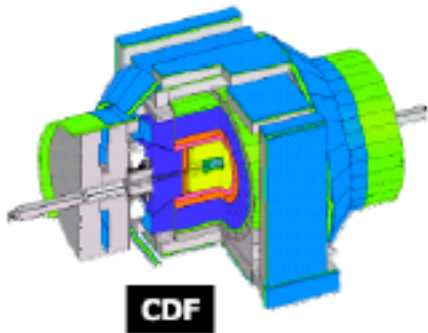
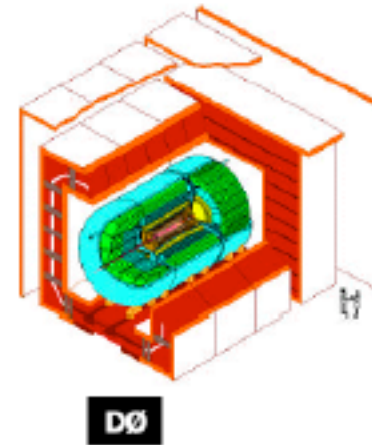
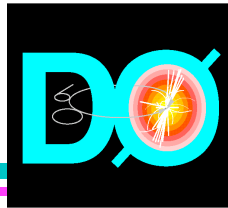

Diboson Physics at the Tevatron



Hadron Collider Physics
July 5, 2005



Al Goshaw , Duke University
for the CDF and D0 Collaborations



The DZero Collaboration

- 19 Countries
- 86 institutions
- ~620 physicists

The D0 Collaboration

 AZ U. of Arizona CA U. of California, Berkeley U. of California, Riverside Cal. State U., Fresno Lawrence Berkeley Nat. Lab. FL Florida State U. IL Fermilab U. of Illinois, Chicago Northern Illinois U. Northwestern U. IN Indiana U. IA U. of Notre Dame IA Iowa State U. KS U. of Kansas Kansas State U. LA Louisiana Tech U. MD U. of Maryland MA Boston U. NE U. of Nebraska Michigan State U. MS U. of Mississippi NE U. of Nebraska NJ Princeton U. NY Columbia U. U. of Rochester SUNY, Stony Brook Brookhaven Nat. Lab. OK Langston U. U. of Oklahoma RI Brown U. TX Southern Methodist U. U. of Texas at Arlington VA U. of Virginia WA U. of Washington	 U. de Buenos Aires	 LAFEX, CBPF, Rio de Janeiro State U. do Rio de Janeiro State U. Paulista, São Paulo	 U. of Alberta McGill U. Simon Fraser U. York U.	 IHEP, Beijing
 U. de los Andes, Bogotá	 Charles U., Prague Czech Tech U., Prague Academy of Sciences, Prague	 LPC, Clermont-Ferrand IN, IN2P3, Grenoble CPPM, IN2P3, Marseille LAL, IN2P3, Orsay LPHE, IN2P3, Paris DAPNIA/SPP, CEA, Saclay IHEP, Strasbourg IPN, IN2P3, Villeurbanne	 U. San Francisco de Quito	 U. of Aachen Bonn U. U. of Freiburg U. of Mainz Ludwig-Maximilians U., Munich U. of Wuppertal
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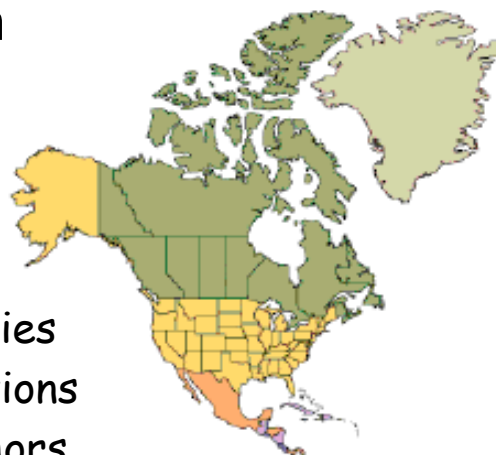
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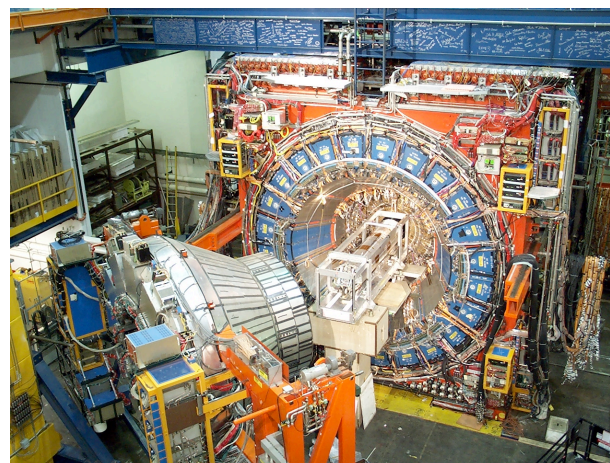
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- 12 countries
- 62 institutions
- ~700 authors



Outline

- Introduction
- Survey of recent measurements
 - $W(l\nu)\gamma$
 - WZ studies using leptonic decays
 - WW and WZ studies using leptonic and hadronic decays
 - $Z(l l)\gamma$
- Summary and outlook

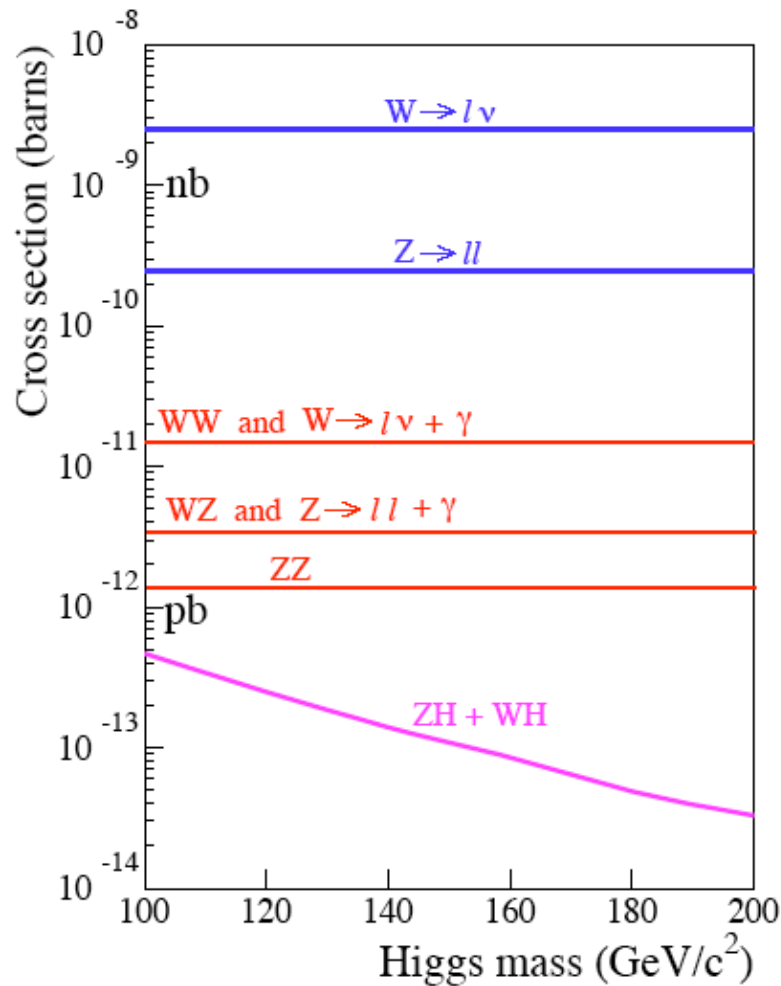
Introduction

Boson production at the Tevatron

- The CDF and DØ collaborations are completing the first round of Run II studies of γ , W , Z and H production.
- Inclusive W and Z bosons
 - W/Z production and decay properties (Serban Protopopescu's talk)
 - W mass and width (Mark Lancaster's talk)
- Vector boson pair production (**this talk**)

$\gamma\gamma$	$W W$
$W\gamma$	$W Z$
$Z\gamma$	$Z Z$
- Higgs boson searches (Anna Goussiou talk)
 - $WH, ZH, H \rightarrow W^*W, \text{SUSY Higgs}$

$p \bar{p}$ production cross sections of W and Z bosons at $\sqrt{s} = 1.96$ TeV



High statistics W/Z inclusive

Lower statistics di-bosons

$E_{T(\gamma)} > 10$ GeV $\Delta R(l\gamma) > 0.7$

Limits on H production

Di-boson physics at the Tevatron

- The study of di-boson production provides a rich source of electroweak Standard Model tests, is sensitive to new physics signatures, and opens a window into the challenges faced in searches for the Higgs boson.
- The CDF and DØ experiments have completed the first analysis phase
 - Results reported here are based upon 200-400 pb⁻¹ of data.
 - Goal is to update with ~ 1 fb⁻¹ of data by winter conferences 2006
 - Ultimate sensitivity based upon 4-8 fb⁻¹
 - And then continuation at the LHC ...
- There are separate talks on electron, photon and muon identification at CDF and DØ, and I will not dwell on details here
 - Electron/photon ID at the Tevatron (Greg Veramendi's talk)
 - Muon ID at the Tevatron (Jeff Temple's talk)

Approach to di-boson studies

- 1. Compare di-boson ($W\gamma$, $Z\gamma$, WW , WZ , ZZ) production properties to Standard Model predictions and measure agreement/deviations.
- 2. Use anomalous coupling parameters as the metric for evaluating the sensitivity to new physics. This assumes the new physics appears as deviations of the W and Z boson from Standard Model point particles. There are of course other sources of new physics that would appear in di-boson production -- perhaps the most likely sources of a discovery.
- 3. Use the advantage of having both $q\bar{q}$ and $q\bar{q}'$ collisions to separate out specific triple gauge couplings where possible:
 - $q\bar{q}' \rightarrow W^* \rightarrow W\gamma$ $WW\gamma$ coupling only
 - $q\bar{q}' \rightarrow W^* \rightarrow WZ$ WWZ coupling only
 - $q\bar{q} \rightarrow Z/\gamma \rightarrow WW$ mix of $WW\gamma$ and WWZ couplings
 - $q\bar{q} \rightarrow Z/\gamma \rightarrow Z\gamma$ mix of $ZZ\gamma$ and $Z\gamma\gamma$ couplings
 - $q\bar{q} \rightarrow Z/\gamma \rightarrow ZZ$ mix of $ZZ\gamma$ and ZZZ couplings

Approach to di-boson studies

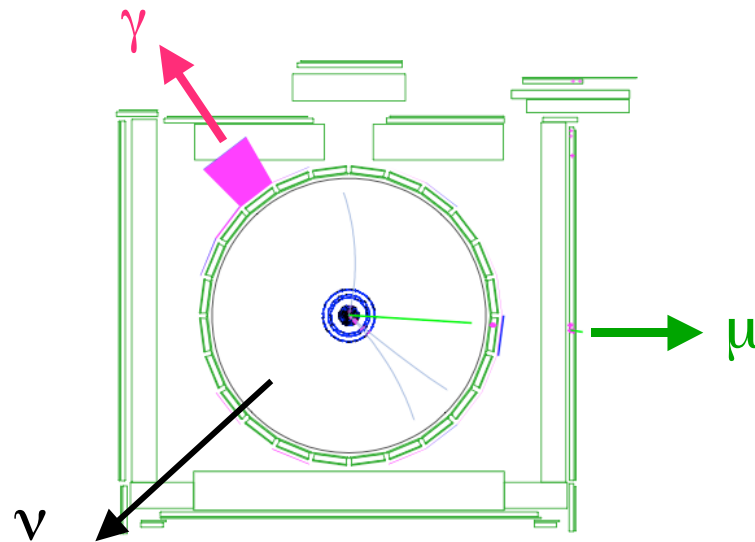
- 4. For the triple gauge coupling studies use (primarily) leptonic decays of the W and Z :
 - $W \gamma \rightarrow l \nu \gamma$
 - $Z \gamma \rightarrow l^+ l^- \gamma$
 - $W^+ W^- \rightarrow l^+ \nu l^- \nu$
 - $W Z \rightarrow l' \nu l^+ l^-$
 - $Z Z \rightarrow l^+ l^- l'^+ l'^-$ and $l^+ l^- \nu \nu$
 - where $l = e$ or μ
- 5. Extend measurements to W/Z hadronic decay channels
 - Specific channels: $W/Z(\text{jet-jet}) + \gamma$ and $W/Z(\text{jet-jet}) + W(l\nu)$
 - Useful for calibration/improvement of di-jet mass resolution
 - Similar to searches for Higgs boson in $W/Z H(b\bar{b})$ searches

W_γ studies

using $p \bar{p} \rightarrow l \nu \gamma + X$

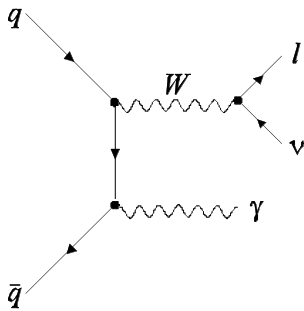
WW_γ coupling

CDF
 $p \bar{p} \rightarrow \mu \nu \gamma + X$
candidate

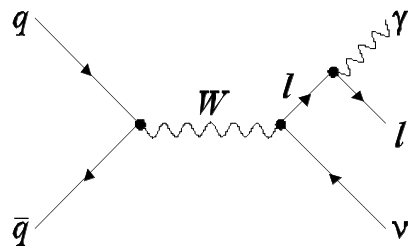


$p \bar{p} \rightarrow l \nu \gamma + X$ Production

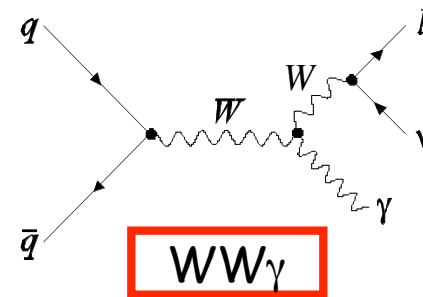
- The $l \nu \gamma$ final states have contributions from quark and lepton bremsstrahlung processes and the direct production $W\gamma \rightarrow l \nu \gamma$



initial state radiation



final state radiation



triple gauge coupling

- The first two diagrams involve W boson coupling only to fermions, and are assumed to be described by the Standard Model.
- The third diagram depends on the $WW\gamma$ coupling
- Therefore the production $p \bar{p} \rightarrow l \nu \gamma + X$ is a measure of this coupling

$p \bar{p} \rightarrow l \nu \gamma + X$ Production

- Under the assumption of Lorentz and electromagnetic gauge invariance, for massless fermions, the $WW\gamma$ coupling can be described in terms of four parameters.
- The effective Lagrangian is [Baur and Berger PRD 41, 1476 (1990)]

$$\begin{aligned} \mathcal{L}_{WW\gamma} = & -ie [(W_{\mu\nu}^\dagger W^\mu A^\nu - W_\mu^\dagger A_\nu W^{\mu\nu}) \\ & + \kappa_\gamma W_\mu^\dagger W_\nu F^{\mu\nu} + \lambda_\gamma / M_W^2 W_{\lambda\mu}^\dagger W_\nu^\mu F^{\nu\lambda} \\ & + 2 \text{ more CP violating terms} \end{aligned}$$

Strong constraints from limits on the neutron's electric dipole moment

- The magnetic dipole and electric quadrupole moments of the W boson are given by:

$$\mu_W = (1 + \kappa_\gamma + \lambda_\gamma)e/2M_W \text{ and } Q_W = -(\kappa_\gamma - \lambda_\gamma)e/M_W^2$$

- In the SM at tree level $\Delta\kappa_\gamma = \kappa_\gamma - 1 = \lambda_\gamma = 0$. Estimates of loop corrections are small: $|\Delta\kappa_\gamma| = 0.008$ and $|\lambda_\gamma| = 0.002$.

$p \bar{p} \rightarrow l \nu \gamma + X$ Production

- Destructive interference of the TGC diagram with the initial state bremsstrahlung process suppresses the $l \nu \gamma$ cross section. For $p \bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV the SM expectations are:

inclusive W and Z boson production:

$$\sigma[W(l\nu)] / \sigma[Z(ll)] \sim 10.7$$

while for $E_T(\gamma) > 10$ (100) GeV, $\Delta R(l\gamma) > 0.7$

$$\sigma[l\nu\gamma] / \sigma[lh\gamma] \sim 4.3 \text{ (1.5)}$$

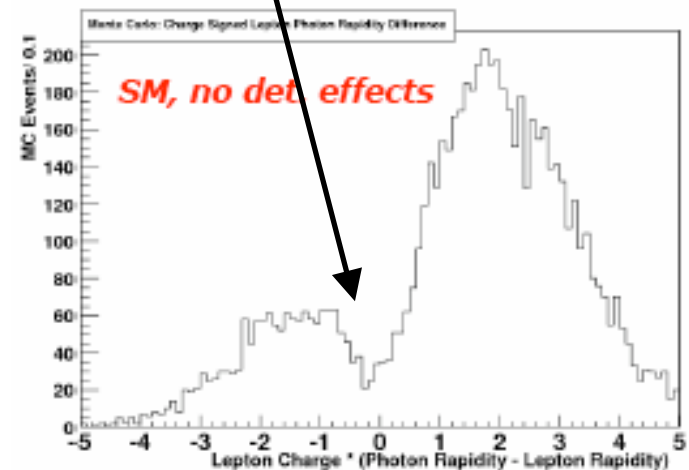
- Radiation amplitude zeros also occur when $\eta(\gamma) - \eta(l^\pm) \sim -0.3$.
[see Baur, Errede, + Landsberg, PRD 50, 1917 (1994)]

Deviations from these SM predictions can be expressed in terms of:

$$\Delta K_\gamma = \Delta K_{0\gamma} / (1 + s/\Lambda^2)^2$$

$$\lambda_\gamma = \lambda_{0\gamma} / (1 + s/\Lambda^2)^2$$

where Λ is the scale of the new physics and \sqrt{s} is the W_γ invariant mass.



Data selection for $l \nu \gamma$ events

- Events triggered on high E_T/P_T central electron/muon
- Selection of leptons** similar to inclusive W/Z measurements

	charged leptons	neutrinos	$l\nu$ transverse mass GeV/c^2
electron channels	$E_T > 25 \text{ GeV}$	$\cancel{E}_T > 25 \text{ GeV}$	30 - 120 (CDF) 40 -> (DØ)
muon channels	$P_T > 20 \text{ GeV}/c$	$\cancel{E}_T > 20 \text{ GeV}$	30 - 120 (CDF) None (DØ)

- Selection of photons**

- central: $|\eta| < \sim 1.0$ photon
- energy: $E_T > 7 - 8 \text{ GeV}$
- isolated: $\Delta R(l \gamma) > 0.7$

Integrated
Luminosity
130-200 pb^{-1}

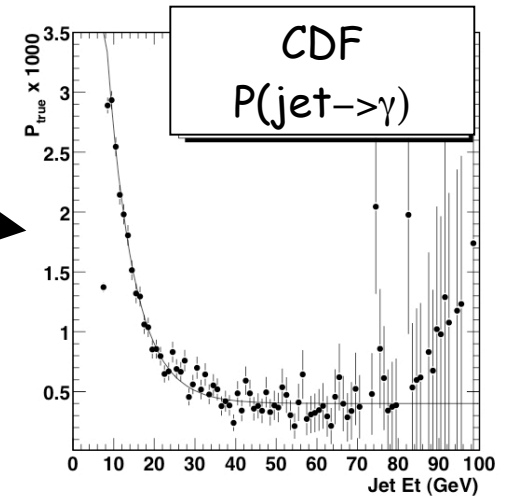
Backgrounds in $l \nu \gamma$ events

- Backgrounds are dominated by $W(l \nu)+\text{jets}$ and $Z/\gamma(l l)+\text{jets}$ with jet \rightarrow fake photon
 - Use jet-triggered data samples to measure the fake rates
 - Correct for real photon content ($\gamma+\text{jet}$ events) in jet data
 - Apply this to jets in $W/Z+\text{jet}$ data

$$P(\text{jet} \rightarrow \text{fake } \gamma) \sim 3 \times 10^{-3} (E_T \sim 10 \text{ GeV})$$

$$\sim 4 \times 10^{-4} (E_T > 50 \text{ GeV})$$

- Other backgrounds determined from SM generators
 - Cross-talk between $Z(l l)\gamma$ and $W(l \nu)\gamma$ channels
 - Feed-down from $W \rightarrow \tau \nu$ decays



backg. source	$e \nu \gamma$	$\mu \nu \gamma$
W + jets	59.5 \pm 18.1	27.6 \pm 7.5
$\tau \nu \gamma$	1.5 \pm 0.2	2.3 \pm 0.2
$ll \gamma$	6.3 \pm 0.3	17.4 \pm 1.0
total backg.	67.3 \pm 18.1	47.3 \pm 7.6
data	195	128

CDF Data

S/B \sim 1.8

Signal acceptance (A) and efficiency (ϵ) for $l \nu \gamma$

- Electron and muon ID efficiencies are evaluated from data using $Z \rightarrow ee$ and $\mu\mu$ decays.
- Photon ID efficiencies are determined from a combination of data (use electrons as proxies for photons) and GEANT-based detector simulations.
- Geometric acceptances determined from SM event generators and detector simulations.
 - Correct for $W \rightarrow l \nu$ decay phase space
 - Quote cross sections for $E_T(\gamma) > 8 \text{ GeV}$ and $\Delta R(l \gamma) > 0.7$

	$e \nu \gamma$	$\mu \nu \gamma$
background	$60.8_{\pm 4.1}$	$71.3_{\pm 5.2}$
data	112	161
$A \times \epsilon$	$0.023_{\pm 0.001}$	$0.044_{\pm 0.002}$

DØ Data

$A \times \epsilon$
 $\sim 2 - 4\%$

Comparison of $p \bar{p} \rightarrow l \nu \gamma + X$ to Standard Model predictions

$\sigma(p \bar{p} \rightarrow l \nu \gamma + X)$ at $\sqrt{s} = 1.96$ TeV with $\Delta R(l \gamma) > 0.7$

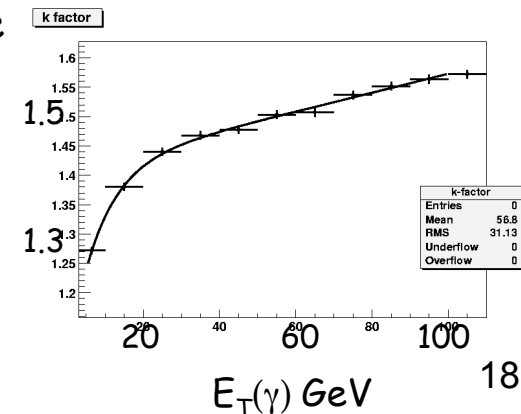
	N_{data} (e + μ)	$\sigma(l \nu \gamma)$ exp.	$\sigma(l \nu \gamma)$ SM theory	$E_T(\gamma)$ cut
CDF	323	18.1 \pm 3.1	19.3 \pm 1.4	> 7 GeV
DØ	273	14.8 \pm 2.1	16.0 \pm 0.4	> 8 GeV

Measured cross sections:

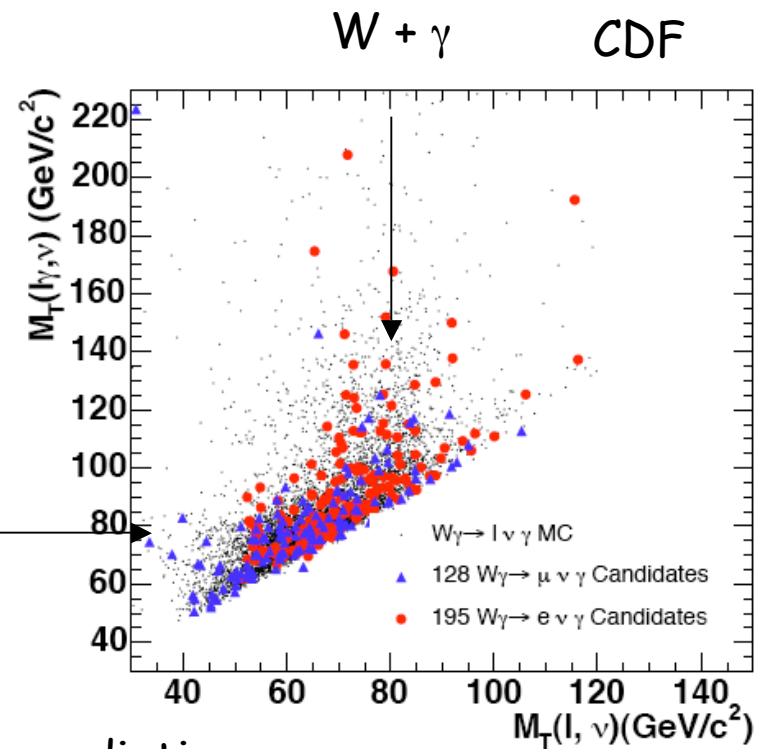
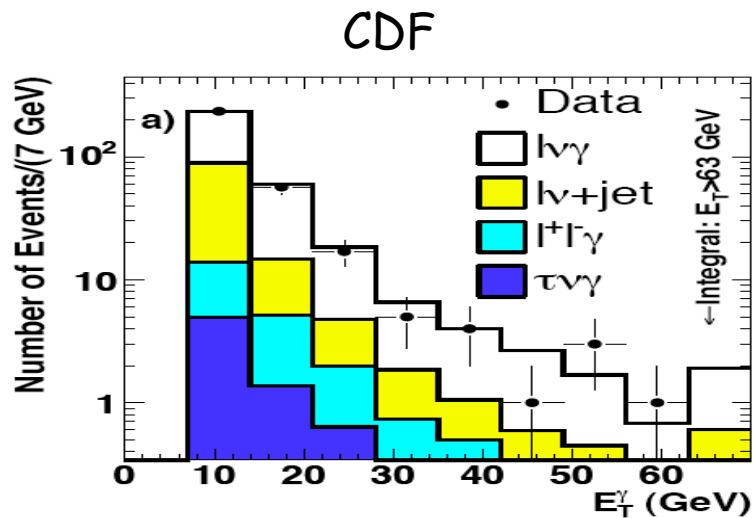
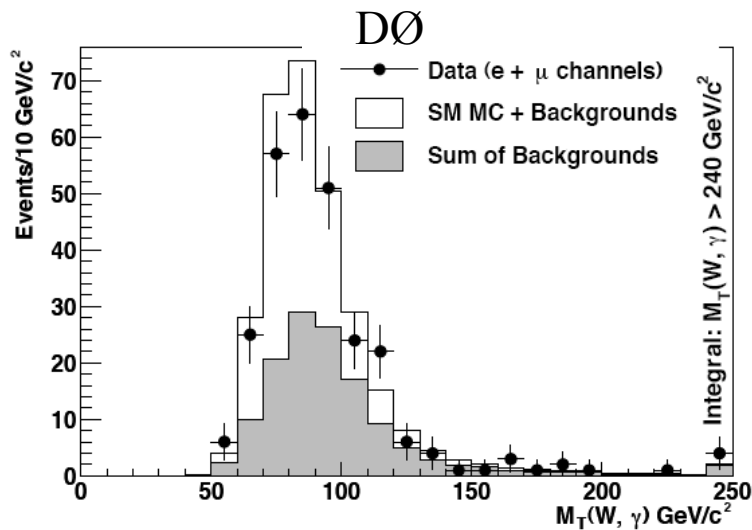
- Systematic and statistical errors ~ equal plus 6% luminosity error
- Systematic errors dominated by jet \rightarrow photon fake rate

Standard Model cross section predictions:

- Use LO matrix elements
[Baur + Berger, PRD 47, 4889 (1993)]
- Correct for NLO k-factor
[Baur, Han, +Ohnemus PRD57, 2823 (1998)]



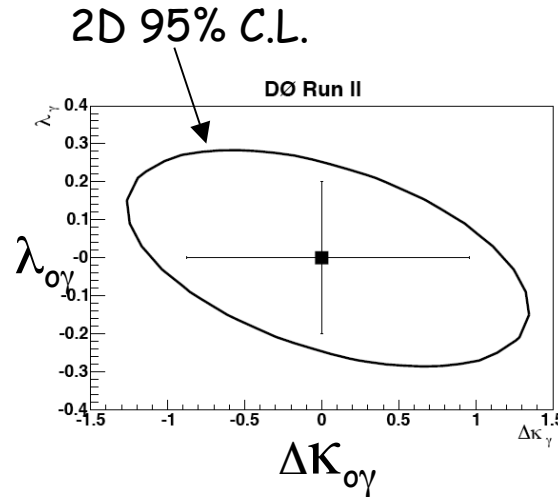
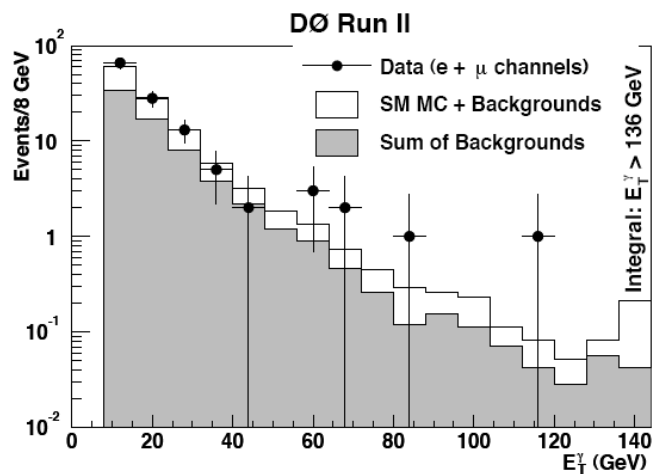
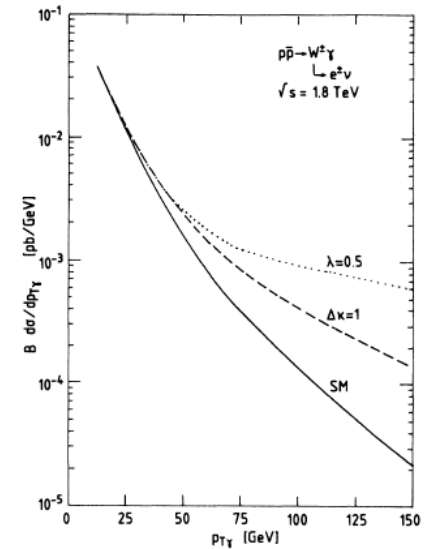
Comparison of $l \nu \gamma$ signal to Standard Model predictions



PRL 94, 040803 (2005) CDF

Using $l \nu \gamma$ events to put limits on $WW\gamma$ couplings (DØ)

- Non-zero $\Delta\kappa_\gamma$ or λ_γ lead to enhancement of high E_T photons above the SM prediction.
- Suppress event with FSR by selecting events with $M_T(l \nu \gamma) > 90 \text{ GeV}/c^2$
- Use binned-likelihood fitting $E_T(\gamma)$ on $\Delta\kappa_{o\gamma}$ vs $\lambda_{o\gamma}$ grid



1D 95% C.L.

$$-0.88 < \Delta\kappa_{o\gamma} < 0.96$$

$$-0.20 < \lambda_{o\gamma} < 0.20$$

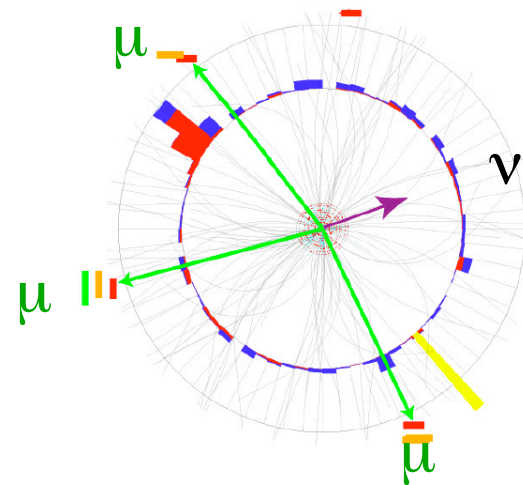
$$\Lambda = 2 \text{ TeV}$$

WZ studies

using $p \bar{p} \rightarrow l' \nu l^+ l^- + X$

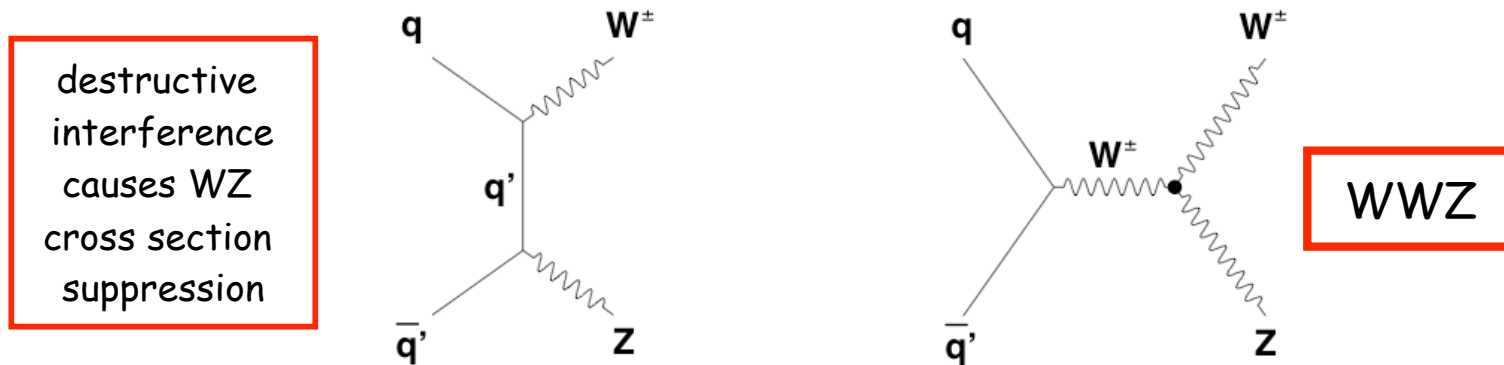
WWZ coupling

DØ
 $p \bar{p} \rightarrow \mu \nu \mu \mu + X$
candidate



$p \bar{p} \rightarrow WZ + X$ Production

- The next logical step in unraveling the di-boson anomalous couplings is to measure WZ production. This isolates the WWZ vertex.



- Under the same assumptions* made for the $WW\gamma$ coupling the effective Lagrangian is [Haigwara, Peccii and Zeppenfeld, NP B282 253 (1987)]:

$$\begin{aligned} \mathcal{L}_{WWZ} = & -ie \cot\theta_W [g_1^Z (W_{\mu\nu} W^\mu Z^\nu - W_\mu Z_\nu W^{\mu\nu}) \\ & + \kappa_Z W_\mu W_\nu Z^{\mu\nu} + \lambda_Z / M_W^2 W_{\lambda\mu} W^\mu_\nu Z^{\nu\lambda} \end{aligned}$$

* in addition to dropping CP violating terms, exclude a term which conserves CP but violates both C and P .

- In the SM at tree level $g_1^Z = \kappa_Z = 1$ and $\lambda_Z = 0$.

Data selection for WZ events

- The SM expectation for $p \bar{p} \rightarrow W Z + x$ at $\sqrt{s} = 1.96$ TeV is ~ 4.0 pb
- But the branching ratios reduce the usable signal:

WZ decay channel	Fraction %
$q \bar{q}' q \bar{q} + q \bar{q} \nu \nu$	61.1
$q \bar{q} l \nu$	22.4
$q \bar{q}' ll$	6.9
$l \nu \nu \nu$	6.4
$l \nu ll$	3.2

← buried in QCD jets

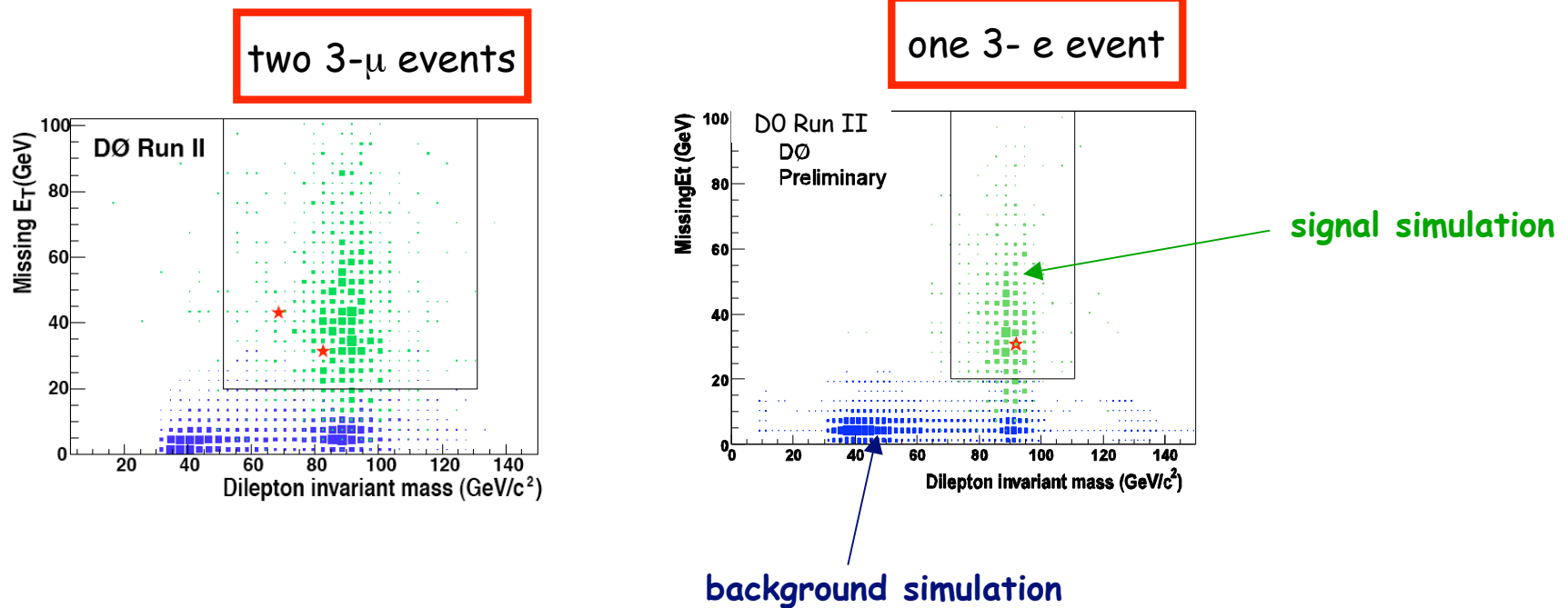
} ← useful (see below)

← buried in inclusive W/Z

← cleanest channel

- Restricting $l = e$ or μ , the total branching ratio for $p \bar{p} \rightarrow l' \nu l^+ l^- + x \sim 1.4\%$.
- Acceptance \times efficiency varying from 10 to 15% (without BR's)
 - Particle ID efficiencies use $Z \rightarrow ee/\mu\mu$ data
 - Acceptance use PYTHIA event generator plus detector simulations

Data selection for $l' \nu l^+ l^-$ events



- Backgrounds with $S/B \sim 3-4$
 - Dominated by jet \rightarrow fake lepton in $W/Z + \text{jet}/\gamma$ events (from data)
 - Feed through from $WW, ZZ, t\bar{t}$ determined from SM event generators plus detector simulations.

Comparison of WZ (and ZZ) to Standard Model predictions using all leptonic decays

$\sigma[p\bar{p} \rightarrow WZ (ZZ) + x]$ at $\sqrt{s} = 1.96$ TeV corrected for W/Z branching ratios

	N_{data}	$N_{\text{background}}$	SM theory cross section (pb)	Experimental cross section (pb)	95% C.L.
CDF(WZ+ZZ)	3	1.02 \pm 0.24	5.0 \pm 0.4	4.3 +5.0 - 2.6	< 15.2 pb
DØ (WZ)	3	0.71 \pm 0.08	3.65 \pm 0.26	4.5 +5.1 - 3.3	< 13.3 pb

- SM theory predictions at NLO with MCFM
- Experimental details can be found at:
 - DØ hep-ex/0504019 (Submitted to PRL)
 - CDF PRD 71, 091105 (2005)

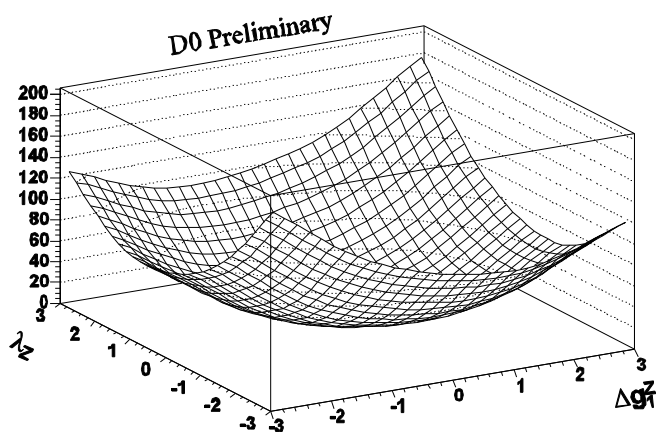
Integrated Luminosity
 ~ 300 pb⁻¹ (DØ)
 ~ 200 pb⁻¹ (CDF)

● Bottom line: consistent with the SM

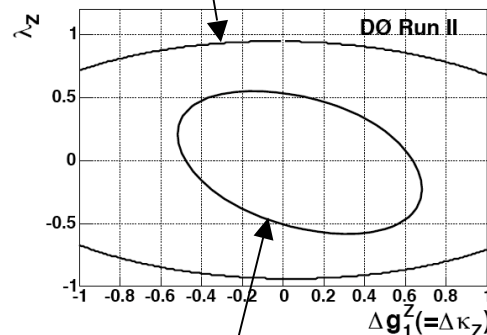
Using $l' \nu l^+ l^-$ events to put limits on WWZ couplings (DØ)

- Non-zero Δg_1^Z , $\Delta \kappa_Z$ or λ_Z lead to enhancement of high Et W/Z above the SM prediction. Here $\Delta g_1^Z = g_1^Z - 1$ and $\Delta \kappa_Z = \kappa_Z - 1$.
- Hold one anomalous coupling parameter at zero and form a 2D grid of predictions versus the other two [Hagiwara, Woodside, Zeppenfeld LO generator].
- Use dipole form factors as for $W\gamma$ analysis:

1D 95% C.L.



Unitarity bound



2D 95% C.L.

$$-0.49 < \Delta g_{01}^Z < 0.66$$

$$-0.48 < \lambda_{0Z} < 0.48$$

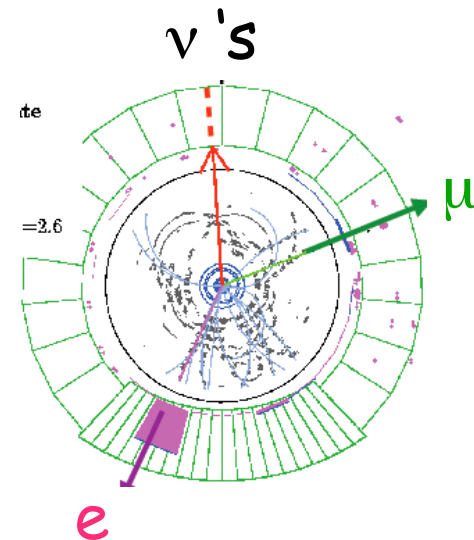
$\Delta \kappa_{0Z}$ no limit

$$\Lambda = 1.5 \text{ TeV}$$

$W^+ W^-$ studies
using $p \bar{p} \rightarrow l^+ \nu l'^- \nu + X$
and $p \bar{p} \rightarrow l \nu q \bar{q} + X$

Mix of WW_γ, WWZ couplings

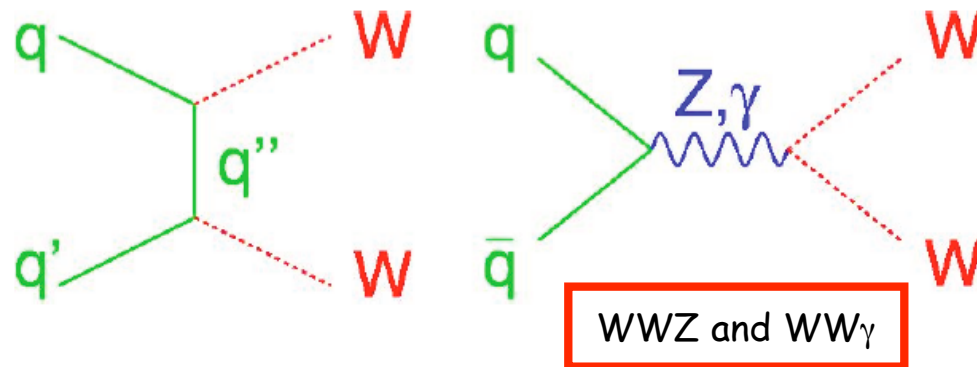
CDF
 $p \bar{p} \rightarrow \mu \nu e \nu + X$
candidate



$p \bar{p} \rightarrow W^+ W^- + X$ Production

- The $WW\gamma$ (WWZ) couplings introduce the parameters discussed above [Haigwara, Woodside and Zeppenfeld, PRD 41, 2113 (1990)]:

destructive interference causes WW cross section suppression



- The SM expectation for $p \bar{p} \rightarrow W W + X$ at $\sqrt{s} = 1.96$ TeV is ~ 12.4 pb

- To date CDF and DØ have used channels.:

- $l \nu l' \nu$ with $l, l' = e$ or μ
BR $\sim 4.6\%$ small BR, good S/B
- $q q' l \nu$ with $l = e$ or μ
BR $\sim 29\%$ good BR, poor S/B

WW	Fraction %
$qq' q q'$	46.2%
$q q' l \nu$	43.5%
$l \nu l' \nu$	10.3%

Comparison of W^+W^- to Standard Model predictions using $l^+ \nu l^- \nu$ decays

- Measurement of $p\bar{p} \rightarrow W^+W^- + X$ at $\sqrt{s} = 1.96$ TeV with BR corrections

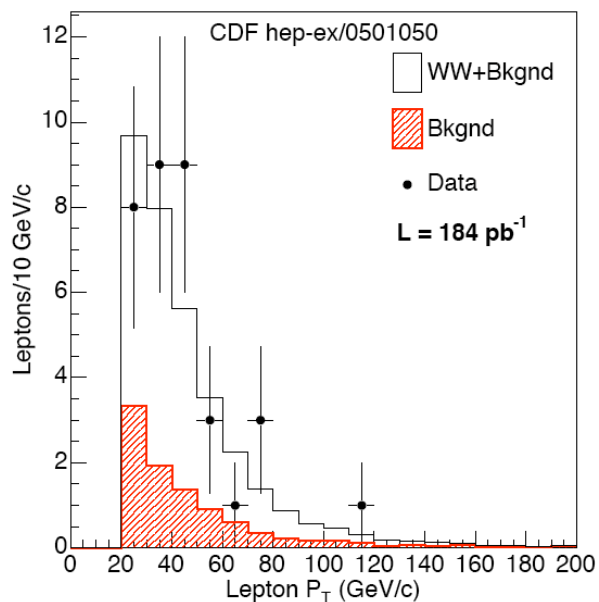
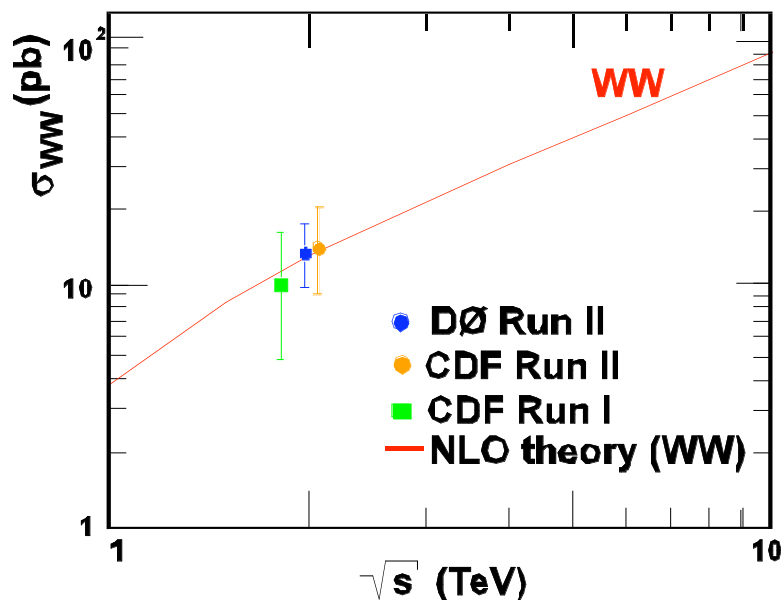
$$\sigma(WW) = 13.8_{-3.8}^{+4.3} (stat.)_{-0.9}^{+1.2} (sys.) \pm 0.9 (lum.) \text{ pb}$$

DØ Run II: PRL 94, 151801 (2005)

$$\sigma(WW) = 14.6_{-5.1}^{+5.8} (stat.)_{-3.0}^{+1.8} (sys.) \pm 0.9 (lum.) \text{ pb}$$

CDF Run II: PRL 94, 211801 (2005)

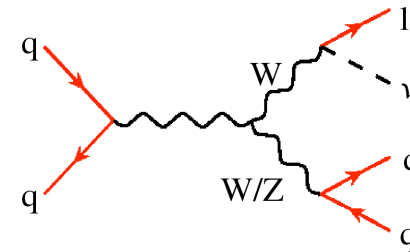
$$SM \sigma(WW) = 12.4 \pm 0.8 \text{ pb}$$



W^+W^- and WZ studies using $l\nu q\bar{q}$ decays

- The channel $p\bar{p} \rightarrow W(l\nu) + W/Z(q\bar{q})$ has been studied by CDF

- Advantages: larger branching ratio
- Disadvantages: much higher backgrounds
- BUT anomalous signals appear at high E_T of the W where backgrounds lower



- $W + \text{jet jet}$ QCD background is constrained by fitting to dijet mass spectrum around M_W plus M_Z peak.

- Fit to data:

$$N_{\text{data}} = 66 \pm 78 \pm 34 \text{ events}$$

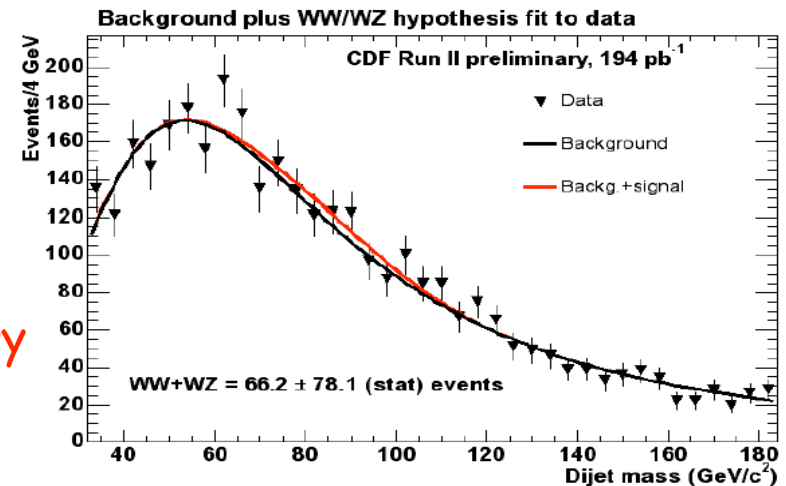
$$\sigma_{\text{data}}(WW + WZ) < 40 \text{ pb}$$

- With the SM expectation

$$N_{\text{SM}} \sim 91 \text{ events}$$

$$\sigma_{\text{SM}}(WW + WZ) = 16.5 \text{ pb}$$

CDF Preliminary



W^+W^- and WZ studies using $l\nu q\bar{q}$ decays

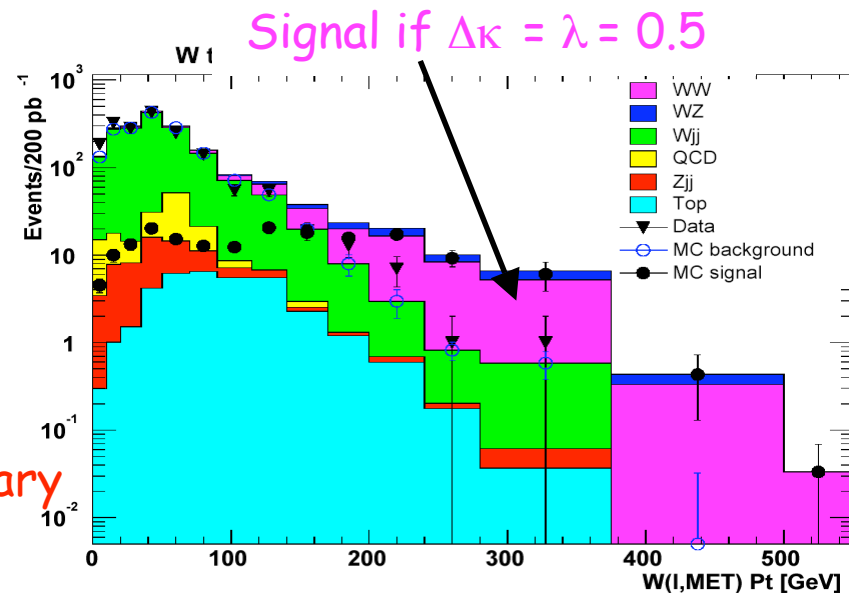
- Fits to anomalous couplings require assumptions here since 5 parameters contribute to WW plus WZ production: Δg_1^Z , $\Delta\kappa_Z$, λ_Z , $\Delta\kappa_\gamma$, and λ_γ
- Assume $\Delta g_{01}^Z = 0$ and let $\Delta\kappa_0 = \Delta\kappa_{0Z} = \Delta\kappa_{0\gamma}$ and $\lambda_0 = \lambda_{0Z} = \lambda_{0\gamma}$
- The P_T of the $W(l\nu)$ is found to be the most sensitive distribution since anomalous VV pairs are produced at high P_T .

- Limits set:

$$-0.42 < \Delta\kappa_0 < +0.58$$

$$-0.32 < \lambda_0 < +0.35$$

CDF Preliminary

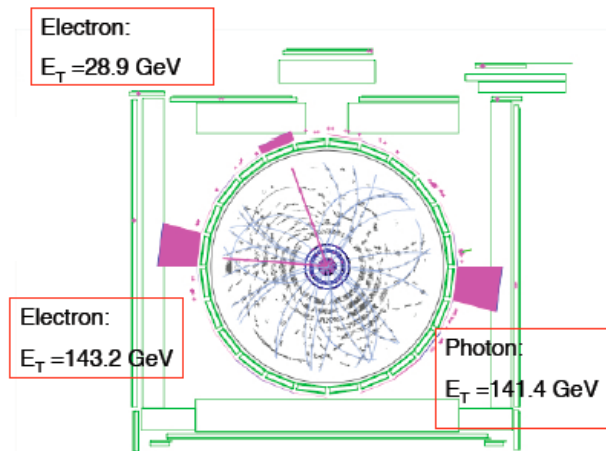


Z_γ studies

using $p \bar{p} \rightarrow l^+ l^- \gamma + X$

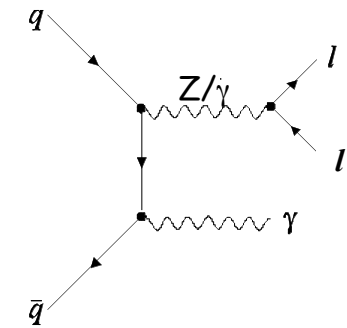
ZZ_γ and $Z\gamma\gamma$ couplings

CDF
 $p \bar{p} \rightarrow e^+ e^- \gamma + X$
candidate

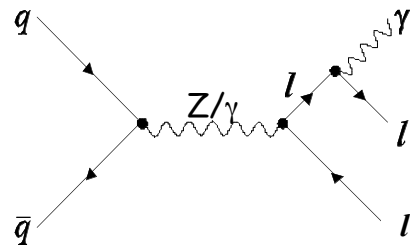


$p \bar{p} \rightarrow l^+ l^- \gamma + X$ Production

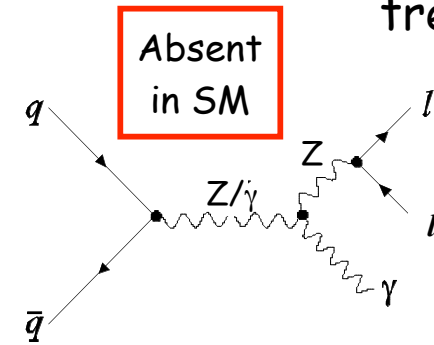
- Within the confines of the SM as the $ZZ\gamma$ and $Z\gamma\gamma$ couplings are zero at tree level.



initial state radiation



final state radiation



triple gauge coupling

Absent
in SM

- Therefore beyond the SM effects appear as deviations from bremsstrahlung radiation.
- Under the assumptions of Lorentz and electromagnetic gauge invariance, the anomalous coupling parameters in this case are:

$$h_i^V = h_{i0}^V / (1 + s/\Lambda^2)^n$$

where $V = s\text{-channel } \gamma \text{ or } Z$

$\sqrt{s} = Z\gamma$ invariant mass

$i = 1,2$ (CP violating), $3,4$ (CP conserving)

\Rightarrow 8 parameters [see e.g. Baur and Berger PRD 47, 4889 (1993)]

Data selection for $l^+ l^- \gamma$ events

- Events triggered on high E_T/P_T central electron/muon. The offline lepton and photon cuts are similar to those for $W\gamma$ events (see page 14).
- DØ selects events with $M(l^+ l^-) > 30 \text{ GeV}/c^2$, CDF 40 - 130 GeV/c^2
- The only significant source of Background is Z + jet events with a jet faking a photon (see p 16).

backg. source	e e γ	$\mu \mu \gamma$
Z + jets	23.6 \pm 2.3	22.4 \pm 3.0
data	138	152

~ 300 pb⁻¹ (DØ)
S/B ~ 6

backg. source	e e γ	$\mu \mu \gamma$
Z + jets	2.8 \pm 0.9	2.1 \pm 0.6
data	36	35

~ 200 pb⁻¹ (CDF)
S/B ~ 15

Comparison of $l^+ l^- \gamma$ signals to Standard Model predictions

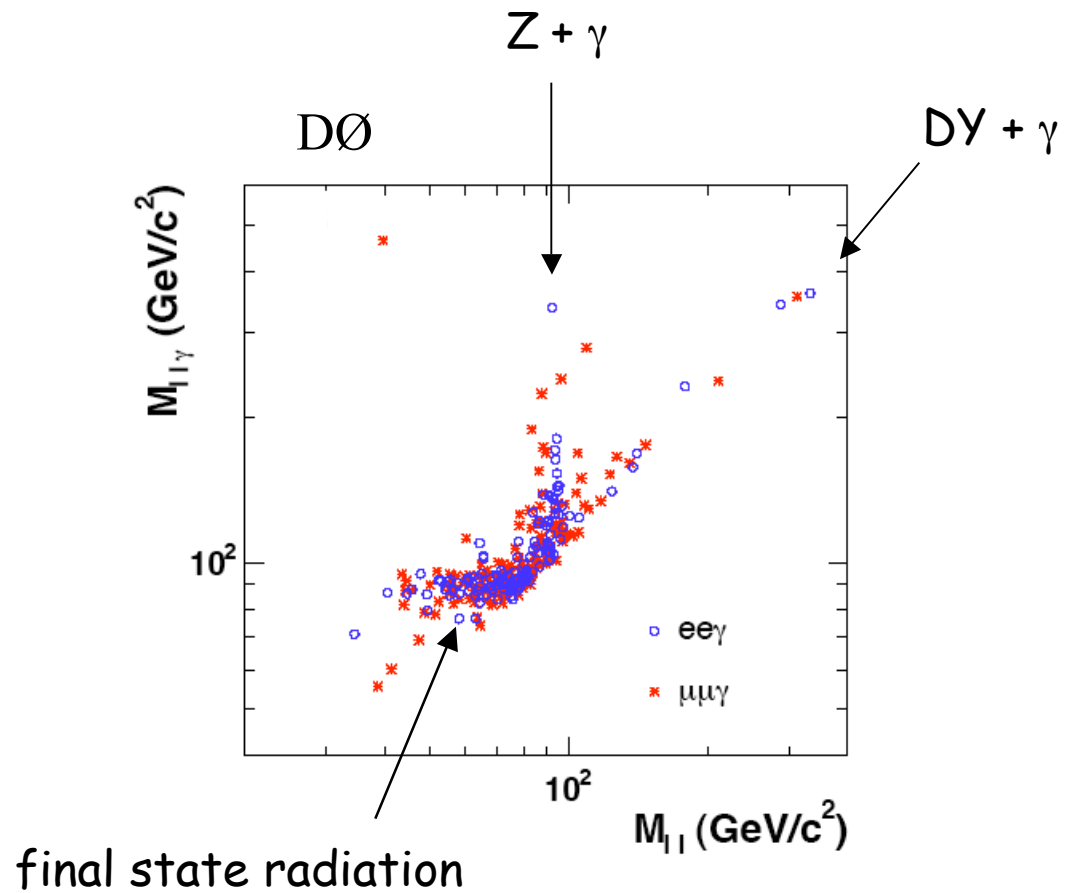
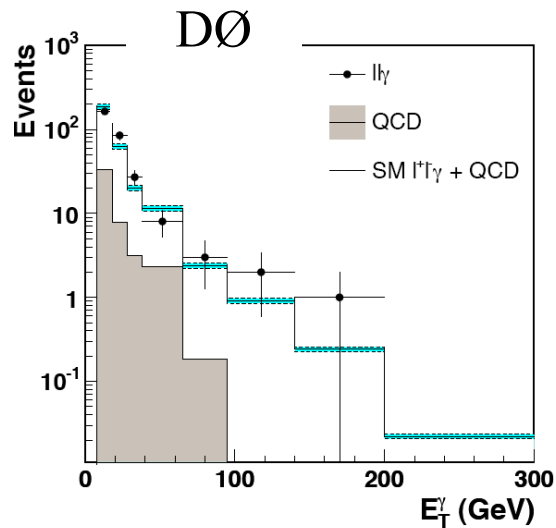
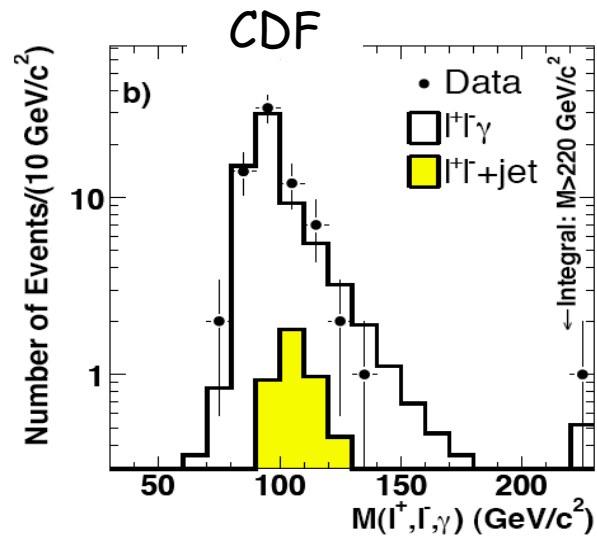
$$\sigma(p \bar{p} \rightarrow ll\gamma + X) \text{ with } \Delta R(l\gamma) > 0.7$$

	N_{data} (e + μ)	$\sigma(ll\gamma)$ exp.	$\sigma(ll\gamma)$ theory	$E_T(\gamma)$ GeV	$M(ll)$ GeV/ c^2
CDF	71	4.6 \pm 0.6	4.5 \pm 0.3	> 7	> 40
DØ	290	4.2 \pm 0.5	3.9 \pm 0.2	> 8	> 30

CDF PRL 94, 041803 (2005)

DØ hep-ex/0502036 (submitted to PRL)

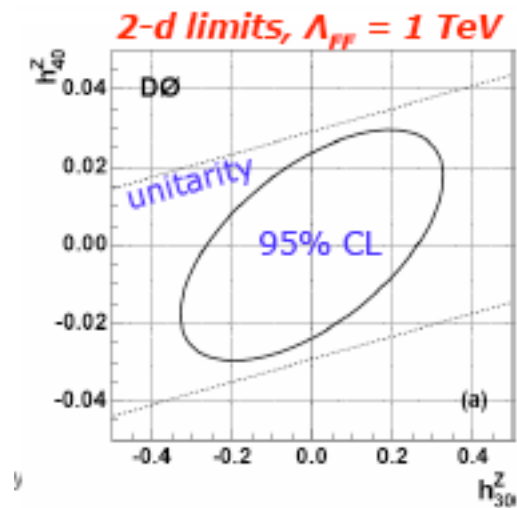
Comparison of $l^+ l^- \gamma$ signal to Standard Model predictions



Using $l^+ l^- \gamma$ events to put limits on $ZZ\gamma$ and $Z\gamma\gamma$ couplings (DØ)

- No deviations from SM predictions, but very clean data samples can be used to put limits on anomalous couplings.
- Form factors imposed to preserve unitarity with $n=3$ for $i=1,3$ and $n=4$ for $i=2,4$ (see form on page 31).
- Use binned-likelihood fitting to $E_T(\gamma)$ on 2D grid of $(h_{10}^V$ vs $h_{20}^V)$ and $(h_{30}^V$ vs $h_{40}^V)$. Set 95% C/L. for $\Lambda = 1$ TeV. (DØ PRL and hep-ex/0502036).

$-0.23 < h_{10,30}^Z < 0.23$
 $-0.020 < h_{20,40}^Z < 0.020$ ←
 $-0.23 < h_{10,30}^\gamma < 0.23$ Best limits to date
 $-0.019 < h_{20,40}^\gamma < 0.019$ ←



Summary and Outlook

SUMMARY: Comparison to SM predictions

Tevatron Run II
 $p\bar{p}$ at $\sqrt{s} = 1.96$ TeV
 200-400 pb⁻¹

All rates and
 kinematic distributions
 are consistent with
 SM predictions

Channel ($l = e, \mu$)	$(\sigma_{\text{data}} - \sigma_{\text{SM}}) / \sigma_{\text{SM}}$
$W\gamma [l\nu\gamma]$	-0.06 ± 0.16 CDF
	-0.06 ± 0.16 DØ
$Z\gamma [ll\gamma]$	$+0.02 \pm 0.13$ CDF
	$+0.08 \pm 0.13$ DØ
$WW [l\nu l\nu]$	$+0.17 \pm 0.42$ CDF
	$+0.10 \pm 0.32$ DØ
cross section limits	$\sigma_{\text{data}} (95\% \text{ C.L.}) / \sigma_{\text{SM}}$
$WZ [l\nu ll]$	3.3 DØ
$WZ + WW [l\nu qq]$	2.4 CDF
$ZW + ZZ [ll (l\nu \text{ or } \nu\nu)]$	3.0 CDF

(table uses nominal SM predictions with no theory uncertainties)

SUMMARY: Limits on anomalous couplings

Analyses just starting on individual channels

Coupling	limits at 95% CL	Energy scale Λ
$WW\gamma$	$-0.88 < \Delta\kappa_{0\gamma} < 0.96$	2 TeV
	$-0.20 < \lambda_{0\gamma} < 0.20$	
WWZ	$-0.49 < \Delta g_{01}^Z < 0.66$	1.5 TeV
	$-0.48 < \lambda_{0Z} < 0.48$	
$ZZ\gamma$	$ h_{10,30}^\gamma < 0.23$	1 TeV
	$ h_{20,40}^\gamma < 0.019$	
ZZZ	$ h_{10,30}^Z < 0.23$	1 TeV
	$ h_{10,30}^Z < 0.020$	
WWZ and $WW\gamma$	$-0.42 < \Delta\kappa_0 < 0.58$	1.5 TeV
	$-0.32 < \lambda_0 < 0.35$	

In future will combine channels and CDF+D0 measurements

SUMMARY: Physics beyond the SM

- New physics sources (anomalous couplings, new fermions or gauge bosons) contribute to the high P_T tails of $W/Z/\gamma$ production.
- At high P_T most sources of background (jets faking photons and leptons) fall rapidly.
- Therefore the sensitivity to new physics is almost entirely statistics driven.
- The Tevatron is ramping up according to its design plan, and the CDF and DØ detectors are operating with good efficiency.
- The data sets presented here represent 2-10% of the potential of the Tevatron.
- We hope to set more than limits ...

Backup Slides

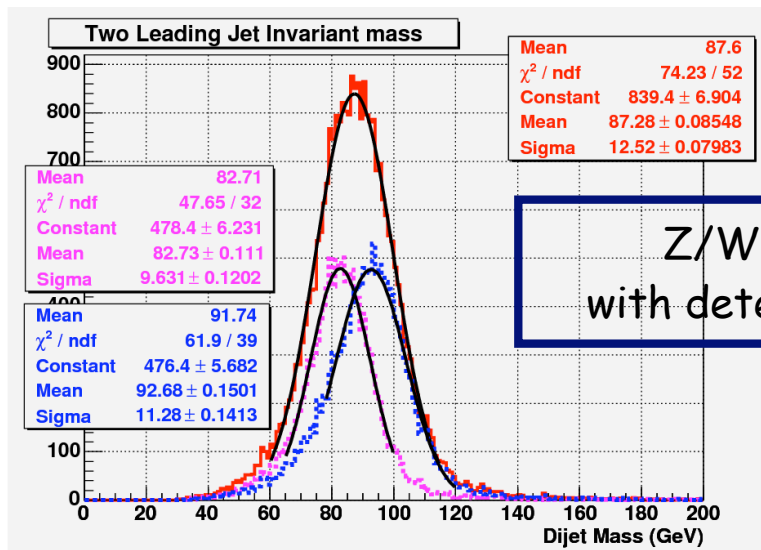
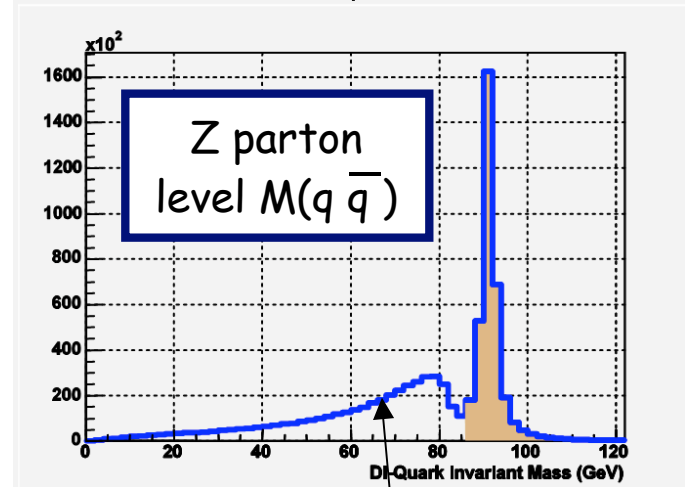
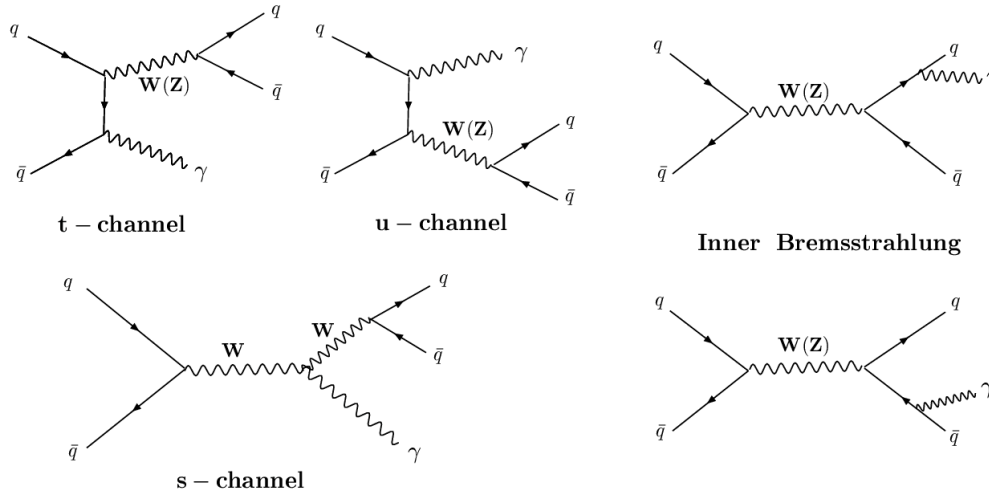
Searches for $W/Z \rightarrow q \bar{q}$
using $W/Z \gamma$ events

Searches for $p \bar{p} \rightarrow W/Z(q \bar{q}) + \gamma$

Motivation

- New physics searches using dijets depend critically on a good understanding of the jet-jet invariant mass resolution.
- One calibration source is $W/Z \rightarrow q \bar{q} \rightarrow \text{jet jet}$
- This requires a trigger that does not bias the W/Z mass peak, and allows low mass side bands for background subtraction.
- Using diboson events of the type $W/Z(q \bar{q}) \gamma$ and $W/Z(q \bar{q}) W(l \nu)$ allows a trigger selection based upon the a high E_t photon or lepton, and provides an unbiased look at the jet-jet spectrum for extraction of $W/Z \rightarrow q \bar{q}$.
- Also, the $W/Z(q \bar{q}) W(l \nu)$ channels are very similar to those used for Higgs searches in $H(b \bar{b}) W(l \nu)$ and provide a SM calibration line.

W/Z(qq̄) signal in W/Z γ events



Z/W di-jet mass with detector resolution

Neural net used for S/B separation

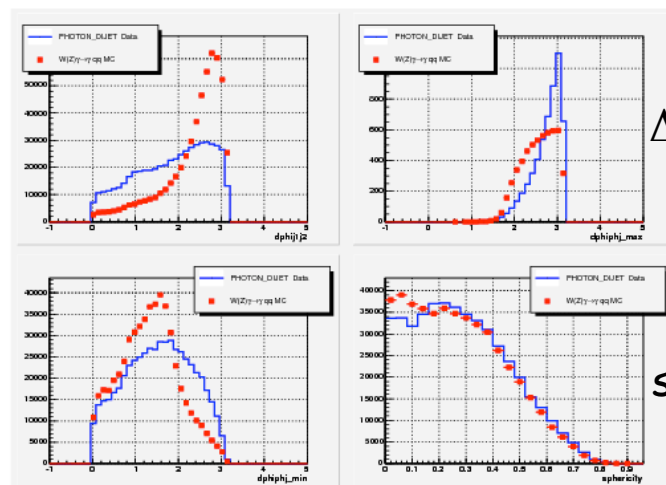
Studied 19
discriminating variable

Found best 10 for
maximizing S/\sqrt{b}

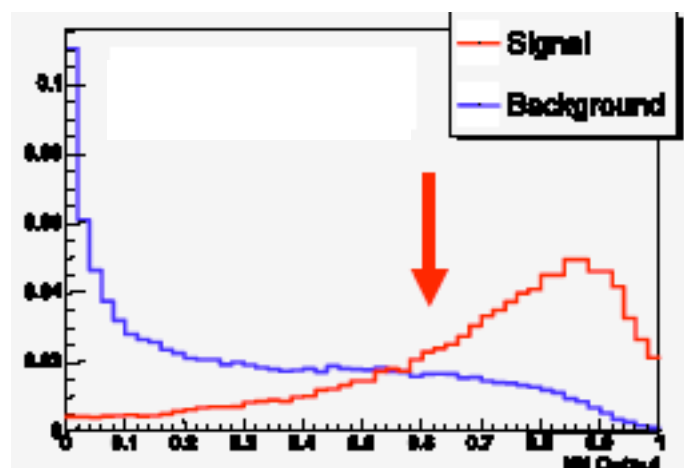
$\Delta M = 72-110 \text{ GeV}$

$S/\sqrt{B} = 2.17$

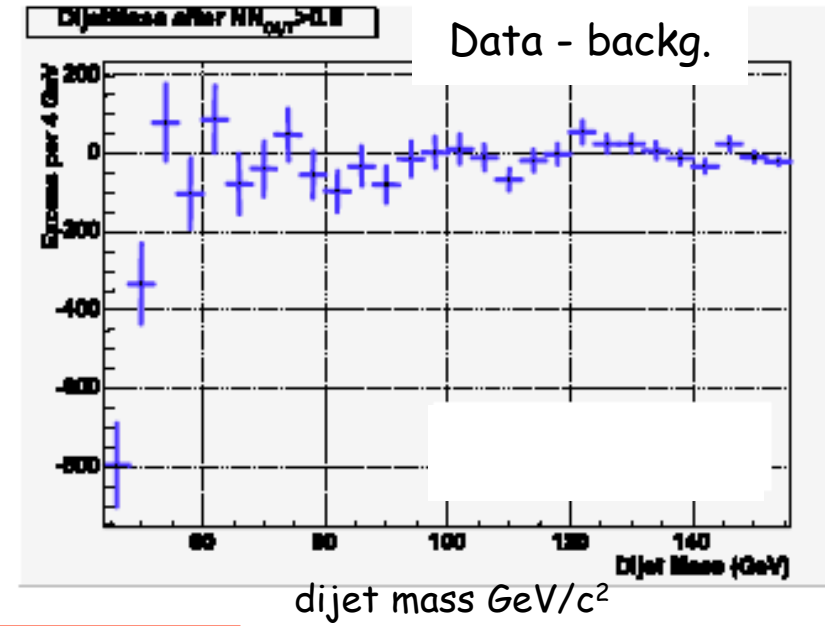
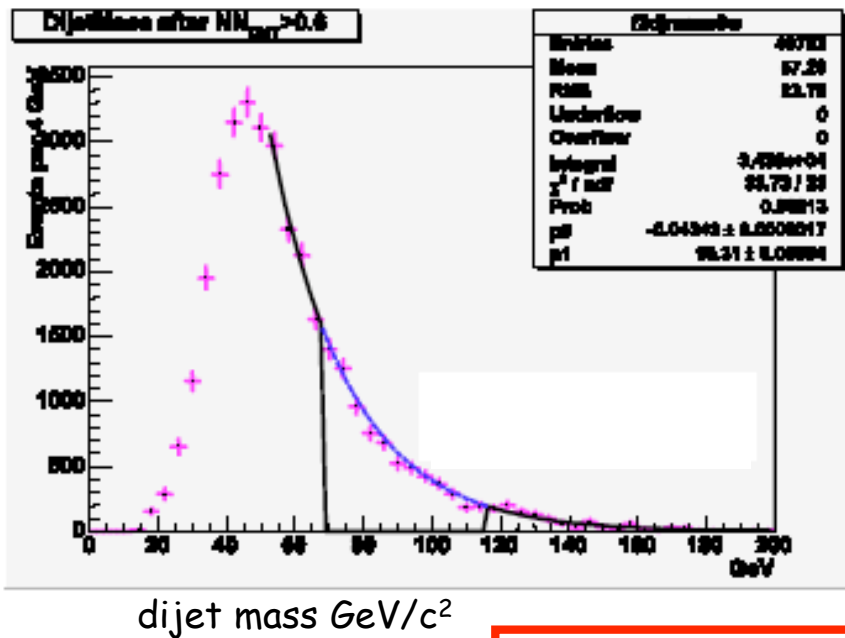
$\Delta\phi(j_1-j_2)$
W/Z γ signal
 γ dijet data
 $\Delta\phi(j_1-\gamma)$



CDF Preliminary



No $W/Z\gamma$ signal seen with 182 pb⁻¹ of data



CDF Preliminary

$$\sigma_{W(Z)\gamma} = \sigma(pp \rightarrow W\gamma) \times \mathcal{B}(W \rightarrow q\bar{q}) + \sigma(pp \rightarrow Z\gamma) \times \mathcal{B}(Z \rightarrow q\bar{q})$$

$$< 54 \text{ pb} \quad \text{at 95\% CL}$$

(SM prediction 20.5 pb)