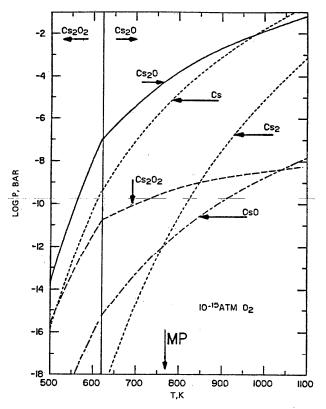


FIG. 16. Rubidium oxide maximum vaporization rates. A—Rb₂O in 10^{-15} bar O₂; B—Congruent vaporization of solid Rb₂O₂; C—Rb₂O₂ in 3.5×10^{-7} bar O₂; RbO₂ in 0.2 bar O₂. Rate units—g cm⁻² s⁻¹.



0.2 BAR O₂ -2 -6 Cs₂0₂ -8 Cs₂0 -10 -12 MP -14 600 500 700 800 900 1000 1100 1200

_Fig. 17. Vaporization of cesium oxides in 10^{-15} bar O_2 . The condensed phases are Cs_2O_2 below 620 K and Cs_2O above this temperature.

-Fig. 19. CsO₂-vaporization in 0.2 bar O₂-

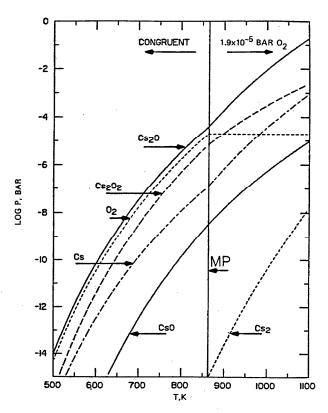


FIG. 18. Congruent vaporization of solid $\rm Cs_2O_2$ and vaporization of liquid $\rm Cs_2O_2$ in $\rm 1.9\times10^{-5}$ bar $\rm O_2$.

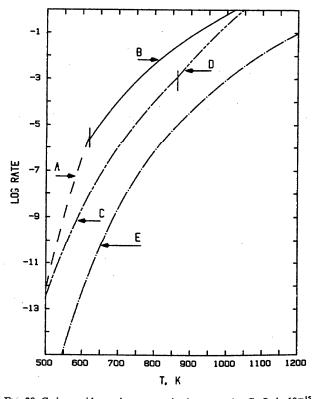


FIG. 20. Cesium oxide maximum vaporization rates. A—Cs₂O₂ in 10^{-15} bar O₂; B—Cs₂O in 10^{-15} bar O₂; C—Congruent vaporization of solid Cs₂O₂; D—Vaporization of liquid Cs₂O₂ in 1.9×10^{-5} bar O₂; E—CsO₂ in 0.2 bar O₂. Rate units—g cm⁻² s⁻¹.

sure of 0.2 bar. Calculated maximum mass loss rates for these vaporization conditions are shown in Fig. 16.

e. Cs-O System

At an oxygen pressure of 10^{-15} bar, Cs_2O_2 is the stable solid phase below about 620 K, and Cs_2O is stable from this temperature to its melting point. Figure 17 shows the partial pressures of species in equilibrium with Cs_2O under these conditions. The partial pressures of species in equilibrium with congruently vaporizing solid Cs_2O_2 are shown in Fig. 18. Since the congruently vaporizing liquid composition is not known, partial pressures of species in equilibrium with liquid Cs_2O_2 shown in this figure were calculated for initial vaporization conditions at the oxygen potential of the congruently vaporizing solid at its melting temperature. Figure 19 shows the partial pressures calculated for vaporization of CsO_2 in 0.2 bar of O_2 . Calculated maximum mass loss rates for the above conditions are shown in Fig. 20.

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