Chameleon scalar fields and the GammeV Experiment

Amol Upadhye Fermilab Munch December 10, 2007

collaborators: A. Chou, J. Steffen, A. Weltman, W.Wester

Outline

Background and motivation Chameleon phenomenology Constraints (lab and astrophysical) Chameleon afterglow and the **GammeV** Experiment

*****Conclusions

Background: PVLAS

- Rotation in polarization direction of photons passing through a region of high magnetic field
- Could be explained by photon coupling to light scalar/pseutoscalar particle
- Recent PVLAS data: effect not reproduced

Background: Cosmology

- Accelerating cosmic expansion could be caused by light scalar field
- Scalar field must be heavy on solar system and laboratory scales to avoid fifth force constraints

Chameleon mechanism: use scalar coupling to fermions/photons to increase mass in high-density regions

Phenomenology: action

$$S = \int d^4x \sqrt{-g} \left(\frac{1}{2M_{\rm Pl}} R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{e^{2g_\gamma \phi/M_{\rm Pl}}}{4} F_{\mu\nu} F^{\mu\nu} \right) + S_m (e^{2g_m \phi/M_{\rm Pl}}, \psi_m)$$

typical matter coupling: $m(\phi) = m_0 e^{g_m \phi/M_{\rm Pl}}$

chameleon potentials:

•power law $V(\phi) = \frac{\xi}{N!} \phi^N$ •exponential $V(\phi) = \Lambda^4 \exp\left(\frac{\Lambda^n}{\phi^n}\right)$

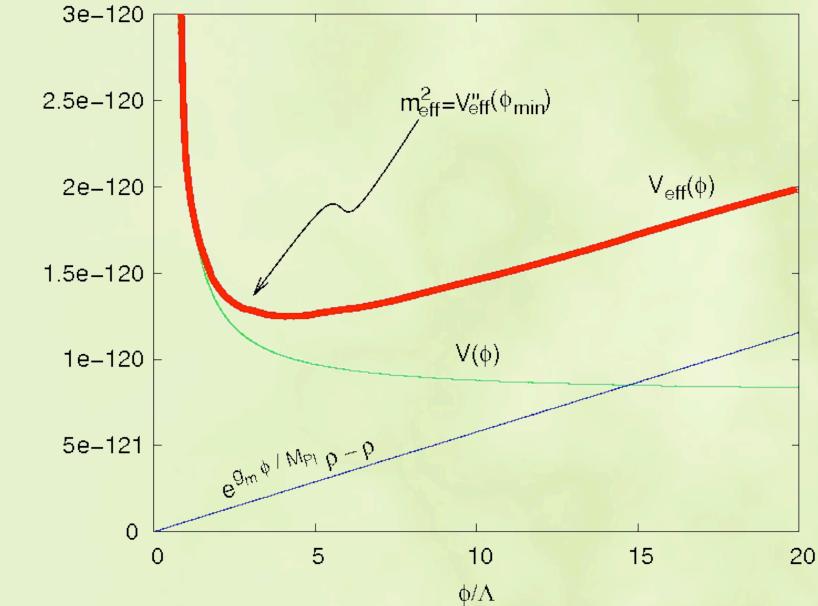
Effective potential:

 $V_{\text{eff}}(\phi, \vec{x}) = V(\phi) + e^{g_m \phi/M_{\text{Pl}}} \rho_m(\vec{x})$ $+e^{g_{\gamma}\phi/M_{\rm Pl}}\rho_{\gamma}(\vec{x})$ $\rho_{\gamma} \equiv \frac{1}{2} (B^2 - E^2)$

Effective mass:

 $m_{\rm eff}^2 = \frac{\partial^2 V_{\rm eff}}{\partial \phi^2} \bigg|_{\rm c}$

 $V_{eff}(\phi)$ [reduced Planck scale]



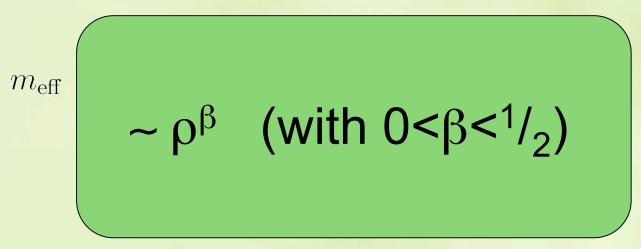
Power law potential:

$$m_{\text{eff}} = \left[\frac{(N-2)!}{\xi} \left(\frac{(N-1)g_m}{M_{\text{Pl}}}\right)^{N-1}\right]^{\frac{1}{2N-4}} \\ \times \left(\rho_m + \frac{g_\gamma}{g_m}\rho_\gamma\right)^{\frac{N-2}{2N-2}}$$

Exponential potential:

$$m_{\rm eff} = \left(\frac{(n+1)^{n+1}g_m^{n+2}}{nM_{\rm Pl}^{n+2}\Lambda^{n+4}}\right)^{\frac{1}{2n+2}} \left(\rho_m + \frac{g_\gamma}{g_m}\rho_\gamma\right)^{\frac{n+2}{2n+2}}$$

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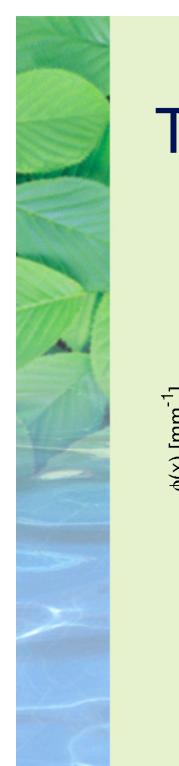
Exponential potential:

$$m_{\rm eff}$$
 ~ ρ^{β} (with $1/2 < \beta < 1$)

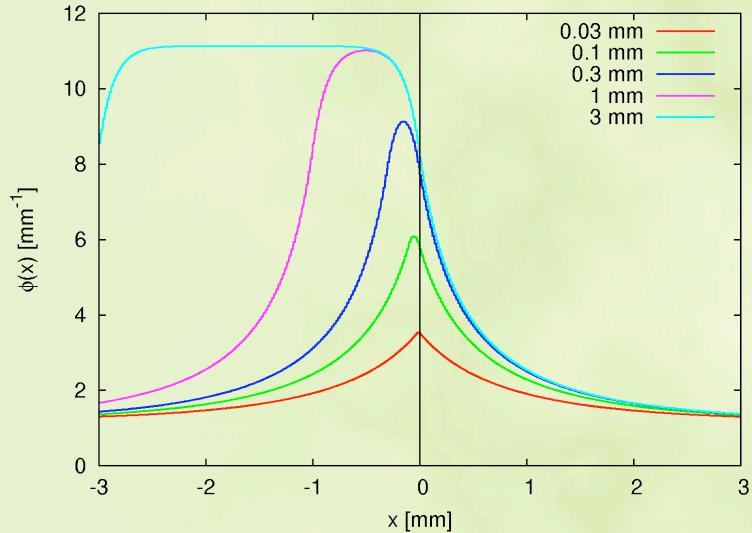
Thin shell effect

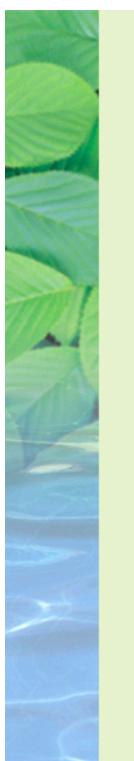
$$\frac{\partial V_{\text{eff}}}{\partial \phi} = \frac{\xi}{(N-1)!} \phi^{N-1} - \frac{g_m}{M_{\text{Pl}}} e^{-g_m \phi/M_{\text{Pl}}} \rho_m$$
$$\Rightarrow \phi_{\min} \approx \left(\frac{(N-1)!g_m \rho_m}{M_{\text{Pl}}\xi}\right)^{\frac{1}{N-1}}$$

- matter coupling g_m≠0 implies a minimum energy in bulk matter
- φ_{min} attained for objects of size R>>m_{eff}⁻¹
- field outside large object couples only to thin outer shell of size ~ m_{eff}⁻¹

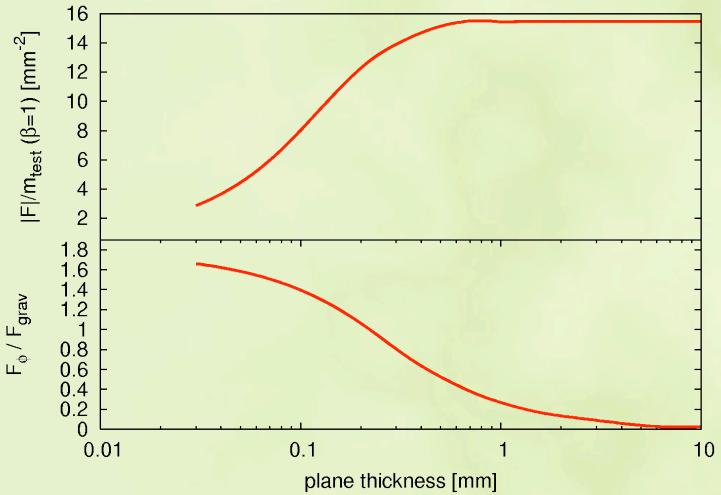


Thin shell effect



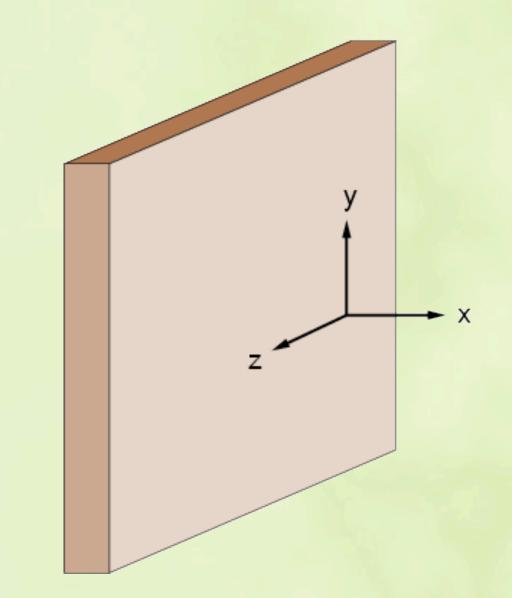


Thin shell effect





Fifth force



Fifth force

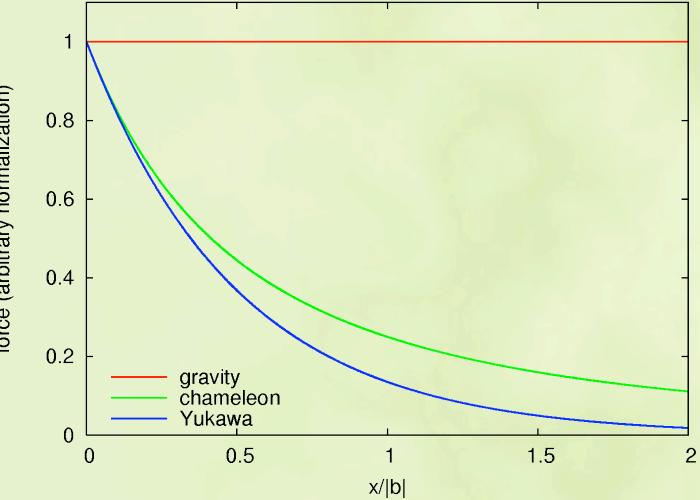
Power law chameleon:

$$\phi(x) = \left(\frac{2N!}{\xi(N-2)^2}\right)^{\frac{1}{N-2}} \frac{1}{(x+b)^{\frac{2}{N-2}}}$$
$$F_{test} = -\frac{2}{N-2} \left(\frac{2N!}{\xi(N-2)^2}\right)^{\frac{1}{N-2}} \frac{g_m m_{test}}{(x+b)^{\frac{N}{N-2}}}$$

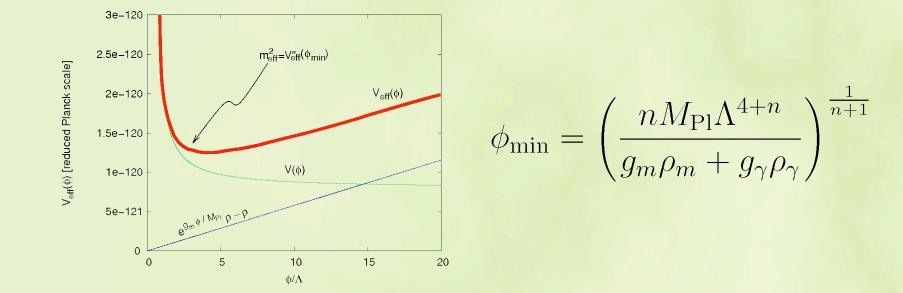
Exponential chameleon $(\phi \gg \Lambda \Rightarrow V(\phi) \approx \Lambda^4 \left(1 + \frac{\Lambda^n}{\phi^n}\right)$): $\phi(x) = \left(\frac{1}{2}(n+2)^2 \Lambda^{4+n}\right)^{\frac{1}{n+2}} (x+b)^{\frac{2}{n+2}}$



Fifth force



Chameleon fragmentation



$$\frac{\partial^m V_{\text{eff}}}{\partial \phi^m} \bigg|_{\phi=\phi_{\min}} = \frac{(-1)^m (n+m-1)! \Lambda^{4+n}}{(n-1)! \phi_{\min}^{n+m}}$$

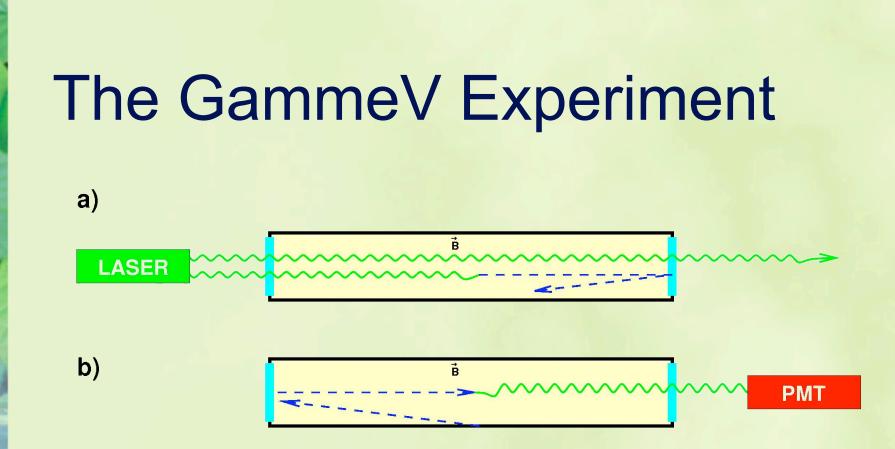
Chameleon fragmentation

$$\sigma_{\text{frag},m} \sim \left(\frac{\partial^m V_{\text{eff}}(\phi_{\min})}{\partial \phi^m}\right)^2 E_{\text{CM}}^{2(m-5)} \sim \left(\frac{E_{\text{CM}}}{\phi_{\min}}\right)^{2m}$$

so $E_{\text{CM}} > \phi_{\min} \Rightarrow \text{ fragmentation!}$

Experimental constraints

- * Laboratory: chameleon Compton wavelength $m_{eff}^{-1} < 10 \ \mu m$ at laboratory densities $\rho \sim 1 \ g/cm^3 (m_{eff} > 0.01 \ eV)$
- Solar cooling: m_{eff}(ρ) >> T_{sun}(ρ) in order to avoid energy loss by chameleon emission



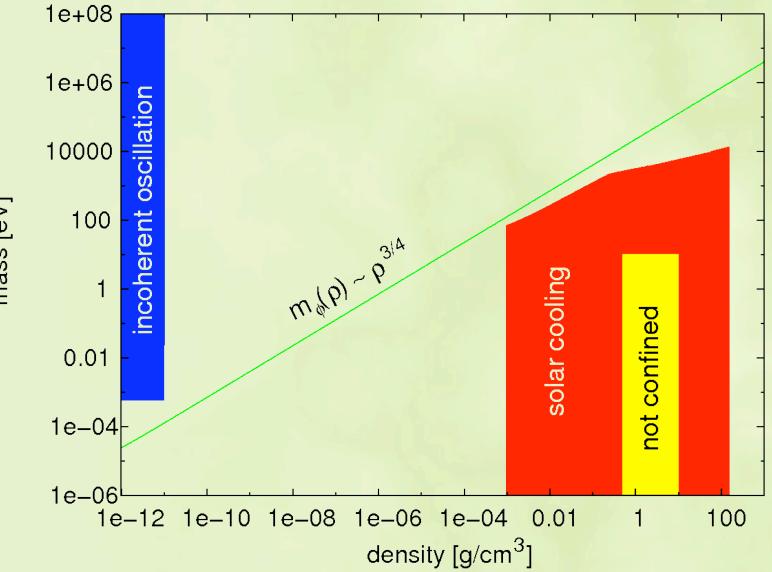
- a) Chameleon production phase: photons propagating through region of magnetic field oscillate into chameleons
- b) Afterglow phase: chameleons in chamber gradually decay back to photons and are detected

The GammeV Experiment

Coherent chameleon-photon oscillations: m_{eff} << 10⁻³ eV inside chamber

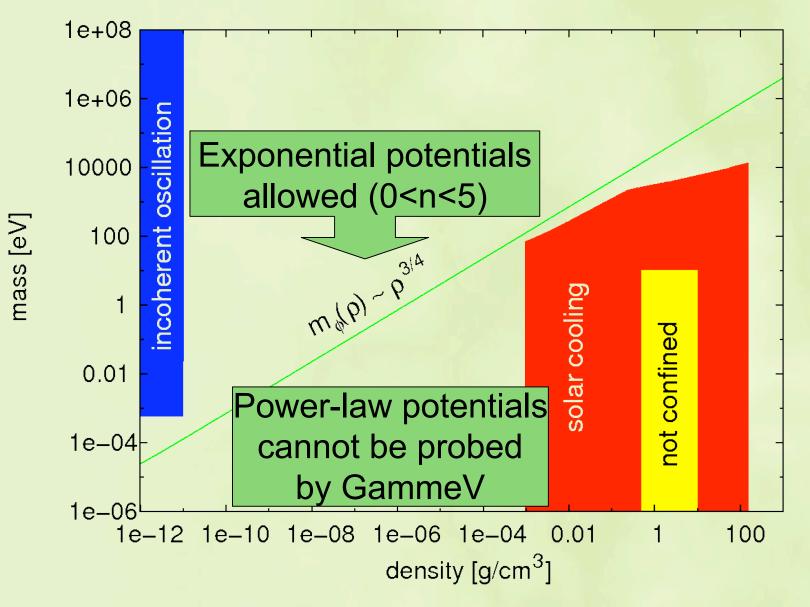
Containment: m_{eff} >> 1 eV in walls and windows of chamber (ρ~1 g/cm³) in order for chameleons to be contained

Requirements for detectability



mass [eV]

Requirements for detectability



Chameleon production

Initial state (entrance window): $|\psi(0)\rangle = |\gamma\rangle$ Photon oscillation into chameleon: $|\psi(z)\rangle = \mathcal{M}_{osc}(z) |\psi(0)\rangle = \kappa z |\phi\rangle + |\gamma\rangle$ $(\text{where } \kappa = \frac{Bg_{\gamma}}{2M_{Pl}}, \mathcal{M}_{osc}(\Delta z) = \begin{pmatrix} 1 & -\kappa \Delta z \\ \kappa \Delta z & 1 \end{pmatrix})$

* Measurement at exit window: $\mathcal{P}_{\gamma \to \phi} = |\langle \phi | \ \psi(L) \rangle|^2 = \frac{B^2 L^2 g_{\gamma}^2}{4M_{\text{Pl}}^2}$

Afterglow

Laser turned off Chameleon momentum isotropized after multiple bounces Measurement at window (z=0): photon escapes, chameleon is reflected Photon regeneration amplitude: $\langle \gamma | \psi(L) \rangle = \kappa \lambda$ geometric factor

Afterglow

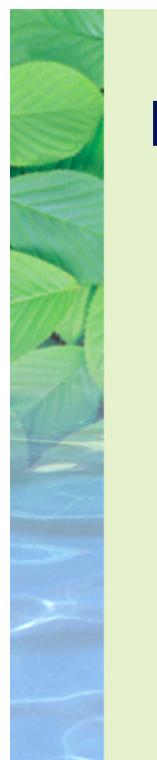
Production phase: $-\tau_{\text{prod}} < t < 0$ *****Afterglow phase: 0 < t*Number of chameleons: $\frac{dN_{\phi}}{dt} = f_{\gamma}P_{\gamma\to\phi}\Theta(-t) - \frac{N_{\phi}}{\tau_{\text{decay}}}$ where $f_{\gamma} \sim 10^{19}$ photons/sec $\tau_{\rm decay} = \frac{L}{\kappa^2 \lambda^2} = \frac{4M_{\rm Pl}^2 L}{\lambda^2 B^2 g_{\gamma}^2}$



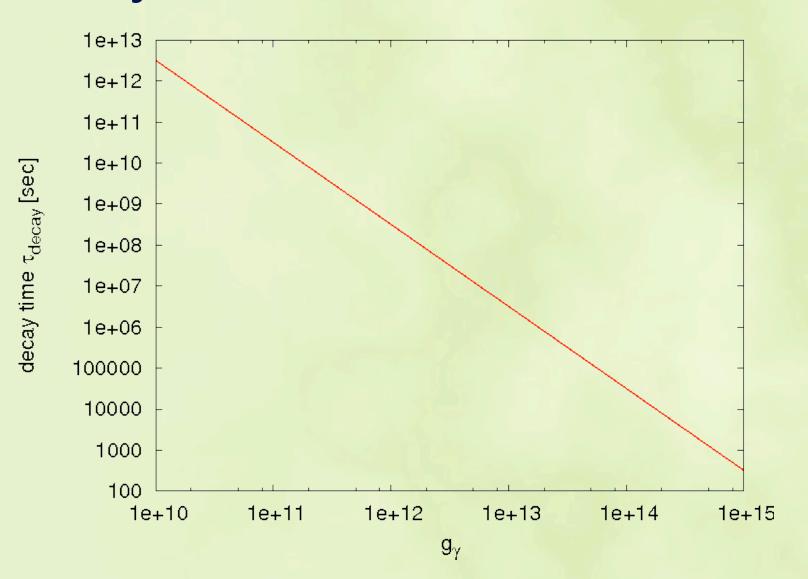
$$N_{\phi}(t) = \begin{cases} \tau_{\text{decay}} f_{\gamma} \mathcal{P}_{\gamma \to \phi} \left(1 - \exp\left(-\frac{\tau_{\text{prod}} + t}{\tau_{\text{decay}}}\right) \right) \\ \text{for } t \leq 0, \\ N_{\phi}(0) \exp\left(-\frac{t}{\tau_{\text{decay}}}\right) \\ \text{for } t > 0 \end{cases}$$

Afterglow signal:

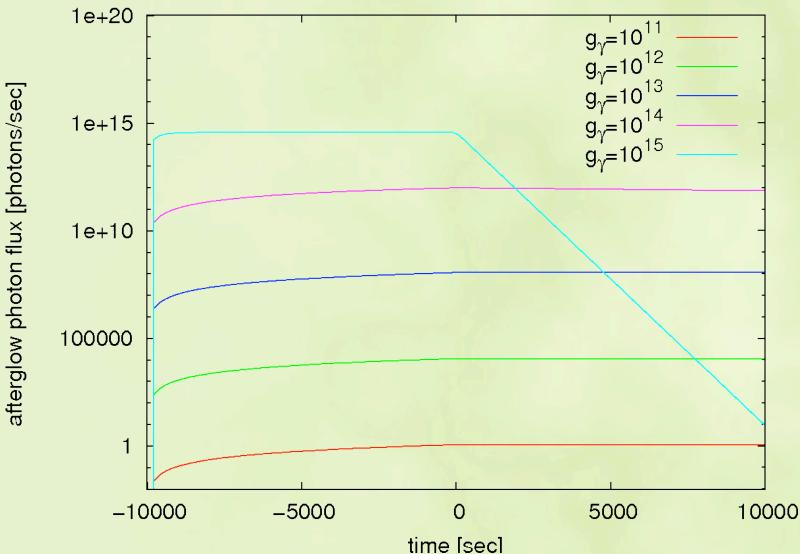
$$f_{\text{afterglow}} = \frac{N_{\phi}(t)}{\tau_{\text{decay}}} = \frac{N_{\phi}(0)}{\tau_{\text{decay}}} \exp\left(-\frac{t}{\tau_{\text{decay}}}\right)$$



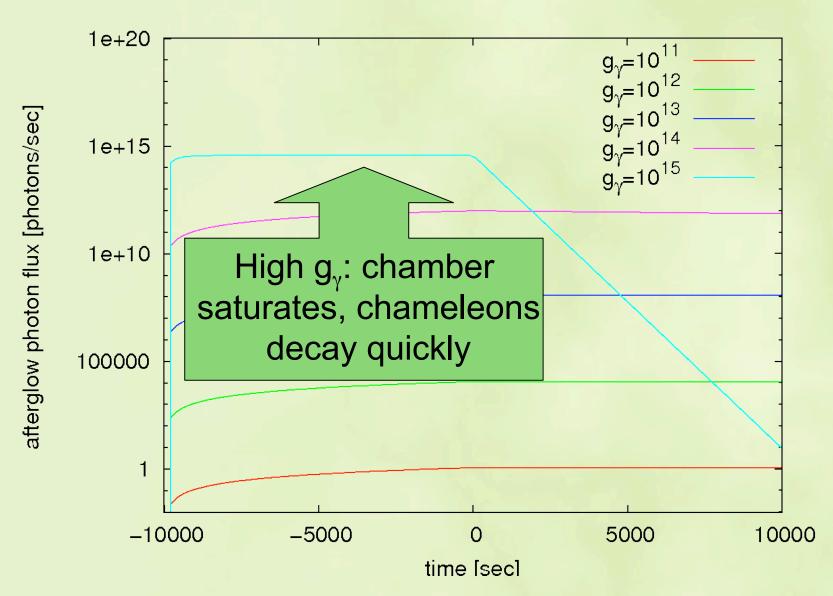
Decay time



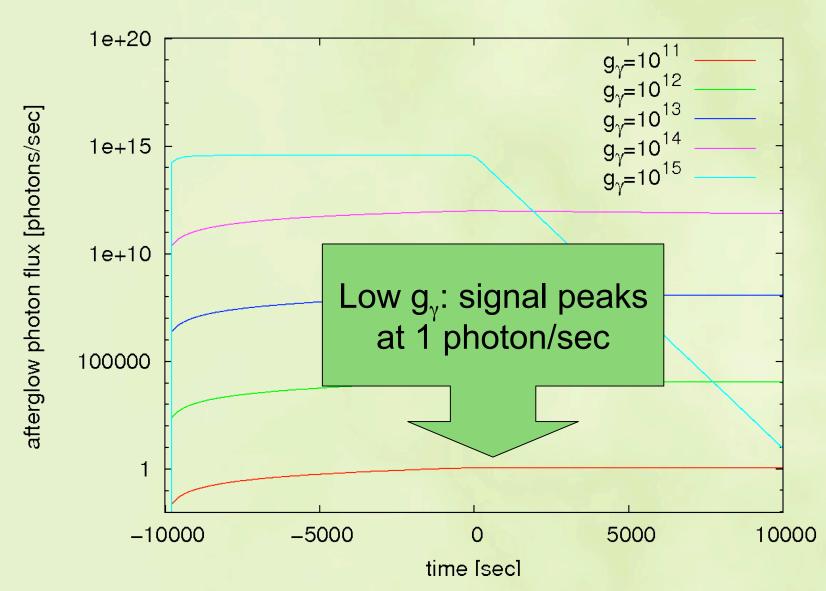
Afterglow signal dN_{\phi}/dt



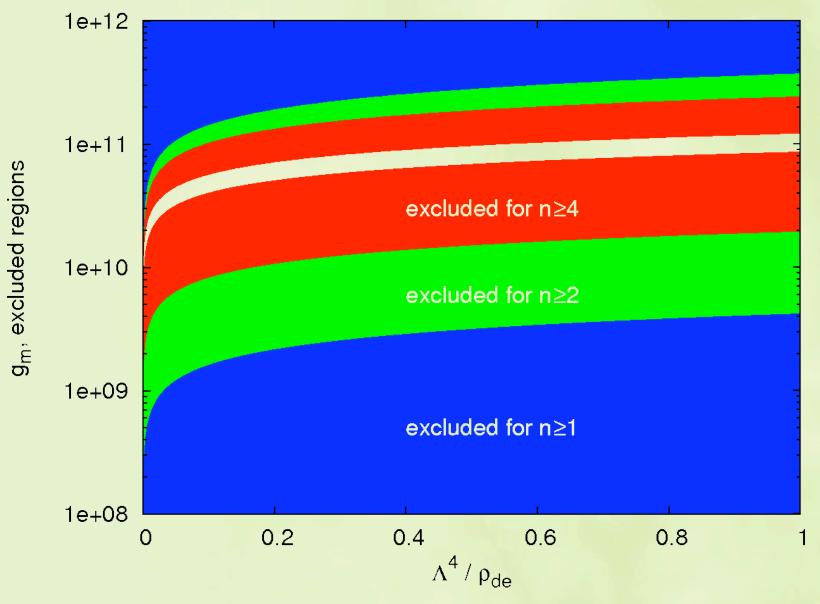
Afterglow signal dN_{\phi}/dt



Afterglow signal dN_{\phi}/dt



Constraining the parameters



Complications

In all cosmologically interesting cases, τ_{decay} >> 1 day, and any observed signal will be constant.

Other types of interactions allow smaller decay times:

- 1. decay to other particles
- 2. fragmentation

Decay to other particles

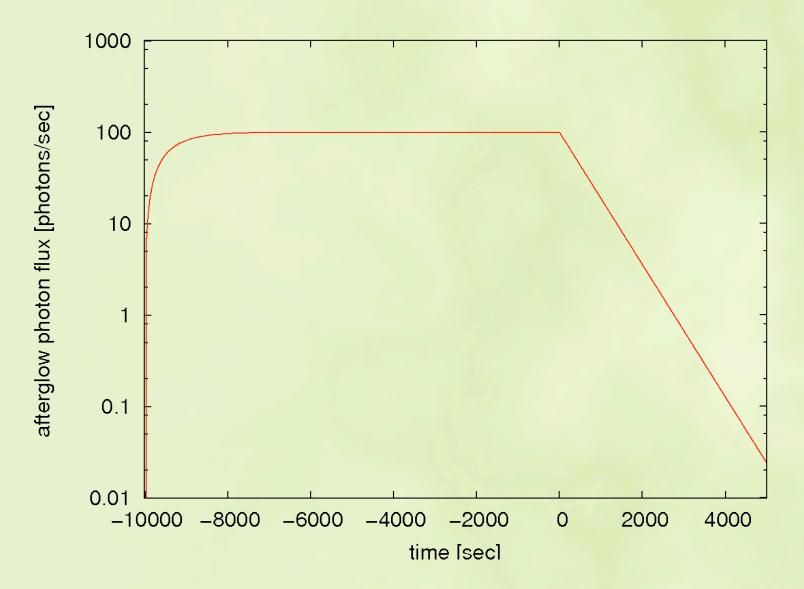
Assume that chameleons can decay to other particles not detectable by GammeV, with decay time τ_{other} .

$$\frac{dN_{\phi}}{dt} = f_{\gamma}P_{\gamma\to\phi}\Theta(-t) - \frac{N_{\phi}}{\tau_{\text{decay}}} - \frac{N_{\phi}}{\tau_{\text{other}}}$$

This alternative decay channel can speed up the decay of the afterglow signal:

$$f_{\text{afterglow}} = \frac{N_{\phi}(0)}{\tau_{\text{decay}}} \exp\left(-\frac{t}{\tau_{\text{decay}}} + \tau_{\text{other}}\right)$$

Decay to other particles



- Fragmentation causes chameleons with detectable energies (E ~ 1 eV) to split into undetectable, lower-energy chameleons.
- Exponential potential: fragmentation cross section diverges if E_{CM}>\u03c6_{min}
- In GammeV, E_{CM} ~ 1 eV.
- ***** For detectable chameleons, $\phi_{min} < 0.1 \text{ eV}$.
- * Assume σ_{frag} remains finite. Fragmentation can change the afterglow time scale.

Fragmentation turns two chameleon particles into many chameleons. Thus the rate at which the number N_{ϕ} of detectable chameleons decreases should be given by:

$$\frac{dN_{\phi}}{dt} = f_{\gamma}P_{\gamma\to\phi}\Theta(-t) - \frac{N_{\phi}}{\tau_{\text{decay}}} - \frac{N_{\phi}^2}{\tau_{\text{frag}}}$$

Production phase ($-\tau_{prod} < t < 0$):

$$N_{\phi}(t) = \frac{\tau_{\text{frag}} f_{\gamma} \mathcal{P}_{\gamma \to \phi} \sinh\left(\frac{a(t + \tau_{\text{prod}})}{\tau_{\text{frag}}}\right)}{a \cosh\left(\frac{a(t + \tau_{\text{prod}})}{\tau_{\text{frag}}}\right) + b \sinh\left(\frac{a(t + \tau_{\text{prod}})}{\tau_{\text{frag}}}\right)}$$

where

$$a = \sqrt{\left(\frac{\tau_{\text{frag}}}{2\tau_{\text{decay}}}\right)^2 + \tau_{\text{frag}} f_{\gamma} \mathcal{P}_{\gamma \to \phi}}, \qquad b = \frac{\tau_{\text{frag}}}{2\tau_{\text{decay}}}$$

Afterglow phase (0 < t):

$$N_{\phi}(t) = \frac{2bN_{\phi}(0)\exp(-t/\tau_{\text{decay}})}{2b + N_{\phi}(0) - N_{\phi}(0)\exp(-t/\tau_{\text{decay}})}$$

In the limit where fragmentation is much quicker than decay, N_{φ} is given by

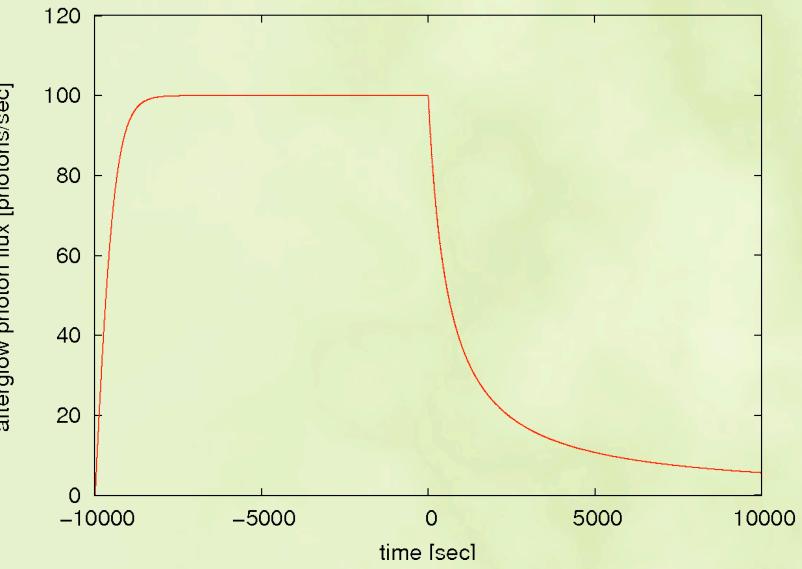
$$N_{\phi}(t) \approx \frac{N_{\phi}(0)}{1 + t/\tau_{\text{half}}}$$

where τ_{half} is the time required for the initial number of chameleons $N_{\phi}(0)$ to be halved.

The afterglow signal is:

$$f_{\rm afterglow}(t) \approx \frac{N_{\phi}(0)/\tau_{\rm decay}}{1+t/\tau_{\rm half}}$$





Conclusions

- Chameleons, self-interacting matter-coupled scalar fields, are candidates for the dark energy.
- GammeV is an experiment searching for afterglow, a unique signature of photoncoupled chameleon particles.
- GammeV can probe cosmologically interesting chameleon theories.
- Data analysis is under way!