Fundamentals of Small-Angle Neutron Scattering

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NCNR Summer School

Neutron Small Angle Scattering and Reflectometry from Submicron Structures June 5 - 9 2000



Outline

- Generalization of Atomic Scattering to a Continuum Model
 - Definitions
 - Scattering Length Density
 - Rayleigh-Gans Equation
 - Transmission
- Two Phase Systems
 - Babinet's Principle
 - Scattering Invariant
 - Porod Limit
- Multi-Phase Systems
 - Approaches to Determining Structures



Constructive interference from structures in the direction of \mathbf{q}

Diffraction length scale
$$d \approx \frac{2\pi}{q}$$

 $2\theta \approx \frac{\lambda}{d} \approx \frac{6\dot{A}}{60 \text{ to } 1000\dot{A}}$
 $2\theta \approx 0.3^{\circ} \text{ to } 5^{\circ}$

Scattering is at small angles - non-zero but smaller than classical diffraction angles





- Previously defined atomic cross sections: $\sigma_{coh}\sigma_{inc}\sigma_{abs}\sigma = 4\pi\pi^2$
- Easier to think in terms of material properties rather than atomic properties
- Define a "Scattering Length Density"

$$\rho(\vec{r}) = b_i \delta(\vec{r} - \vec{r}_i)$$

or
$$\rho = \frac{\sum_{i=1}^{n} b_i}{\overline{v}}$$

 $\overline{\mathbf{V}}$ is the volume containing the n atoms



• We can use material properties rather than atomic properties when doing <u>small-angle</u> scattering



What Length Scales Are Probed?



- Then from: $\frac{d\sigma}{d\Omega}(\vec{q}) = \frac{1}{N} \left| \sum_{i}^{N} b_{i} e^{i\vec{q}\cdot\vec{r}} \right|^{2}$
- We can replace the sum over atoms with integral over the scattering length density $\sum_{i}^{N} b_{i} \rightarrow \int_{V} \rho(\vec{r}) d\vec{r}$
- Normalizing by sample volume and introducing scattering length density:

$$\frac{\mathrm{d}\Sigma}{\mathrm{d}\Omega}(\vec{q}) = \frac{\mathrm{N}}{\mathrm{V}}\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega}(\vec{q}) = \frac{1}{\mathrm{V}}\left|\int_{\mathrm{V}} \rho(\vec{r}) \mathrm{e}^{\mathrm{i}\vec{q}\cdot\vec{r}} \mathrm{d}\vec{r}\right|^{2}$$

- "Rayleigh-Gans Equation"
- Inhomogeneities in $\rho(\vec{r})$ give rise to small angle scattering
- $\Sigma = \sigma / V$ is the "macroscopic cross section"





- This is especially true if the scattering is from
 - "countable" units:

$$\left| \int_{V} f(\vec{r}) d\vec{r} \right|^{2} \rightarrow \sum_{i}^{N} \sum_{j}^{N} f(\vec{r}_{i} - \vec{r}_{j})$$

Polymers Particulates Proteins

monomer unit per particle polypeptide subunits

• A statistical description may also be appropriate:

$$\rho(\mathbf{r}) \rightarrow \gamma(\mathbf{r})$$

Non-Particulate

correlation function



$$\frac{d\sigma}{d\Omega}(\vec{q}) = differential cross section$$

• normalize by scattering volume

$$\frac{d\Sigma}{d\Omega}(\vec{q}) = \frac{N}{V}\frac{d\sigma}{d\Omega}(\vec{q}) = \text{scattering per unit volume}$$

• Two contributions to measured signal:

$$\frac{d\Sigma}{d\Omega}(\vec{q}) = \frac{d\Sigma_{coh}}{d\Omega}(\vec{q}) + \frac{d\Sigma_{inc}}{d\Omega}$$

• Incoherent scattering is not q-dependent and contributes only to the noise level, while absorption reduces the overall signal

$$\frac{d\Sigma}{d\Omega}(\vec{q}) \propto I(q) = \frac{\text{"Scattered Intensity"}}{\text{(measured quantity)}}$$







Macroscopic Sample

• Integrating over the sample thickness:

$$\int_{\phi_0}^{\phi} \frac{\mathrm{d}\phi}{\phi} = -\Sigma \int_0^t \mathrm{d}x \qquad \ln \frac{\phi}{\phi_0} = -\Sigma \cdot t$$
$$\frac{\phi}{\phi_0} = \text{transmission} = \mathrm{e}^{-\Sigma \cdot t}$$

• Tempting to solve for optimal sample thickness:

$$t = 1/\Sigma$$
 or $T_{opt} = 0.37$

- but in practice ...
 - incoherent scattering background
 - multiple scattering
- Higher transmissions are desired



$$\rho_1$$
 ρ_2

• Incompressible phases of scattering length density ρ_1 and ρ_2

$$V = V_1 + V_2$$

$$\rho(\mathbf{r}) = \begin{cases} \rho_1 \text{ in } V_1 \\ \rho_2 \text{ in } V_2 \end{cases}$$

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Then from the Rayleigh-Gans equation:

• break the total volume into two sub-volumes

$$\frac{d\Sigma}{d\Omega}(\vec{q}) = \frac{1}{V} \left| \int_{V_1} \rho_1 e^{i\vec{q}\cdot\vec{r}} d\vec{r}_1 + \int_{V_2} \rho_2 e^{i\vec{q}\cdot\vec{r}} d\vec{r}_2 \right|^2$$
$$\frac{d\Sigma}{d\Omega}(\vec{q}) = \frac{1}{V} \left| \rho_1 \int_{V_1} e^{i\vec{q}\cdot\vec{r}} d\vec{r}_1 + \rho_2 \left\{ \int_{V} e^{i\vec{q}\cdot\vec{r}} d\vec{r} - \int_{V_1} e^{i\vec{q}\cdot\vec{r}} d\vec{r}_1 \right\} \right|^2$$

• So at non-zero q-values:

$$\frac{d\Sigma}{d\Omega}(\vec{q}) = \frac{1}{V}(\rho_1 - \rho_2)^2 \left| \int_{V_1} e^{i\vec{q}\cdot\vec{r}} d\vec{r}_1 \right|^2$$
Material Properties
+
Radiation Properties
$$\int_{V_1} e^{i\vec{q}\cdot\vec{r}} d\vec{r}_1 | d\vec{r}_1 |$$
Spatial Arrangement of Material



Babinet's Principle





Two structures give the same scattering

$$\frac{\mathrm{d}\Sigma}{\mathrm{d}\Omega}(\vec{\mathbf{q}}) \propto \left(\rho_1 - \rho_2\right)^2$$

*incoherent scattering may be different

- Contrast is relative
- Loss of phase information is $\rho_1 > \rho_2$?
- Very important in multi-phase systems
 - contrast matching / variation

Scattering Invariant



10% black / 90% white in each square

- Scattered intensity for each would certainly be different $\widetilde{Q} = \int \frac{d\Sigma}{d\Omega} (\vec{q}) d\vec{q} = (2\pi)^3 \overline{(\rho(\vec{r}) - \overline{\rho})^2}$
- For an incompressible, two-phase system:

$$\frac{Q}{4\pi} \equiv Q^* = 2\pi^2 \phi_b (1 - \phi_b) (\rho_w - \rho_b)^2$$

•Domains can be in any arrangement

Guinier & Fournet, pp. 75-81.





• At large q: $I(q) \propto q^{-4}$

$$\frac{\pi}{Q^*} \cdot \lim_{q \to \text{large}} (I(q) \cdot q^4) = \frac{S}{V}$$

• S/V = specific surface area of sample



Porod Scattering



*Glatter & Kratky pp. 30-1.



Multi-Phase Materials



• "contrast" and "structure" terms can still be factored as for 2-phase system

$$\frac{\mathrm{d}\Sigma}{\mathrm{d}\Omega}(\vec{\mathbf{q}}) \rightarrow \frac{\mathrm{d}\Sigma}{\mathrm{d}\Omega}(\mathbf{q}, \rho_{\mathrm{i}}, S_{\mathrm{ij}})$$





$$\frac{d\Sigma}{d\Omega}(q) = \sum_{i=1}^{p} (\rho_{i} - \rho_{0})^{2} S_{ii}(q) + \sum_{i < j} (\rho_{i} - \rho_{0}) (\rho_{j} - \rho_{0}) S_{ij}(q)$$

• Scattering is now a sum of several terms with possibly many unknowns $(S_{ij}'s)$

*Higgins & Benoit pp. 121-2.



Solving Multi-Phase Structures

- Contrast Matching
 - reduce the number of phases 'visible'



• The two distinct two-phase systems can be easily understood



Solving Multi-Phase Structures



- A set of scattering experiments can yield a set of equations
- of known contrasts and unknown 'partial structure functions'
- Sturhmann Analysis

Determine structure from $R_g = F(contrast)$







- Rayleigh-Gans equation
- Scattering length density
- Specialized for specific systems of interest
- Two-phase systems:
 - Relative scattering length density
 - Model independent results
- Multi-phase systems:
 - Advanced techniques
 - Control of contrast